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6. Instability Stability and Instability of Standing Waves for the Nonlinear Fractional Schrödinger Equation

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$$iu_t - (-\triangle + k^2)^s u + (\frac{1}{|x|^{\gamma}} * |u|^2)u = 0, \quad x \in \mathbb{R}^N,$$
 (1)

$$\begin{split} & u = u(t,x) : [0,T) \times \mathbb{R}^N \to \mathbb{C}, \ 0 < T \le +\infty; \\ & (-\triangle + k^2)^s u = \mathcal{F}^{-1}[(|\xi|^2 + k^2)^s \mathcal{F}[u](\xi)], \ s > 0; \\ & (\frac{1}{|x|^{\gamma}} * |u|^2)(x) = \int \frac{|u(y)|^2}{|x-y|^{\gamma}} dy; \ 0 < \gamma < 4s. \end{split}$$

 \diamondsuit Boson star, fractional quantum mechanics etc.

Known Studies

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6. Instability In particular, when N = 3, $s = \frac{1}{2}$ and $\gamma = 1$, Eq.(1) is the Boson star equation.

• Fröhlich, Lenzmann, Lewin(2007-2009): well-posedness, orbital stability, singularity of blow-up solutions.

• Bao and Dong(2011): efficient and accurate numerical methods to compute the ground states and dynamics.

Known Studies

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6. Instability For general 0 < s < 1, Laskin(2000,2002): expanding the Feynman path integral from the Brownian-like to the Lévy-like quantum mechanical paths.

• Cho, Hwang, Hajaiej and Ozawa(2012): local well-posedness, existence of blow-up solutions.

• Guo, Huo, Huang(2011-2012): global well-posedness.

Known Studies

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6. Instability Existence and stability of standing waves for the classic nonlinear Schrödinger equation (i.e. s = 1)

- Strauss(1977): existence of standing waves.
- Cazenave and Lions(1982): orbital stability of standing waves.
- Weinstein(1983): strongly instability of standing waves.

Existence and stability of standing waves for the fourth-order nonlinear Schrödinger equation (i.e. s = 2)

- Zhu, Zhang and Yang(2010): existence of standing waves and ground state.
- Fibich, Baruch and Mandelbaum(2011): numerical studies for standing waves.
- Levandosky(1998): orbital stability of standing waves.

Problems and Arguments

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\bullet Problems

1. Existence of standing waves for (1).

 ${\it 2. Stability of these standing waves.}$

• Arguments

 \diamond Weinstein (1983): blow-up argument

 $\diamondsuit\,$ Gérard (1998) Hmidi & Keraani (2005): profile decomposition

 \diamondsuit Cazenave and Lions (1982), Zhang(2000): variational argument

Local Well-posedness

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Denote $H^s := \{v \in S^{'}(\mathbb{R}^N) \mid \int (1+|\xi|^2)^s |\widehat{v}(\xi)|^2 d\xi < +\infty\}$ and

$$E(v) := \frac{1}{2} \int \overline{v} (-\Delta)^s v dx - \frac{1}{4} \int \overline{v} (\frac{1}{|x|^{\gamma}} * |v|^2) v dx.$$

Impose

$$u(0,x) = u_0 \in H^s.$$

$$\tag{2}$$

Proposition 1 (Cho, Hwang, Hajaiej and Ozawa, 2012) Let $N \ge 2, 0 < s < 1$ and $0 < \gamma \le 2s$. If $u_0 \in H^s$, then $\exists! u(t, x)$ of (1)-(2) on I = [0, T) such that $u(t, x) \in C(I; H^s) \cap C^1(I; H^{-s})$. Moreover, for all $t \in I$, u(t, x) satisfies the following conservation laws.

(i) Conservation of mass $||u(t)||_2 = ||u_0||_2$.

(ii) Conservation of energy $E(u(t)) = E(u_0)$

Profile Decomposition

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6. Instability **Proposition 2** Let $N \ge 2$ and 0 < s < 1. If $\{v_n\}_{n=1}^{+\infty}$ is a bounded sequence in H^s , then there exists a subsequence of $\{v_n\}_{n=1}^{+\infty}$ (still denoted $\{v_n\}_{n=1}^{+\infty}$), a family $\{x_n^j\}_{j=1}^{+\infty}$ of sequences in \mathbb{R}^N and a family $\{V^j\}_{j=1}^{+\infty}$ of H^s functions satisfying the following.

(i) For every $k \neq j, \, |x_n^k - x_n^j| \to +\infty$ as $n \to +\infty$.

