## Continuation in dimension

Recall that a Q-curvature $\mathbf{Q}$ in even dimension $n$ has

$$
\widehat{\mathbf{Q}}=\mathbf{Q}+\mathbf{P} \omega
$$

where $\mathbf{P}: \mathcal{E} \rightarrow \mathcal{E}[-n]$ is an operator of the form

$$
\mathbf{P}=\delta\left\{(d \delta)^{n / 2-1}+\mathrm{LOT}\right\} d
$$

$\mathbf{P}$ is automatically conformally invariant.

Let's do the original Q-curvature construction in the case $n=2$, where $\mathbf{Q}=\mathrm{J}$. This is really the derivation of the Gauss curv. presc. eqn. from the higher-dimensional Yamabe eqn. The packaging in terms of stabilization (inflating the dimension by taking the product with flat tori) was suggested by Robin Graham. The original treatment inflates the dimension in an invariant-theoretic sense.

Start with a Riemannian 2-manifold ( $M, g$ ), and take the product with a standard flat $m$-torus ( $T, k$ ) for $m=0,1,2, \cdots$. Then change the metric

$$
g+k=\left(\begin{array}{cc}
g_{a b} & 0 \\
0 & k_{\alpha \beta}
\end{array}\right)
$$

conformally, using a conformal factor $\omega$ that depends only on the $M$-parameter $x$.

The Yamabe equation

$$
\left(\Delta+\frac{n-2}{2} \mathrm{~J}\right) e^{(n-2) \omega / 2}=\frac{n-2}{2} \widehat{\mathrm{~J}}^{(n+2) \omega / 2}
$$

in our situation says that

$$
\begin{aligned}
& \left\{\Delta[g+k]+\frac{m}{2} \mathrm{~J}[g+k]\right\} e^{m \omega / 2}= \\
& \frac{m}{2} \mathrm{~J}\left[e^{2 \omega}(g+k)\right] e^{(m+4) \omega / 2}
\end{aligned}
$$

Dividing by $e^{m \omega / 2}$, this may be written

$$
\begin{aligned}
& \frac{m}{2}\left(\mathrm{~J}\left[e^{2 \omega}(g+k)\right] e^{2 \omega}-\mathrm{J}[g+k]\right)= \\
& e^{-m \omega / 2} \underbrace{\Delta[g+k] e^{m \omega / 2}}_{\Delta[g] e^{m \omega / 2}} .
\end{aligned}
$$

The RHS in this last eqn. is polyn. in $m$, and in fact (since $\Delta=\delta d$ ) one with no constant term. After evaluation at some $x \in M$, these are true polynomials (with numerical coefficients). In particular, everything is indep. of $y \in T$. This establishes the LHS (eval. at $x$ ) as a polyn. in $m$, indep. of the $T$-parameter $y$.

What's happening here is: there was an $(m+1)^{-1}$ involved in defining J from K. But the formula for the difference of $J$ quantities at conformally related metrics is supplying an $m+1$ factor to keep things polynomial.

When going to higher order, it's important to be able to conclude the polynomial nature as above, rather than by explicit formulas.

We have either

- $\infty$ many eqns. (param. by $m$ ) on $M$; or
- a polynomial eqn. (in the vbl. m) at each $x \in M$.

Taking the 2nd viewpoint and harvesting the termwise eqns., the $m^{0}$ level gives

$$
\Delta[g] 1=0
$$

What the $m^{1}$ level provides can be seen by truncating power series:

$$
\begin{aligned}
& \left(1-\frac{m \omega}{2}\right) \Delta[g]\left(1+\frac{m \omega}{2}\right)= \\
& \frac{m}{2}\left(\mathrm{~J}\left[e^{2 \omega}(g+k)\right] e^{2 \omega}-\mathrm{J}[g+k]\right)+O\left(m^{2}\right),
\end{aligned}
$$

where $O\left(m^{2}\right)$ denotes a polynomial with a factor of $m^{2}$. (Recall that all this is happening at some $x \in M$.)

