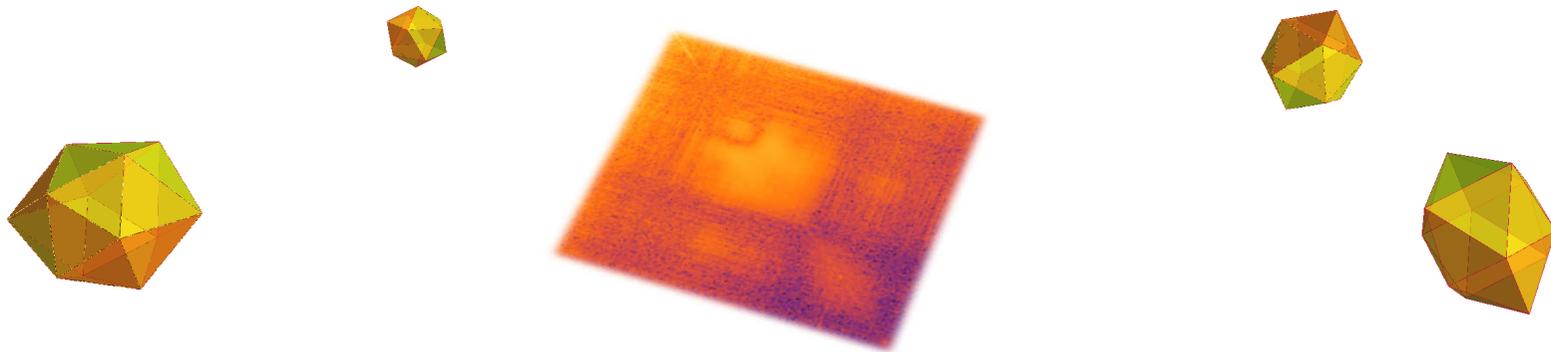


# A fluctuation x-ray scattering approach for biological imaging

Ruslan Kurta, European XFEL, Germany





- Introduction: x-ray scattering approaches for biological structure determination
  - Serial femtosecond crystallography (SFX)
  - Single particle coherent diffractive imaging (SPI)
  - Small angle x-ray scattering (SAXS)
  
- Fluctuation x-ray scattering (FXS)
  - Angular cross-correlation function (CCF): from SAXS to higher-order SAXS
  - Two- and three-point CCFs
  
- Applications of FXS
  - Single particle structure recovery from 2D disordered ensembles
  - 3D structure of viruses from single-particle imaging experiments
  
- Summary



**FERMI 2011**



**SACLA 2011**



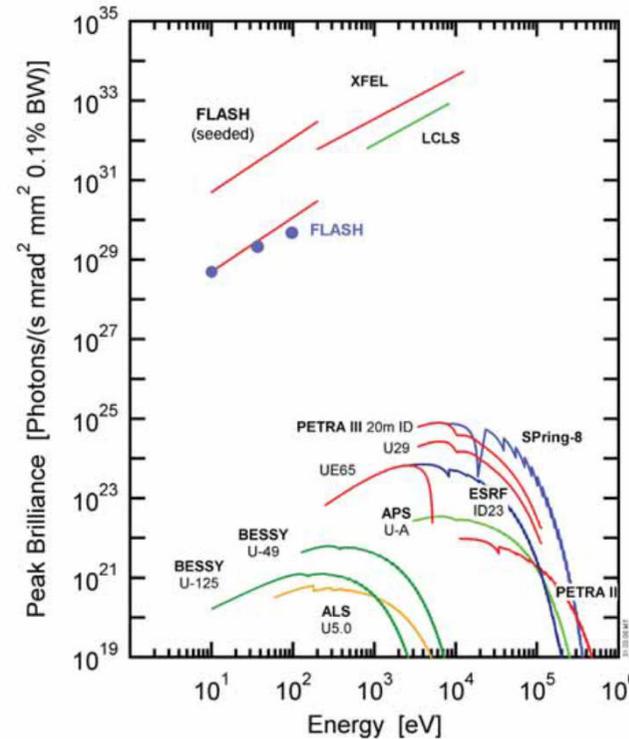
**PAL XFEL 2016**



**LCLS 2009**



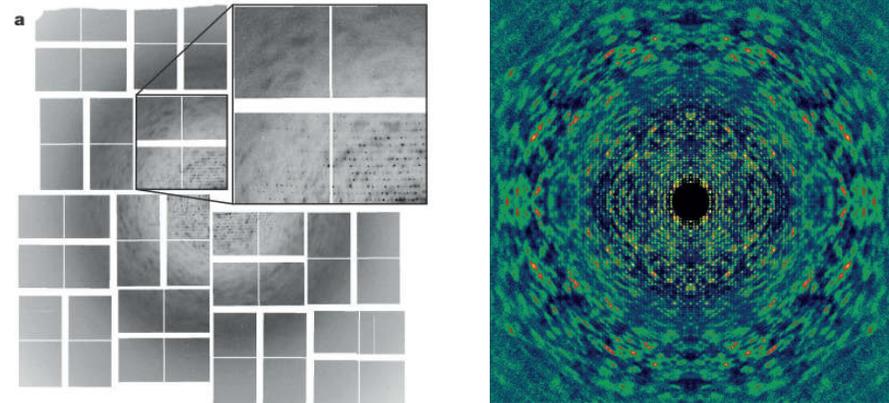
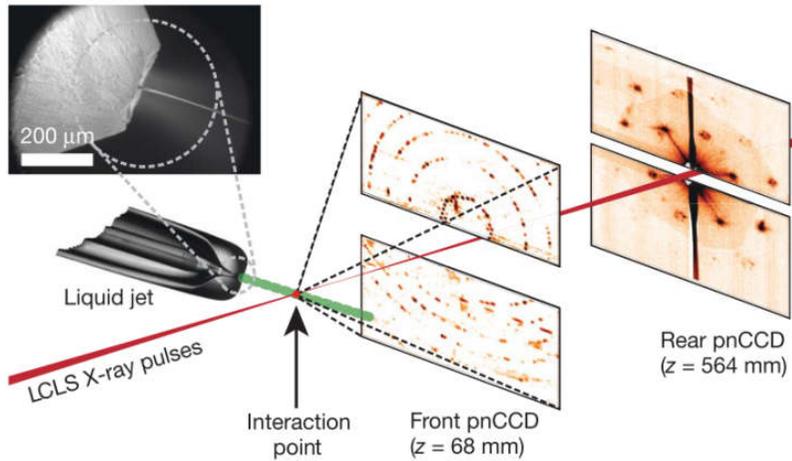
**FLASH 2005**



