



Vortices in superconductors: V. 3D mesoscopic structures

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Sphere

Expansion method

Consider only the first GL-equation (i.e. $\kappa \gg 1$)

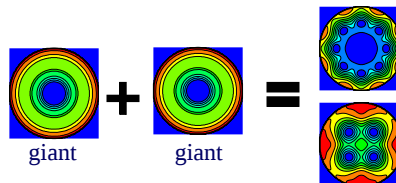
$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial \Psi}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Psi}{\partial \varphi^2} - 2is \frac{\partial \Psi}{\partial \varphi} - s^2 \rho^2 \Psi + \frac{\partial^2 \Psi}{\partial z^2} = - \left(1 - \frac{T}{T_c} \right) \Psi$$

$$\Psi(\rho, \varphi, z) = e^{iL\varphi} \psi(\rho, z)$$

$$\left(-i\hbar \vec{\nabla} - \frac{2e}{c} \vec{A} \right)_n \Psi = 0$$

$\Rightarrow$ 2D differential equation

$$\psi(r, \varphi) = \sum_{L_j=0}^L \sum_n C_{n,L_j} f_{n,L_j}(r) \exp(iL_j \varphi)$$



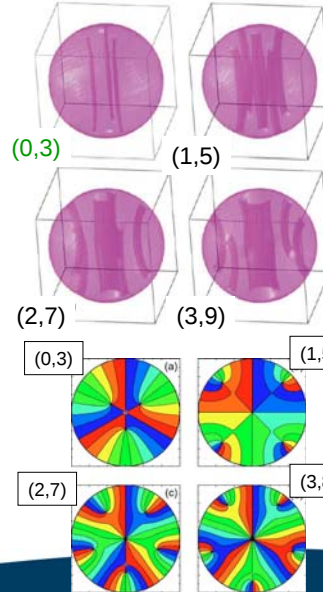
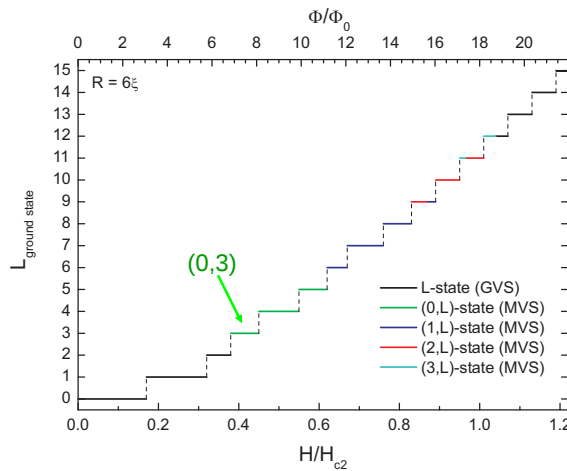
B.J. Baelus, D. Sun and F.M. Peeters, Phys. Rev. B **75**, 174523 (2007)

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Ground state configurations



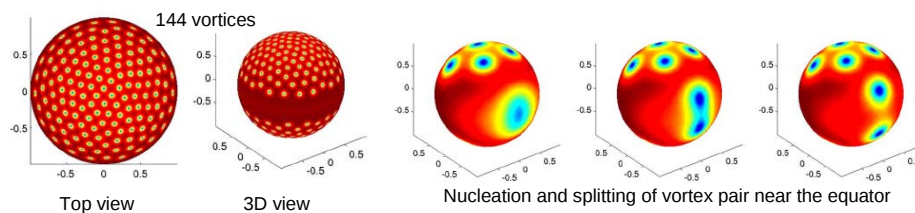
Hollow sphere

Numerical simulations of the quantized vortices on a thin superconducting hollow sphere

Qiang Du, Lili Ju

Journal of Computational Physics 201 (2004) 511–530

⇒ Numerical solution of the first GL equation (thin sphere shell approximation)



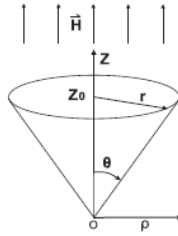
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Cone

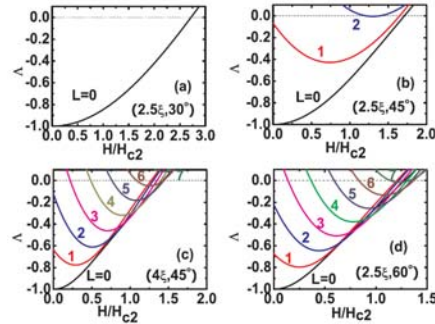
Y. Chen (UA)