(ii) For every $l \ge 1$ and every $x \in \mathbb{R}^N$, $v_n(x)$ can be decomposed as

$$v_n(x) = \sum_{j=1}^{l} V^j(x - x_n^j) + v_n^l(x)$$

with $\lim_{l \to +\infty} \limsup_{n \to +\infty} \|v_n^l\|_p = 0 \quad \text{ for every } p \in (2, \tfrac{2N}{(N-2s)^+}).$

Moreover, we have, as $n \to +\infty$,

$$||v_n||_2^2 = \sum_{j=1}^l ||V^j||_2^2 + ||v_n^l||_2^2 + o(1),$$

$$||v_n||^2_{\dot{H}^s} = \sum_{j=1}^l ||V^j||^2_{\dot{H}^s} + ||v_n^l||^2_{\dot{H}^s} + o(1).$$

Orbital Stability of Standing Waves

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6. Instability Let $k = 0, N \ge 2$ and M > 0. For any $0 < \gamma < 2s$, we define the following variational problem

$$d_M := \inf_{\{v \in H^s \mid \|v\|_2^2 = M\}} E(v)$$
(3)

where $E(v) := \frac{1}{2} \int |\xi|^{2s} |\hat{v}|^2 d\xi - \frac{1}{4} \int \int \frac{|v(x)|^2 |v(y)|^2}{|x-y|^{\gamma}} dx dy$ is the energy functional. Define the set

 $S_M := \{ v \in H^s | v \text{ is the minimizer of the variational problem (3)} \}.$ (4) From the Euler-Lagrange Theorem, we see that for any $v \in S_M$, there exists $\omega \in \mathbb{R}$ such that

$$(-\triangle)^{s}v + \omega v - (\frac{1}{|x|^{\gamma}} * |v|^{2})v = 0, \quad v \in H^{s}.$$

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6. Instability **Theorem 1** Let k = 0, $N \ge 2$ and M > 0. Assume 0 < s < 1 and $0 < \gamma < 2s$. Then for arbitrary $\varepsilon > 0$, there exists $\delta > 0$ such that for any $u_0 \in H^s$, if the initial data u_0 satisfies

$$\inf_{v\in S_M} \|u_0 - v\|_{H^s} < \delta,$$

then the corresponding solution u(t,x) of the Cauchy problem (1)-(2) is such that

$$\inf_{v \in S_M} \|u(t,x) - v(x)\|_{H^s} < \varepsilon$$

for all t > 0, where S_M is defined in (4).

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6. Instability Key: The existence of minimizer of Variational Problem (3) i.e. Let $N \ge 2, 0 < \gamma < 2s$ and M > 0. Then,

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$$d_M := \inf_{\{v \in H^s \mid \|v\|_2^2 = M\}} E(v) = \min_{\{v \in H^s \mid \|v\|_2^2 = M\}} E(v).$$

Firstly, the variational problem (3) is well defined. Using the Hölder inequality with $1 = \frac{\gamma}{2s} + \frac{2s - \gamma}{2s}$, we deduce that

$$\int \frac{|v(x)|^2}{|x-y|^{\gamma}} dx \le C |||x|^{-s} v(x+y)||_2^{\frac{\gamma}{s}} ||v||_2^{\frac{2s-\gamma}{s}} \le C ||v||_{\dot{H}^s}^{\frac{\gamma}{s}} ||v||_2^{\frac{2s-\gamma}{s}}.$$

Thus,

 $\int (\frac{1}{|x|^{\gamma}} * |v(x)|^2) |v(x)|^2 dx \le \| \int \frac{|v(x)|^2}{|x-y|^{\gamma}} dx \|_{\infty} \|v(y)\|_2^2 \le C \|v\|_{\dot{H}^s}^{\frac{\gamma}{s}} \|v\|_2^{\frac{4s-\gamma}{s}}.$

From Young inequality, $\exists \; C(\varepsilon,\gamma,s,\sqrt{M})>0$ such that for all $0<\varepsilon<\frac{1}{2}$

$$E(v) \ge \frac{1}{2} \|v\|_{\dot{H}^{s}}^{2} - C \|v\|_{2}^{\frac{4s-\gamma}{s}} \|v\|_{\dot{H}^{s}}^{\frac{\gamma}{s}} \ge (\frac{1}{2} - \varepsilon) \|v\|_{\dot{H}^{s}}^{2} - C(\varepsilon, \gamma, s, \|v\|_{2}).$$
(5)