**European XFEL (2017)**

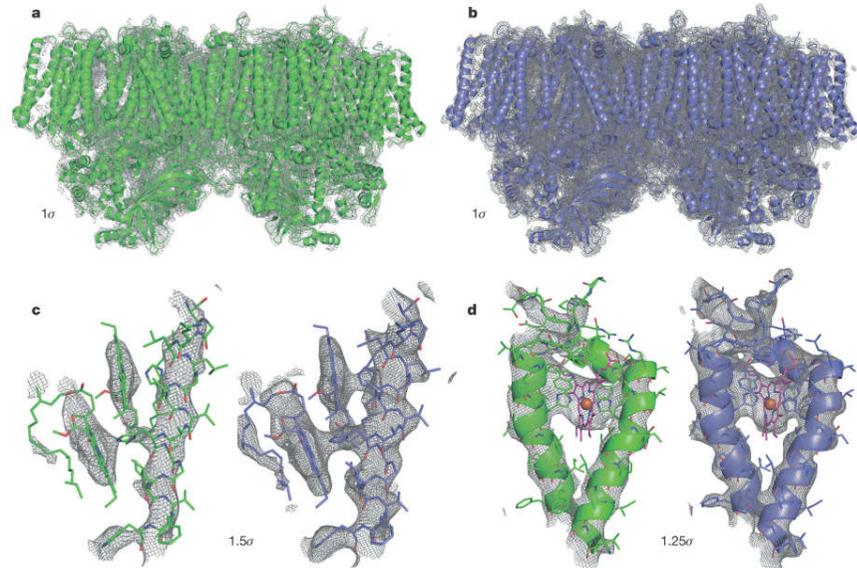


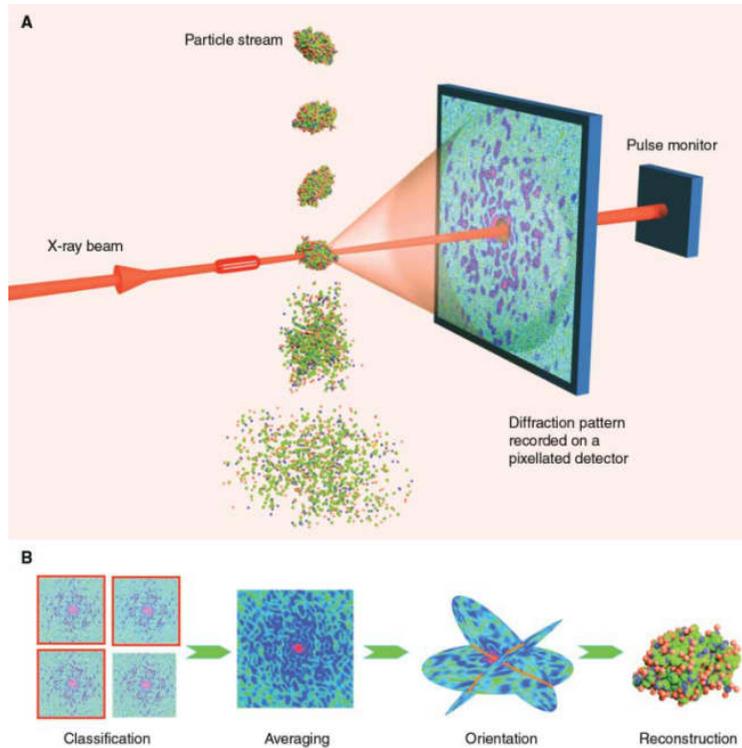
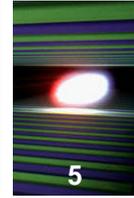
- ✓ Unprecedented peak brilliance
- ✓ Ultra-short pulse duration
- ✓ High repetition rate
- ✓ High coherence



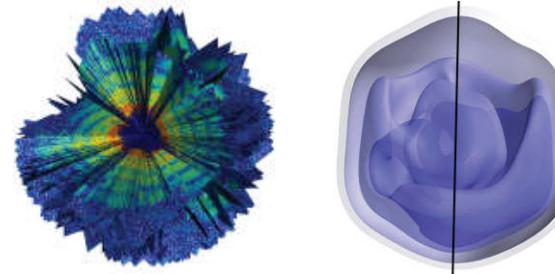
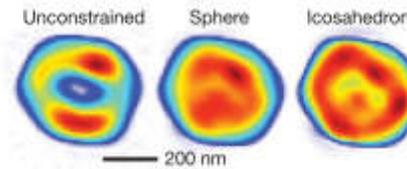
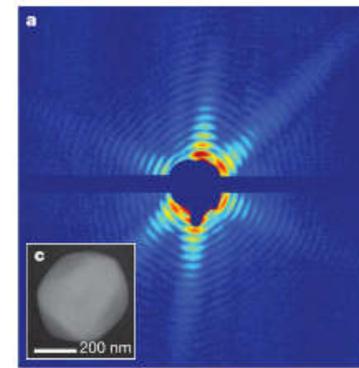
- SFX is an excellent tool for high-resolution structure studies of macromolecules and their complexes.
- Requires crystalline samples.