M.M. Doria (Rio de Janeiro)



$$\Psi(\rho, \varphi, z) = \sum_{L_j=0}^L C_{L_j} \psi_{L_j}(\rho, z) \exp(i L_j \varphi).$$

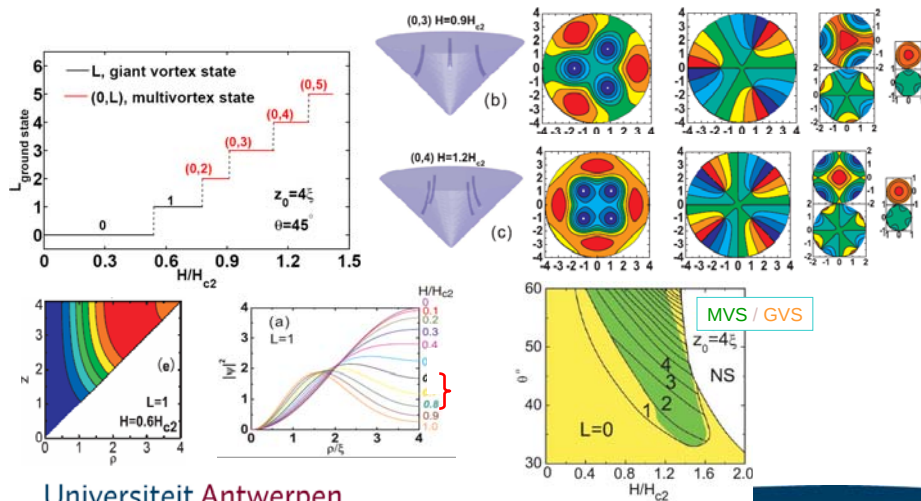
$$\frac{\partial F}{\partial C_{L_j}} = 0 \quad \text{and} \quad \frac{\partial F}{\partial C_{L_j}^*} = 0$$



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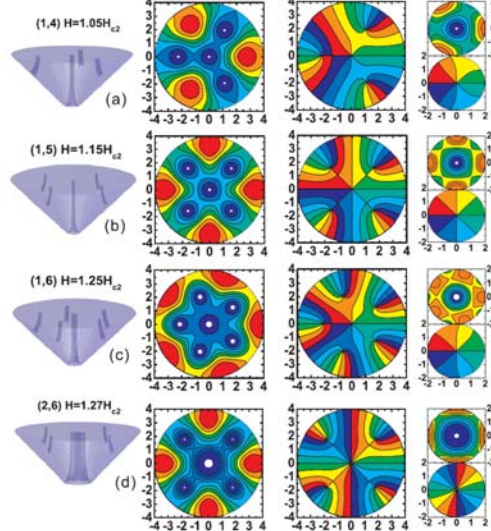
Phase diagram



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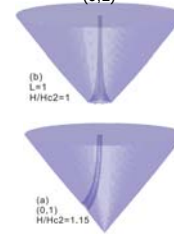


$(z_0=4\xi, \theta=45^\circ)$



Bend vortices

$(z_0=4\xi, \theta=38^\circ)$

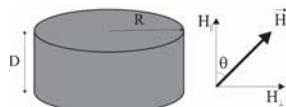


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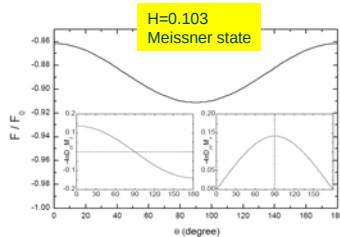
Tilted vortices in cylinder



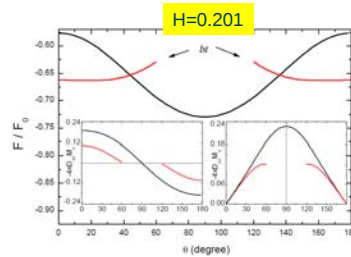
Numerical solution of first GL-equation (i.e. $\lambda \gg 1$)

$R=4\xi; D=4\xi$

Magnetic field units: $H_{ref} = H_{c3//}$, max. field of short cylinder for || orientation



