Langevin equations for landmark image registration with uncertainty

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Image Registration with Landmarks

- Define a set of landmarks (automatically or manually) on two or more images (reference q_i^R and template q_i^T)
- Find deformation (warp) ϕ such that $\phi(\mathbf{q}_i^R) = \mathbf{q}_i^T$
- Apply to rest of image

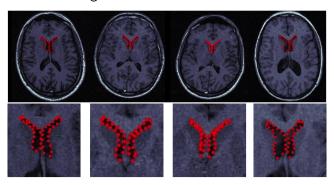


Image Registration with Landmarks

- Many choices of ϕ , usually desire a diffeomorphism
- Can define a velocity field $v(\mathbf{q},t)$ that takes the landmarks from \mathbf{q}_i^R at t=0 to \mathbf{q}_i^T at t=1
- Expand $v(\mathbf{q}, t) = \sum_{i=1}^{n} \mathbf{p}_i(t) G(\mathbf{q}, \mathbf{q}_j(t))$

Image Registration with Landmarks

Leads to reduced optimisation problem

$$\min \int_0^1 \sum_{i=1}^n \mathbf{p}_i(t) \mathbf{p}_j(t) G(\mathbf{q}_i(t), \mathbf{q}_j(t)) dt$$

with velocity constraints

$$\frac{d\mathbf{q}_i}{dt} = \sum_{i=1}^n \mathbf{p}_i(t) G(\mathbf{q}_i(t), \mathbf{q}_j(t))$$

and landmark constraints

$$\mathbf{q}_i(0) = \mathbf{q}_i^R$$
 and $\mathbf{q}_i(1) = \mathbf{q}_i^T$.

- Usually take G to be Gaussian: $G(\mathbf{q}_i, \mathbf{q}_j) = \exp(-\|\mathbf{q}_i \mathbf{q}_j\|^2/l^2)$
- Finite dimensional in space

Optimal Control Formulation

In Hamiltonian formulation we are solving a boundary-value problem:

$$\frac{d\mathbf{q}}{dt} = \nabla_p H(\mathbf{q}, \mathbf{p}),
\frac{d\mathbf{p}}{dt} = -\nabla_q H(\mathbf{q}, \mathbf{p})$$

for

$$H = \frac{1}{2} \sum_{i,j} \mathbf{p}_i^T \mathbf{p}_j G(\mathbf{q}_i, \mathbf{q}_j),$$

with boundary conditions

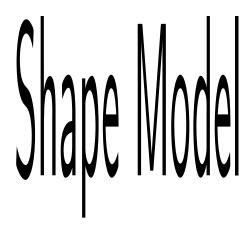
$$\mathbf{q}_i(0) = \text{reference}, \mathbf{q}_j(1) = \text{target}.$$

Defines a shooting method, need appropriate time-stepping method (Euler) and initial momenta $\mathbf{p}(0)$.

Uncertainty in Landmarks

- Two major sources of uncertainty
 - landmark positions
 - velocity field far from landmarks
- How can you incorporate this uncertainty?

Uncertainty in Landmarks



Langevin Equation

- Formulation is based on conservation of energy (the Hamiltonian)
- · Could enclose system in a heat bath
 - Langevin equation
 - constant temperature in equilibrium
- Classically:
 - system of SDEs for potential energy *V*:

$$d\mathbf{q}_{i} = \mathbf{p}_{i}dt$$

$$d\mathbf{p}_{i} = [-\lambda \mathbf{p}_{i} - \nabla_{\mathbf{q}_{i}}V]dt + \sigma d\mathbf{W}_{i}(t),$$

where $W_i(t)$ is *iid* \mathbb{R}^2 Brownian motion.

• This has invariant measure $\exp(-\beta H)$ where $H = \frac{1}{2} \sum_i \|\mathbf{p}_i\|^2 + V(q)$ and $\beta = 2\lambda/\sigma^2$ is the inverse temperature

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Noisy Landmark Registration

The generalised version is:

$$d\mathbf{q}_{i} = \nabla_{\mathbf{p}_{i}}Hdt$$

$$d\mathbf{p}_{i} = [-\lambda\nabla_{\mathbf{p}_{i}}H - \nabla_{\mathbf{q}_{i}}H]dt + \sigma d\mathbf{W}_{i}(t).$$