H. N. Chapman *et al.*, Nature 470, 73 (2011)  
K. Ayer *et al.*, Nature 530, 202 (2016)

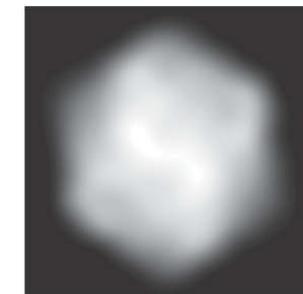
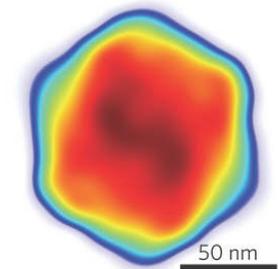
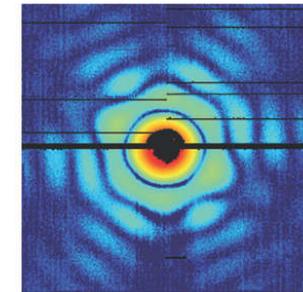




Mimivirus



Coarboxysome  
(cell organelle)



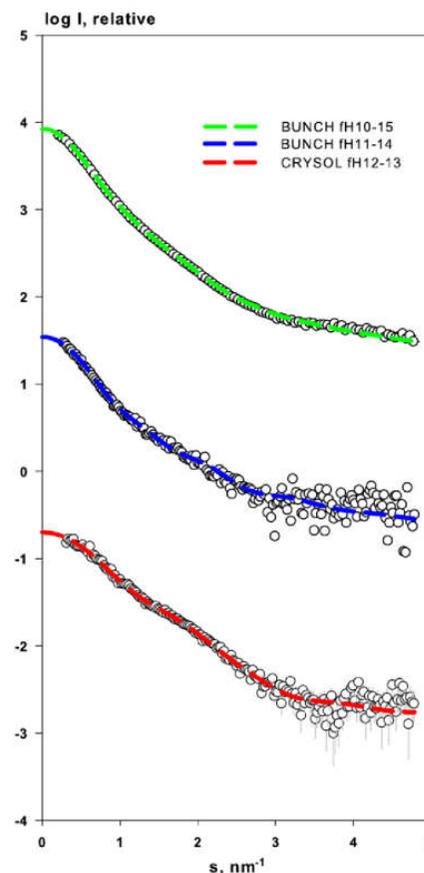
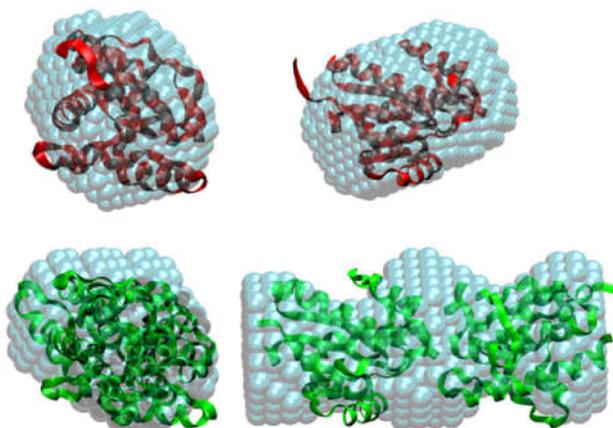
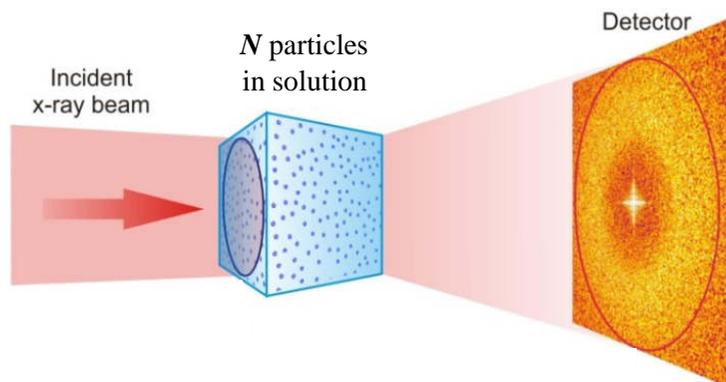
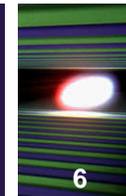
■ Requires extremely high photon flux (XFELs)

K. J. Gaffney, H. N. Chapman, *Science* 316, 5830 (2007)

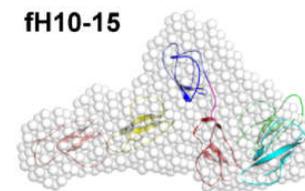
M. M. Seibert *et al.*, *Nature* 470, 78 (2011)

F. M. Hantke *et al.*, *Nat. Phot.* 8, 943 (2014)

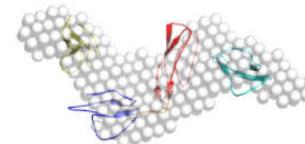
T. Eckeberg *et al.*, *PRL* 114, 098102 (2015)