This is

$$
\begin{aligned}
& -\frac{m \omega}{2} \underbrace{\Delta[g] 1}_{0}+\frac{m}{2} \Delta[g] \omega= \\
& \frac{m}{2} \mathrm{~J}\left[e^{2 \omega} g\right] e^{2 \omega}-\frac{m}{2} \mathrm{~J}[g]+O\left(m^{2}\right),
\end{aligned}
$$

since to get the $m^{0}$ term in the polynomial $J\left[e^{2 \omega}(g+k)\right] e^{2 \omega}$, we may evaluate at $m=0$.

This last equation really doesn't have anything (even typographical) to remind us of the torus $T$. The $m^{1}$ coef. in the last eqn. is Gauss curvature prescription,

$$
\Delta[g] \omega+\mathrm{J}[g]=\mathrm{J}\left[e^{2 \omega} g\right] e^{2 \omega} \quad \text { (dim. 2) }
$$

Now generalize this to a way of getting the critical Q-curvature prescription

$$
P_{n}[g] \omega+Q_{n}[g]=Q_{n}\left[e^{2 \omega} g\right] e^{n \omega}
$$

(in the non-density version) from the subcritical (Yamabe type) Q-prescription equations in higher dimensions:

$$
\begin{aligned}
& \left(P_{n}^{0}[g]+\frac{N-n}{2} Q_{n}[g]\right) e^{(N-n) \omega / 2}= \\
& \frac{N-n}{2} Q_{n}\left[e^{2 \omega} g\right] e^{(N+n) \omega / 2} .
\end{aligned}
$$

Here $n$ is the target dim. and $N$ is the running dim.

This is the conformal covariance law

$$
\widehat{P}_{n}=\Omega^{-(N+n) \omega / 2} P_{n} \Omega^{(N-n) \omega / 2}
$$

for the GJMS operator

$$
P_{n}=P_{n}^{0}+\frac{N-n}{2} Q_{n}
$$

applied to the function 1 .

We want to use this in the situation where the manifold is a product of the $n$-dim. $M$ and a flat $m$-torus $T$ (so $N=n+m$ ):

$$
\begin{aligned}
& \left\{P_{n}^{0}[g+k]+\frac{m}{2} Q_{n}[g+k]\right\} e^{m \omega / 2}= \\
& \frac{m}{2} Q_{n}\left[e^{2 \omega}(g+k)\right] e^{(m+2 n) \omega / 2}
\end{aligned}
$$

We divide by $e^{m \omega / 2}$ to get

$$
\begin{aligned}
& \frac{m}{2}\left(Q_{n}\left[e^{2 \omega}(g+k)\right] e^{n \omega}-Q_{n}[g+k]\right)= \\
& e^{-m \omega / 2} P_{n}^{0}[g+k] e^{m \omega / 2}
\end{aligned}
$$

Since $P_{n}^{0}$ has the form $\delta(\bullet) d$, this establishes

$$
Q_{n}\left[e^{2 \omega}(g+k)\right] e^{n \omega}-Q_{n}[g+k]
$$

as a rational function in $m$ (at each $x \in M$, indep. of $y \in T$ ). (Recall that we can keep the poles under control; the rightmost one is at $n-2$.) This is just as good, for our purposes, as the polynomial behavior we had in the Yamabe $\rightarrow$ GCP case, since
rat'l. fens. agree at $\infty$ many points

they agree

their power series expansions
(at some regular point) agree termwise.
And the engine behind these equivalences is (not complex variables but) polynomial continuation.

At the $m^{0}$ level we have $P_{n}^{0} 1=0$. At the $m^{1}$ level, it's

$$
\frac{m}{2} P_{n}^{0} \omega=\frac{m}{2}\left(Q_{n}\left[e^{2 \omega} g\right] e^{n \omega}-Q_{n}[g]\right) .
$$

But this is the critical Q-curvature prescription equation

$$
P_{n}[g] \omega+Q_{n}[g]=Q_{n}\left[e^{2 \omega} g\right] e^{n \omega} \quad(\operatorname{dim} . n),
$$

since $P_{n}=P_{n}^{0}$ in dim. $n$.