H=0.103
Meissner state



H=0.201

Meissner state



59°

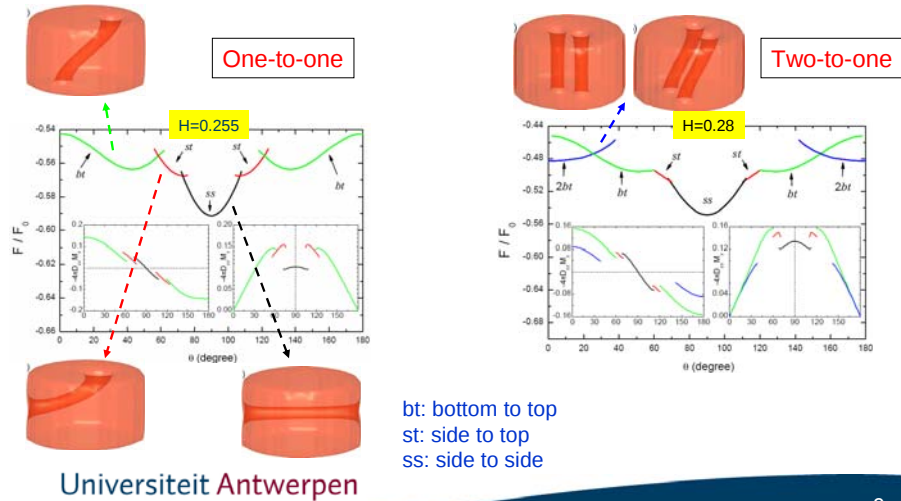
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A.R. de C. Romaguera, M.M. Doria, and F.M. Peeters, Phys. Rev. B 75, 184525 (2007)

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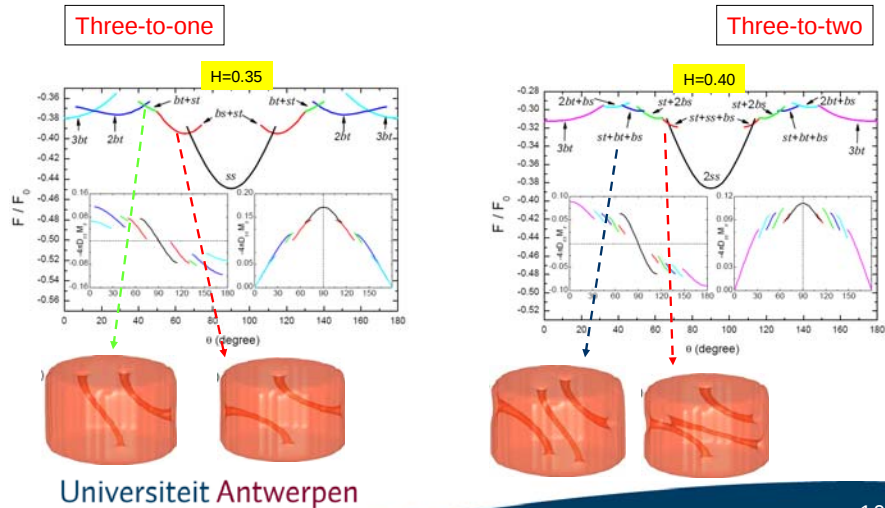
Tilted vortices in cylinder



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Tilted vortices in cylinder



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Transverse magnetization and torque in asymmetrical mesoscopic superconductors

Asymmetrical coexistence of vortices with different lengths $\rightarrow$ transverse magnetization
 $\rightarrow$ torque

Sphere with $R=6.0\xi$ - lacking one of his quadrants

$\mu=VM$: magnetic moment of a sample with volume V

$\tau=\mu \times H$: torque

$\tau=\mu H \sin\theta \mathbf{e}_\perp$: where θ is the angle between $\mathbf{H}$ and $\mathbf{M}$



For example:

$\xi=100\text{nm} \rightarrow H_{c2} = \Phi_0 / 2\pi\xi^2 = 300\text{G}$

If $R=0.6\mu\text{m}$ then transverse magnetization $M_x = 10^{-3} H_{c2} = 0.3\text{G}$,