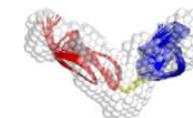
**fH10-15**



**fH11-14**



**fH12-13**

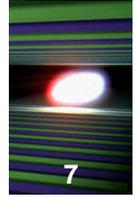


- Particles in native environment
- Low-resolution particle structure

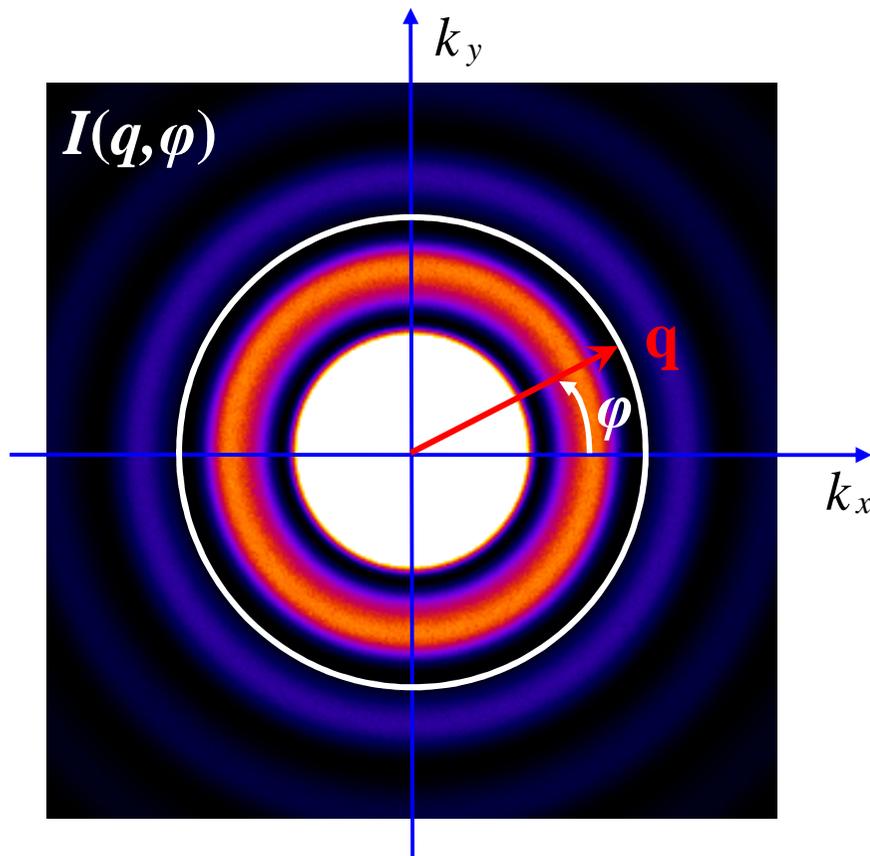
H. D. T. Merens and D. I. Svergun, *J. Struct. Biol.* 172, 128-141 (2010)

# Small-angle x-ray scattering (SAXS):

Dilute system of particles



7



Scattered intensity from  $N$  particles:

$$I(q, \varphi) = \sum_{i=1}^N I_i(q, \varphi) \quad (1)$$

where  $I_i(q, \varphi)$  is an individual contribution from the  $i$ -th particle.

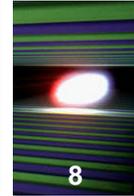
SAXS intensity:

$$I(q) \equiv \langle I(q, \varphi) \rangle_{\varphi} = \frac{1}{2\pi} \int_0^{2\pi} I(q, \varphi) d\varphi,$$

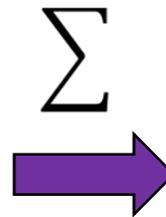
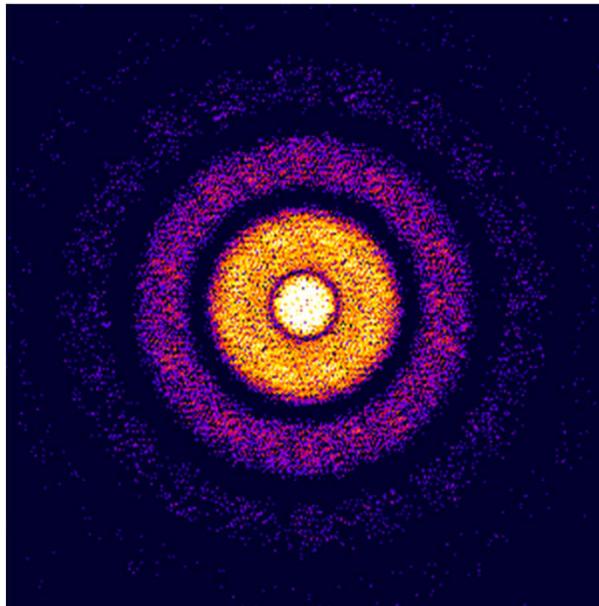
or

$$I(q) = \left\langle \sum_{i=1}^N I_i(q, \varphi) \right\rangle_{\varphi} \quad (2)$$

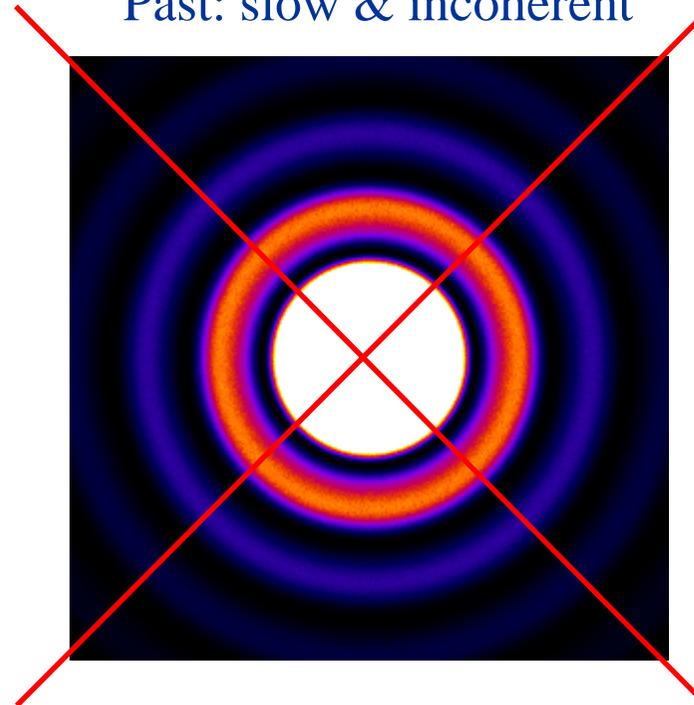
# Scattering of coherent ultrashort pulses from a system of $N$ particles

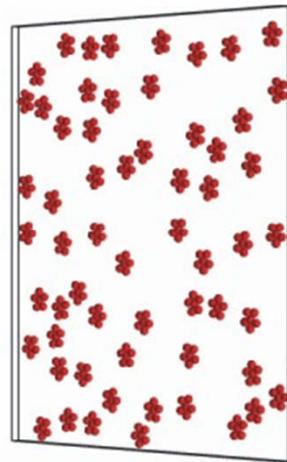
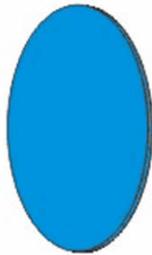
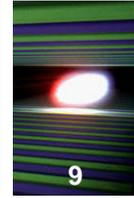