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6. Instability Secondly, there exists $C_0 > 0$,

$$d_M \le -C_0 < 0.$$

Indeed, take $v_n = \rho_n^{\frac{N}{2}} R(\rho_n x)$, where $\rho_n > 0$ and R is a function such that $||v_n||_2^2 = ||R||_2^2 = M$. We deduce that

$$E(v_n) = \frac{1}{2} \int |\xi|^{2s} |\widehat{v_n}|^2 d\xi - \frac{1}{4} \int \int \frac{|v_n(x)|^2 |v_n(y)|^2}{|x-y|^{\gamma}} dx dy$$
$$= \frac{\rho_n^{2s}}{2} \int |\xi|^{2s} |\widehat{R}|^2 d\xi - \frac{\rho_n^{\gamma}}{4} \int \int \frac{|R(x)|^2 |R(y)|^2}{|x-y|^{\gamma}} dx dy.$$

Since $0 < \gamma < 2s$, we can choose $\rho_n > 0$ sufficiently small such that there exists $C_0 > 0$ such that $E(v_n) \leq -C_0 < 0$.

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6. Instability Thirdly, we prove the existence of minimizers of (3). Take the minimizing sequence $\{v_n\}_{n=1}^{+\infty} \subset H^s$ of (3) satisfying that as $n \to +\infty$,

$$E(v_n) \to d_M \text{ and } \|v_n\|_2^2 \to M.$$
 (6)

Then, for *n* large enough, $E(v_n)$ satisfies $E(v_n) < d_M + 1$. From (5), for any $0 < \varepsilon < \frac{1}{2}$ and $n \ge 1$ large enough, we deduce that

$$\left(\frac{1}{2}-\varepsilon\right)\|v_n\|_{\dot{H}^s}^2 \le d_M + 1 + C(\varepsilon,\gamma,s,M).$$

Hence, $\{v_n\}_{n=1}^{+\infty}$ is bounded in H^s . Apply Proposition 2, we see that

$$v_n(x) = \sum_{j=1}^{l} V^j(x - x_n^j) + v_n^l,$$
(7)

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Thus, we have

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6. Instability $E(v_n) = \sum_{j=1}^{l} E(V^j(x - x_n^j)) + E(v_n^l) + o(1).$ (8)

Take the scaling transformation $V_{\rho_j}^j = \rho_j V^j (x - x_n^j)$ with $\rho_j = \frac{\sqrt{M}}{\|V^j (x - x_n^j)\|_2} \ge 1$. We have $\|V_{\rho_j}^j\|_2^2 = M$. We deduce that as $n \to +\infty$ and $l \to +\infty$

$$\begin{split} d_M &\geq E(v_n) \\ &= \sum_{j=1}^l \left(\frac{E(V_{\rho_j}^j)}{\rho_j^2} + \frac{\rho_j^2 - 1}{4} \int \int \frac{|V^j(x - x_n^j)|^2 |V^j(y - x_n^j)|^2}{|x - y|^{\gamma}} dx dy \right) + E(v_n^l) + o(1) \\ &\geq \sum_{j=1}^l \frac{d_M}{\rho_j^2} + \inf_{j \geq 1} \frac{\rho_j^2 - 1}{4} \left(\sum_{j=1}^l \int \int \frac{|V^j(x - x_n^j)|^2 |V^j(y - x_n^j)|^2}{|x - y|^{\gamma}} dx dy \right) + \frac{\|v_n^l\|_2^2}{M} dx dy \\ &\geq d_M + C_0 \left(\frac{M}{\|V^{j_0}\|_2^2} - 1 \right) + o(1), \end{split}$$

where $C_0 > 0$. Therefore, there exists only one term $V^{j_0} \neq 0$ in the decomposition (7) such that $\|V^{j_0}\|_2^2 = M$. V^{j_0} is the minimizer of (3).

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6. Instability **Theorem 2** Let k = 0 and $N \ge 2$. Assume $0 < \frac{\gamma}{2} = s < 1$. Then, the ground state solitary waves $e^{i\omega t}Q(x)$ of the fractional nonlinear Schrödinger equation (1) are unstable in the following sense: For arbitrary $\varepsilon > 0$, there exist the radial initial data $\{u_{0,n}\}_{n=1}^{+\infty} \subset H^{s_0}$ with $s_0 = \max\{2s, \frac{2s+1}{2}\}$ satisfying $|x|u_{0,n} \in L^2$, $x \cdot \nabla u_{0,n} \in L^2$,

$$\|u_{0,n} - Q\|_{H^s} < \varepsilon \tag{9}$$

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and the corresponding solution $\{u_n(t,x)\}_{n=1}^{+\infty}$ blows up in the finite time, where Q is the ground state solution of

$$(-\Delta)^{s}Q + Q - (\frac{1}{|x|^{2s}} * |Q|^{2})Q = 0.$$
(10)

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Proposition 3 Let
$$N \ge 2$$
 and $0 < s < 1$. Then,

$$\int |v|^2 dx \leq \frac{2}{N} \left(\int \overline{v} x (-\triangle)^{1-s} x v dx \right)^{\frac{1}{2}} \left(\int \overline{v} (-\triangle)^s v dx \right)^{\frac{1}{2}}, \quad \forall \; v \in H^s.$$

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Proof.