while the longitudinal is 10 times larger: $M_z = 10^{-2} H_{c2} = 3\text{G}$

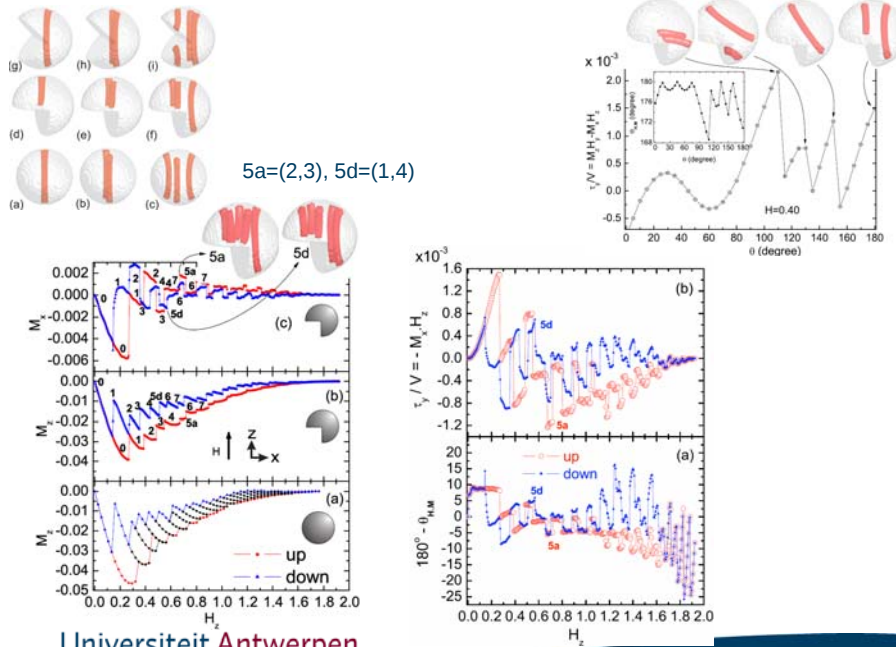
The magnetic moment of this sphere is: $\mu = M_z V = 7.0 \cdot 10^{-12} \text{emu}$

and the torque $\tau_y = 4.0 \cdot 10^{-12} \text{erg}$

A.R. de C. Romaguera, AmM. Doria, and F.M. Peeters, Phys. Rev. B **76**, 020505(R) (2007)

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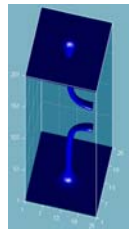
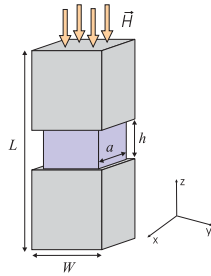


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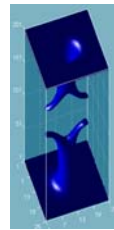
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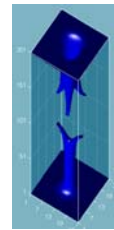
SC wire with a constriction



(1,0,1)



(2,0,2)



(3,0,3)

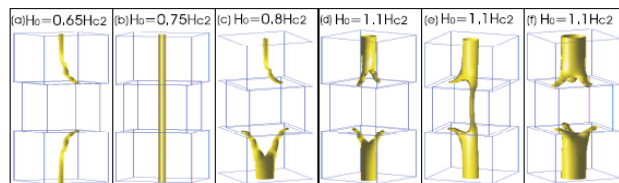
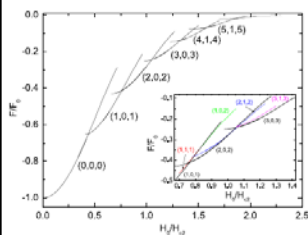
A. K. Elmurodov, D. Y. Vodolazov, and F. M. Peeters, *Europhys. Lett.*, **74**, 151 (2006)

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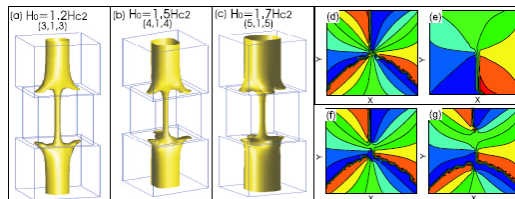


Vortex states



(1,0,1) (1,1,1) (1,0,2) (2,0,2) (2,1,2) (3,0,3)

Spatial break-up of a GVS:
 Inside thick wire:
 GVS(L=3)
 near constriction $\rightarrow$ GVS(L=2)+1V
 Inside constriction
 $\rightarrow$ 1 vortex



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Conclusion

- 3D structures: vortices are 3D objects (tubes instead of disks)
- Added possibility to manipulate the vortices in the third dimension:
 - bend vortice → see e.g. sphere, cone, cylinder in a tilted magnetic field
 - spatial break-up of a GVS (splitting of a vortex) → see wire with a constriction
 - Vortices can exert a torque on an asymmetric sample