Present : fast & coherent



Past: slow & incoherent



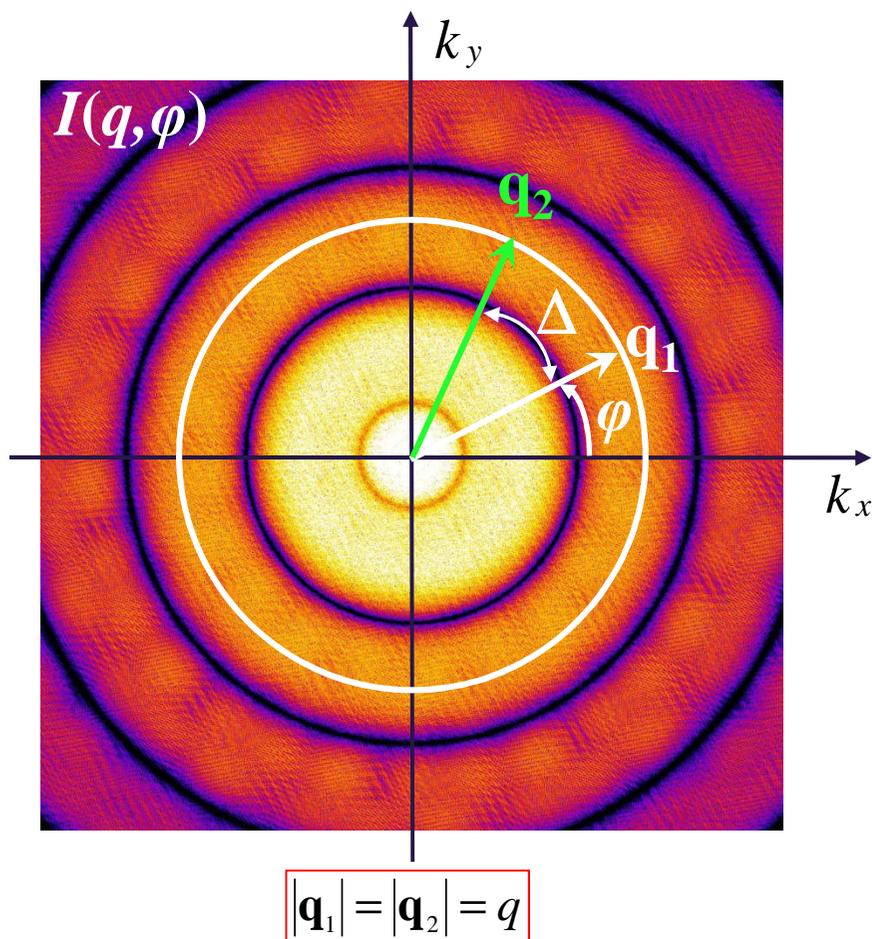
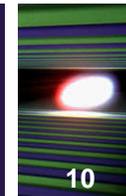


- Avoids intensity smearing due to rotational diffusion by applying ultrashort x-ray pulses (XFELs)
- Advantageous for weakly-scattering particles, since multiple-particle hits can be used along with the single hits

Z. Kam, *Macromolecules* 10, 927 (1977)

R.P. Kurta, M. Altarelli, I.A. Vartanyants, *Adv. Chem. Phys.* 161, Ch.1 (2016)

# X-ray cross-correlation analysis: two-point cross-correlation function



Two-point cross-correlation function (CCF):

$$C(q, \Delta) = \langle I(q, \varphi) I(q, \varphi + \Delta) \rangle_{\varphi} \quad (1)$$

Fourier components:

$$C^n(q) = \frac{1}{2\pi} \int_0^{2\pi} C(q, \Delta) \exp(-in\Delta) d\Delta \quad (2)$$

$$I^n(q) = \frac{1}{2\pi} \int_0^{2\pi} I(q, \varphi) \exp(-in\varphi) d\varphi \quad (3)$$



$$\langle \langle I(q, \varphi) \rangle_{\varphi} \rangle_M = \langle I^0(q) \rangle_M - \text{conventional SAXS} \quad (4)$$

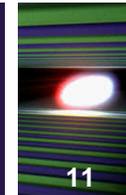
$$\langle C^n(q) \rangle_M = \langle |I^n(q)|^2 \rangle_M - \text{“higher-order SAXS”} \quad (5)$$

$\langle \rangle_M$  - averaging over diffraction patterns.

Z. Kam, *Macromolecules* 10, 927 (1977)

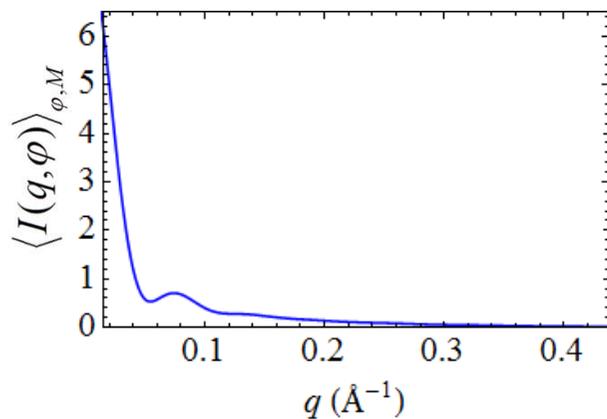
R.P. Kurta, M. Altarelli, I.A. Vartanyants, *Adv. Chem. Phys.* 161, Ch.1 (2016)

# Single-particle structure recovery from solution x-ray scattering

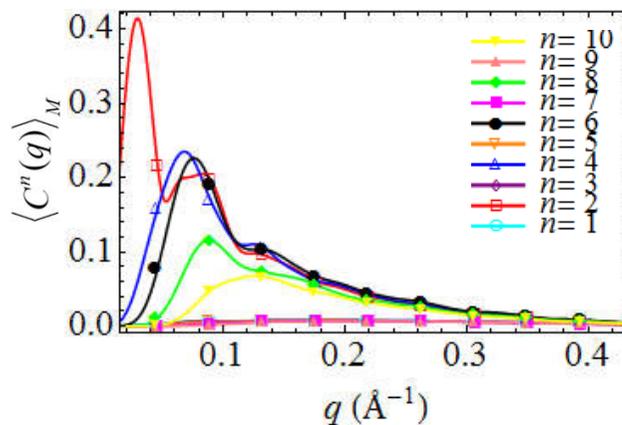


$$\langle I(q, \varphi) \rangle_{\varphi, M} = \langle I^0(q) \rangle_M$$

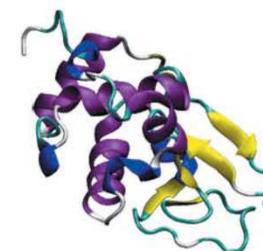
$$\langle C^n(q) \rangle_M \Rightarrow \langle |I^n(q)|^2 \rangle_M, n = 1, 2, 3, \dots$$