- Gibbs distribution still invariant
- For large β , system is concentrated at small H
- So solve this system subject to

$$\mathbf{q}_i(0) = \mathbf{q}_i^R + \text{noise}, \mathbf{q}_i(1) = \mathbf{q}_i^T + \text{noise}$$

and $\mathbf{p}_i(0) \sim \exp(-\beta H(\mathbf{q}_i^R, \cdot))$. Noise modelled as zero mean, *iid* Gaussian

For H separable this was studied by Hairer et al., 2011

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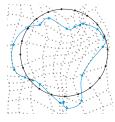
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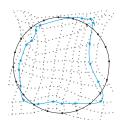
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Inverse Temperature Effect

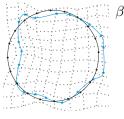




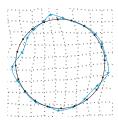
$$\beta = 20$$



$$\beta = 40$$

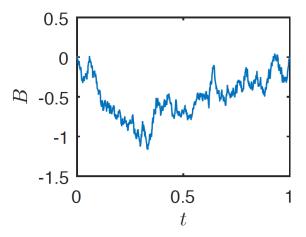






Brownian Bridges

- B is a Brownian Bridge if its distribution is W(t) conditioned by W(T) = b
- No noise on boundary conditions
- Example of a Brownian bridge conditioned on W(1)=0



Brownian Bridges

- Brownian motion is a Gaussian random field, defined by mean and covariance
- Can condition this on observations, get a second Gaussian random field with well-known formulae for mean and covariance
- Covariance does not depend on b

- Bridges very useful for parameter estimation for SDEs with discrete data
- Diffusion process X(t) with SDE

$$dX = f(X)dt + \sigma(X)dW(t), X(0) = a, t \in [0, T]$$

Bridge process satisfies a modified SDE

$$dX = \tilde{f}(X, t)dt + \sigma(X)dW(t), t \in [0, T]$$

where $\tilde{f}(x,t) = f(x) + [\sigma(x)\sigma^T(x)]\nabla \log p(t,x;T,b)$, where p is the transition density

- This is the Doob *h*-transform
- But hard to find p
- Change of measure relative to a linear diffusion bridge. Girsanov formula gives the density and hence a formula for rejection sampling

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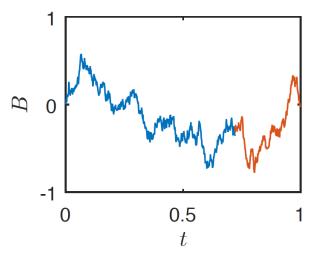
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- Simulate forward and backward, and reject samples until they cross
- Approximates the diffusion bridge
- Uses a Metropolis-Hastings iteration



 Dimensionality problem: paths less likely to cross as number of dimensions increases. Can sometimes find clever coupling of forward and backward Brownian motions.

- If H is separable, can replace H by $\frac{1}{2} \sum_{i} \|\mathbf{p}_{i}\|^{2}$
- Then have a linear SDE
- We are in high dimensions and have a non-separable Hamiltonian

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Approximation

- One good thing: Langevin trajectories are close to the Hamiltonian ones since have short time interval and small perturbation
- So let's approximate:

Approximation I: Linearisation

- Linearise about solution of Hamiltonian BVP
- Can discretise resulting linear SDE and compute covariance structure.
- Solve Hamiltonian Boundary Value Problem. Use mean position of uncertain landmarks
- 2 Linearise Langevin equation about this solution
- 3 For initial data at $t=\frac{1}{2}$, propagate forward and backward to find a Gaussian prior distribution, i.e., compute mean and covariance for process $(\mathbf{q}_i(t),\mathbf{p}_i(t))$
- Use explicit Euler
- Condition on boundary data (Euler-Maruyama)

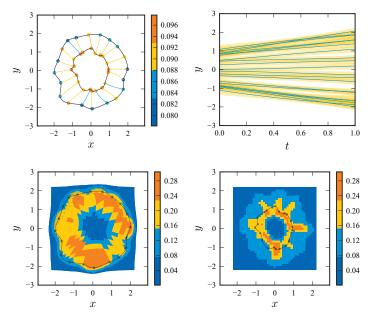
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Results: Linearisation



Approximation II: Operator Splitting

• Write generator $\mathcal{L} = \mathcal{L}_0 + \sigma^2 \mathcal{L}_1$, where

$$\mathcal{L}_{0} = H_{p} \frac{\partial}{\partial q} - H_{q} \frac{\partial}{\partial p} \text{(Liouville operator)}$$

$$\mathcal{L}_{1} = -\frac{\beta}{2} H_{p} \frac{\partial}{\partial p} + \frac{1}{2} \frac{\partial^{2}}{\partial p^{2}}$$

Both have Gibbs invariant and a BCH formula.