For all $v \in H^s$, we have

$$\begin{split} \int |v|^2 dx &= \frac{2}{N} \int \mathcal{F}[\nabla \cdot (xv)] \mathcal{F}^{-1}[\overline{v}] d\xi \\ &= \frac{2}{N} \int \overline{\mathcal{F}[v]} \xi \cdot \nabla_{\xi} \mathcal{F}[v] d\xi \\ &\leq \frac{2}{N} \int |\xi|^s |\overline{\mathcal{F}[v]}| |\xi|^{1-s} |\nabla_{\xi} \mathcal{F}[v]| d\xi \\ &\leq \frac{2}{N} \left(\int \overline{v} (-\Delta)^s v dx \right)^{\frac{1}{2}} \left(\int \overline{v} x (-\Delta)^{1-s} x v dx \right)^{\frac{1}{2}}. \end{split}$$

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6. Instability **Proposition 4** (see [Cho, Hwang, Kwon and Lee, 2012]) Let $k = 0, N \ge 2, 0 < s < 1$ and $\gamma = 2s$. Assume that $u_0 \in H^{s_0}$ with $s_0 = \max\{2s, \frac{\gamma+1}{2}\}$ is radial symmetric, and $|x|u_0 \in L^2$ and $x \cdot \nabla u_0 \in L^2$. If u(t, x) is the solution of the Cauchy problem (1)-(2), then for all $t \in I$ (the maximal time interval), $\int \overline{u}x(-\Delta)^{1-s}udx$ is nonnegative and

 $\int \overline{u}x(-\Delta)^{1-s}xudx \le 2sE(u_0)t^2 + Ct + C.$

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6. Instability Let $\{c_n\}_{n=1}^{+\infty} \subset \mathbb{C} \setminus \{0\}$ be such that $|c_n| > 1$ and $\lim_{n \to +\infty} |c_n| = 1$, and $\{\rho_n\}_{n=1}^{+\infty} \subset \mathbb{R}^+$ be such that $\lim_{n \to +\infty} \rho_n = 1$. We take the initial data

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$$u_{0,n} = c_n \rho_n^{\frac{N}{2}} Q(\rho_n x),$$

where Q is the ground state solution of (10). We see that for all $n \ge 1$, $||u_{0,n}||_2 > ||Q||_2$ and

$$\lim_{n \to +\infty} \|u_{0,n}\|_2 = \lim_{n \to +\infty} |c_n| \|Q\|_2 = \|Q\|_2$$

and

$$\lim_{n \to +\infty} \|u_{0,n}\|_{\dot{H}^s} = \lim_{n \to +\infty} |c_n| \ \rho_n^s \ \|Q\|_{\dot{H}^s} = \|Q\|_{\dot{H}^s}.$$

Hence, from the Brézis-Lieb Lemma, $\forall \varepsilon > 0$, $||u_{0,n} - Q||_{H^s} < \varepsilon$ as n is sufficiently large.

On the other hand, we see that

n

$$E(u_{0,n}) = \frac{(|c_n|^2 - |c_n|^4)\rho_n^{2s}}{2} \|Q\|_{\dot{H}^s}^2 \le -C_0 < 0.$$

Then,

$$\int \overline{u_n} x (-\Delta)^{1-s} x u_n dx \le -2sC_0 t^2 + Ct + C$$

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which implies that there exists $0 < T < +\infty$ such that

$$\lim_{t \to T} \int \overline{u_n} x (-\Delta)^{1-s} x u_n dx = 0.$$

Finally, using the conservation of mass and applying the inequality in Proposition 3 to u_n , we see that for all time t

$$\begin{aligned} \|u_{0,n}\|_{2}^{2} &= \|u_{n}(t)\|_{2}^{2} \quad \leq \frac{2}{N} \left(\int \overline{u_{n}} x(-\Delta)^{1-s} x u_{n} dx\right)^{\frac{1}{2}} \left(\int \overline{u_{n}} (-\Delta)^{s} u_{n} dx\right)^{\frac{1}{2}} \\ &\leq \frac{2}{N} \left(\int \overline{u_{n}} x(-\Delta)^{1-s} x u_{n} dx\right)^{\frac{1}{2}} \|u_{n}(t)\|_{\dot{H}^{s}}. \end{aligned}$$

Therefore, there exists $0 < T < +\infty$ such that

$$\lim_{t \to T} \|u_n(t)\|_{\dot{H}^s} = +\infty.$$

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6. Instability Thanks!