+



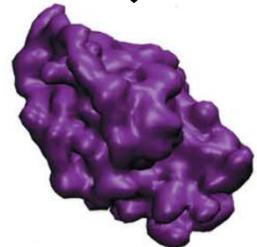
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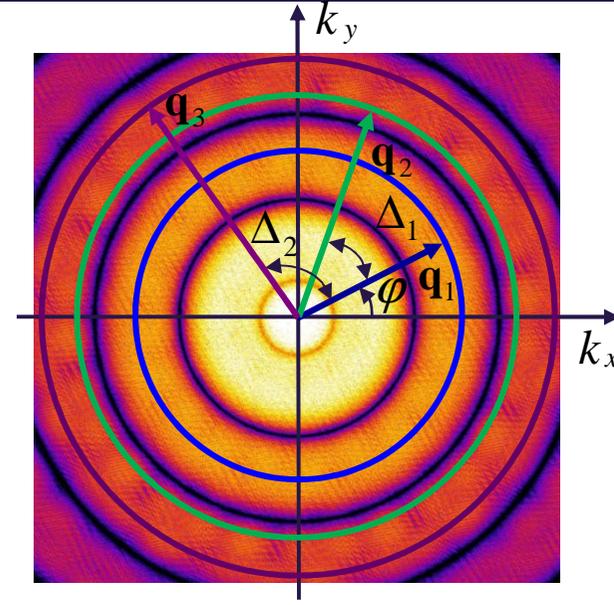
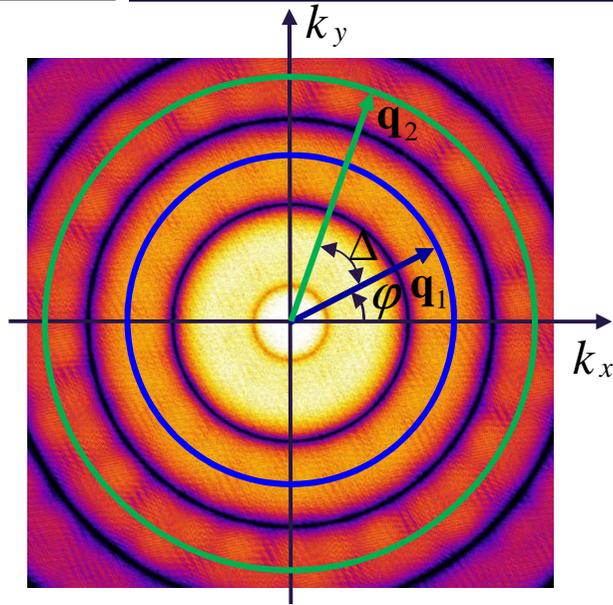
Bead model, etc



Conventional SAXS analysis



Higher-order SAXS analysis



Two-point CCF:

$$\langle C(q_1, q_2, \Delta) \rangle = \left\langle \left\langle I(q_1, \varphi) I(q_2, \varphi + \Delta) \right\rangle_{\varphi} \right\rangle_M$$

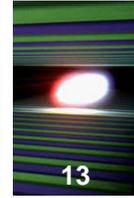
Three-point CCF:

$$\langle C(q_1, q_2, q_3, \Delta_1, \Delta_2) \rangle = \left\langle \left\langle I(q_1, \varphi) I(q_2, \varphi + \Delta_1) I(q_3, \varphi + \Delta_2) \right\rangle_{\varphi} \right\rangle_M$$

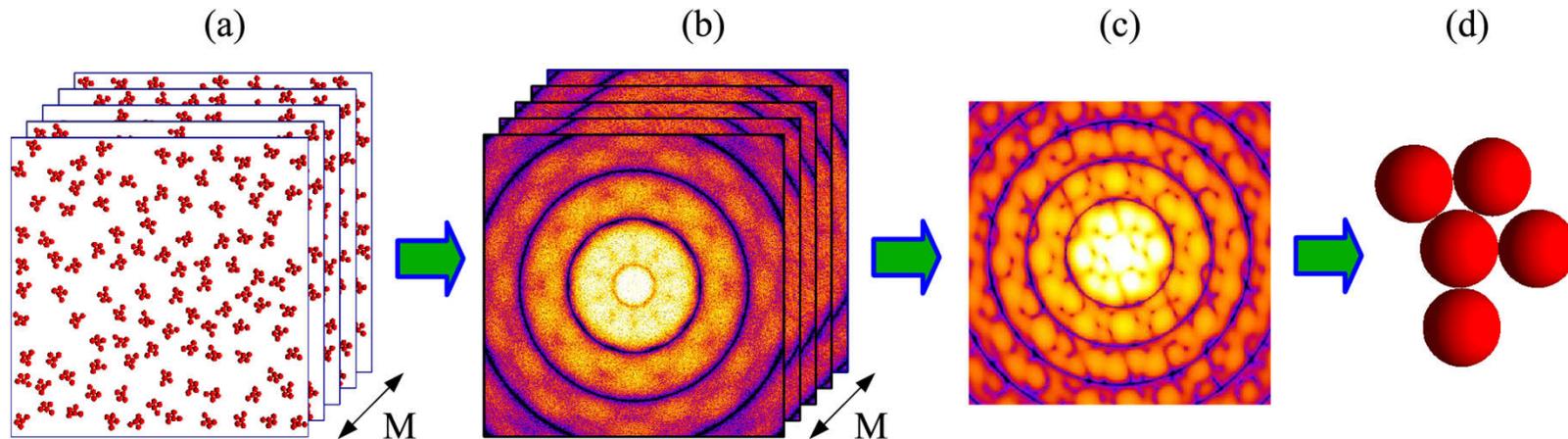
where  $\langle f(q, \varphi) \rangle_{\varphi} = \frac{1}{2\pi} \int_0^{2\pi} f(q, \varphi) d\varphi$  - angular average,

$\langle \rangle_M$  - statistical average (over diffraction patterns).