Approximation II: Operator Splitting

Two choices of Strang splitting:

$$e^{L^*} \approx e^{\sigma^2 L_1^*/2} e^{L_0^*} e^{\sigma^2 L_1^*/2}$$

- Gibbs distribution on initial data
- Hamiltonian flow
- Posterior is not Gaussian, but can compute MAP estimate (Laplace approximation)

$$e^{L^*} \approx e^{L_0^*/2} e^{\sigma^2 L_1^*} e^{L_0^*/2}$$

- Gibbs distribution on midpoint
- Hamiltonian flow
- Parameterised by one set of landmarks and two sets of momenta
- Can compute averages of sets of landmarks

Approximation II: Operator Splitting

Lead to different semigroups:

$$[\textbf{p}(0),\textbf{q}(0)] \underset{e^{\sigma^2 L_1^*/2}}{\longmapsto} [\textbf{p}(1/2),\textbf{q}(0)] \underset{e^{L_0^*}}{\longmapsto} [\tilde{\textbf{p}}(1/2),\textbf{q}(1)] \underset{e^{\sigma^2 L_1^*/2}}{\longmapsto} [\textbf{p}(1),\textbf{q}(1)].$$

$$[\mathbf{p}(0),\mathbf{q}(0)] \underset{e^{L_0^*/2}}{\longmapsto} [(\mathbf{p}(1/2),\mathbf{q}(1/2)] \underset{e^{\sigma^2L_1^*}}{\longmapsto} [(\tilde{\mathbf{p}}(1/2),\mathbf{q}(1/2)] \underset{e^{L_0^*/2}}{\longmapsto} [\mathbf{p}(1),\mathbf{q}(1)].$$

Time-half steps governed by Ornstein-Uhlenbeck SDE

$$d\vec{p} = -\lambda \nabla_{\mathbf{p}} H(\mathbf{p}, \mathbf{q}_0) dt + \sigma d\mathbf{W}(t), \qquad \mathbf{p}(0) = \mathbf{p}_0,$$

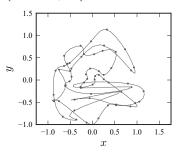
Laplace Approximation

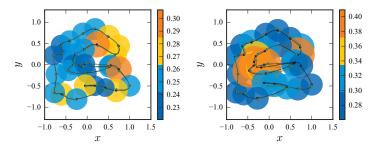
- Approximate posterior covariance max by second-order approximation to likelihood at MAP point
- Unconstrained optimisation by Gauss-Newton

For a change of optimisation function, can construct average over sets of flows (where $(\mathbf{p}^*, \mathbf{q}^*)$ is the average)

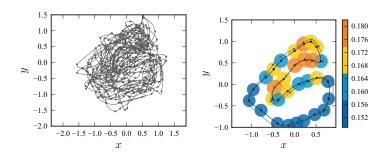
$$F(\mathbf{p}^*, \mathbf{q}^*, \mathbf{p}^j) := \beta H(\mathbf{p}^*, \mathbf{q}^*) + \frac{\beta}{4\lambda} \sum_{j=1}^J ||\mathbf{p}^j - e^{-\lambda \mathcal{G}(\mathbf{q}^*)} \mathbf{p}^*||^2 + \frac{1}{2\delta^2} \sum_{j=1}^J ||\mathbf{q}^j - S_q(1/2; 0, [\mathbf{p}^j, \mathbf{q}^*])||^2.$$

Results: Averaging (2 sample)

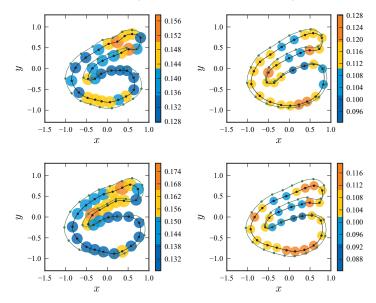




Results: MAP average (16 samples)



Results: MAP average (64 and 256 samples)



Conclusions

- Introduced Langevin equation to define prior distribution on set of diffeomorphisms
- Computationally difficult to sample the diffusion bridge
- Hence, have linearised the equations, and also made an operator splitting
- Second approach enables averaging of multiple landmark sets