# Single particle structure recovery from a disordered system of 2D particles



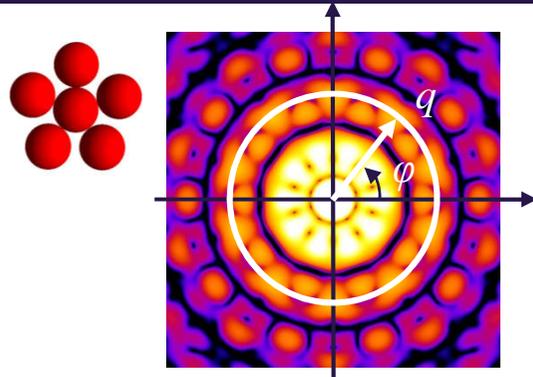
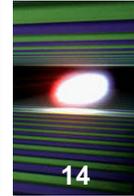
Experiment → XCCA → Phase retrieval → Structure



- (a) A large number  $M$  of realizations of a disordered system composed of  $N$  reproducible particles;
- (b) Measured set of diffraction patterns;
- (c) Diffraction pattern corresponding to a single particle;
- (d) Structure of the single particle.

R.P. Kurta, R. Dronyak, M. Altarelli, E. Weckert and I.A. Vartaniants, *New J. Phys.* 15, 013059 (2013)  
 R.P. Kurta, M. Altarelli, I.A. Vartaniants, *Adv. Cond. Matt. Phys.* 2013, 959835 (2013)

# Recovery of the scattered intensity from a single particle by means of XCCA



Fourier series of the scattered intensity from a single particle:

$$I(q, \varphi) = \sum_{n=-\infty}^{\infty} |I_q^n| \exp(i\phi_q^n) \exp(in\varphi)$$

Fourier components of the CCF:

$$\langle C_{q_1, q_2}^n \rangle_M = \frac{1}{2\pi} \int_0^{2\pi} \langle C(q_1, q_2, \Delta) \rangle_M \exp(-in\Delta) d\Delta$$

2-point:

$$\langle C_{q_1, q_2}^n \rangle_M = I_{q_1}^{n*} \cdot I_{q_2}^n \cdot N$$

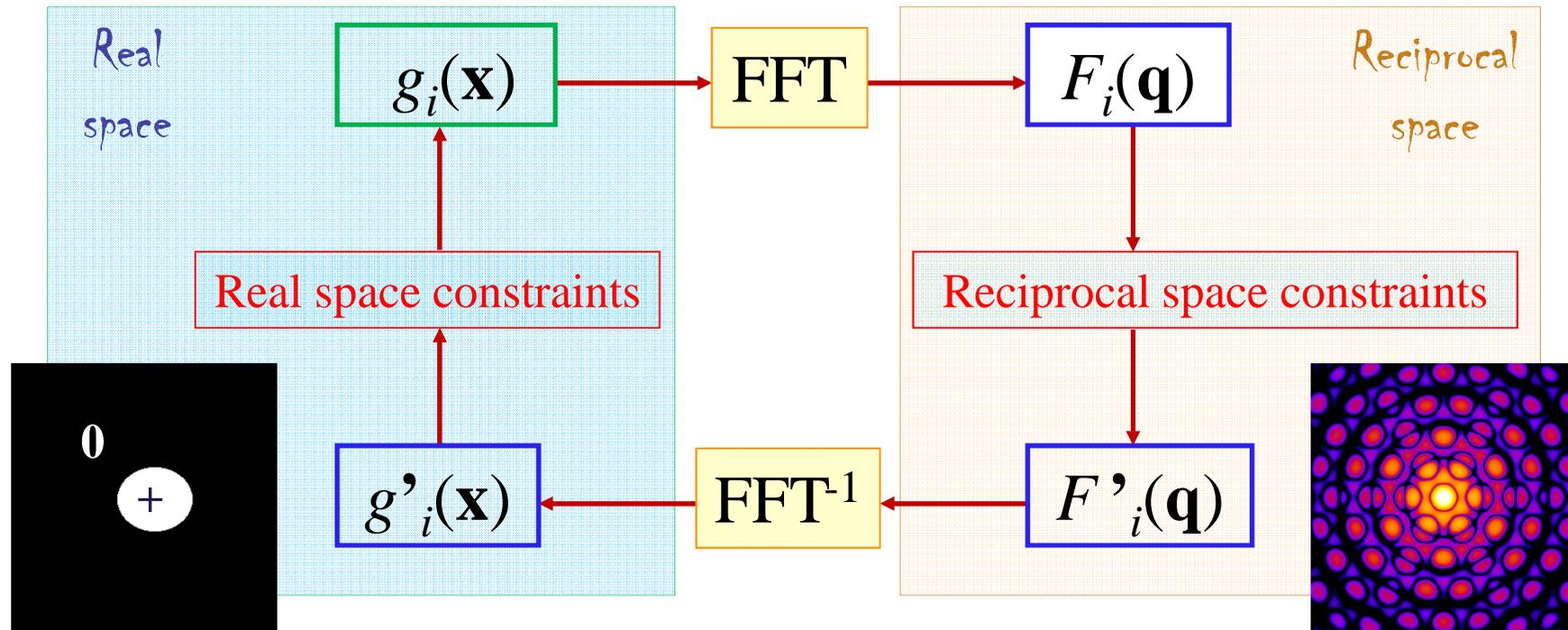
Many-particle quantities

single-particle quantities

3-point:

$$\langle C_{q_1, q_2, q_3}^{n_1, n_2} \rangle_M = I_{q_1}^{(n_1+n_2)*} I_{q_2}^{n_1} I_{q_3}^{n_2} \cdot N$$

$$\begin{cases} |I_q^n| \\ \arg(\langle C_{q_1, q_2}^n \rangle_M) = \phi_{q_2}^n - \phi_{q_1}^n \\ \arg(\langle C_{q_1, q_2, q_3}^{n_1, n_2} \rangle_M) = \phi_{q_2}^{n_1} + \phi_{q_3}^{n_2} - \phi_{q_1}^{n_1+n_2} \end{cases}$$



## Real space constraints:

- Finite support
- Positivity

## Reciprocal space constraints:

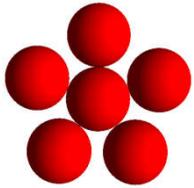
- Intensity constraint (taking into account missing data due to a beamstop, detector gaps, etc)

$$|F_i(\mathbf{q})| \rightarrow \sqrt{I_{\text{exp}}(\mathbf{q})}$$

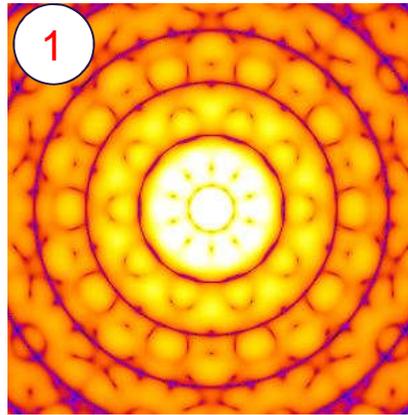
R.W.Gerchberg, W.O. Saxton, Optik 35, 237 (1972)

J.R. Fienup, Appl. Opt. 21, 2758 (1982)

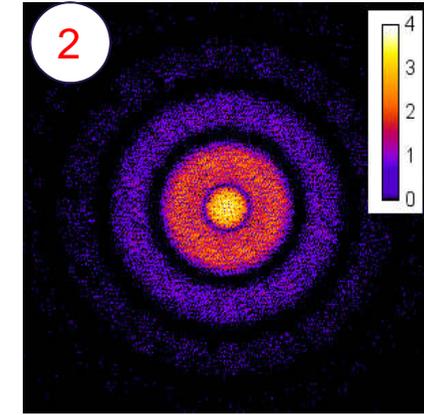
V. Elser, J. Opt. Soc. Am. A 20, 40 (2003)



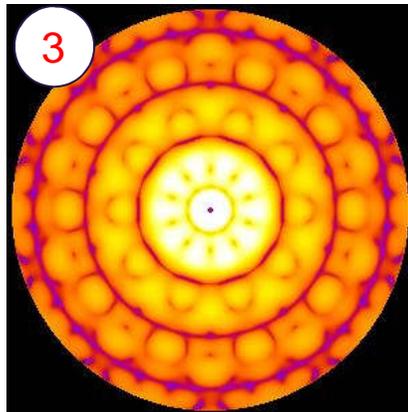
Single pentagonal  
cluster ( $N=1$ )



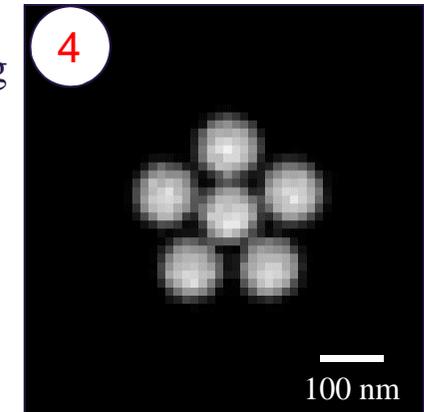
Coherent scattering,  
 $N=10$  clusters,  
Fluence:  $4 \cdot 10^{10}$  ph/ $\mu\text{m}^2$ ,  
Poisson noise,  
 $M=10^5$  diffraction patterns.



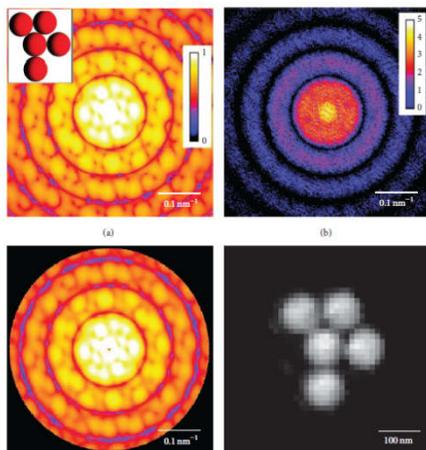
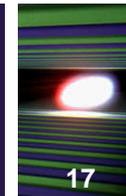
Diffraction pattern of  
a single cluster,  
recovered from  
 $M=10^5$  patterns of the  
form (2)



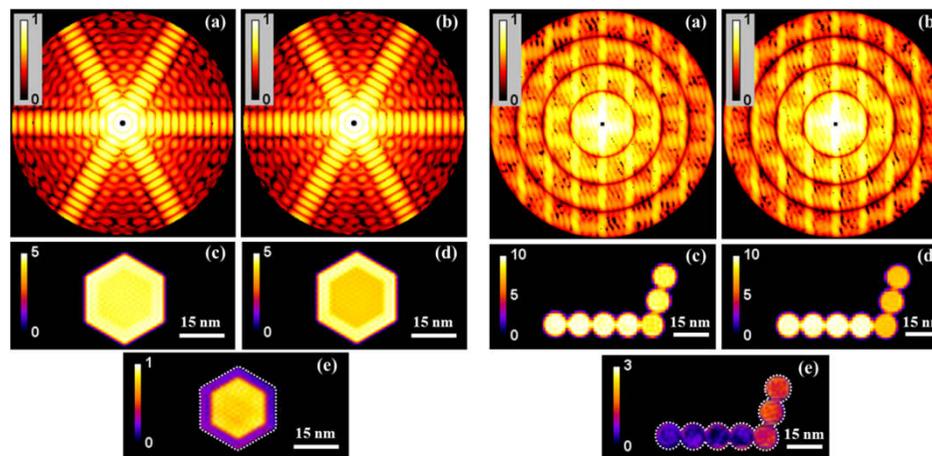
Projected electron density of  
the cluster, reconstructed using  
iterative phase retrieval  
algorithms.



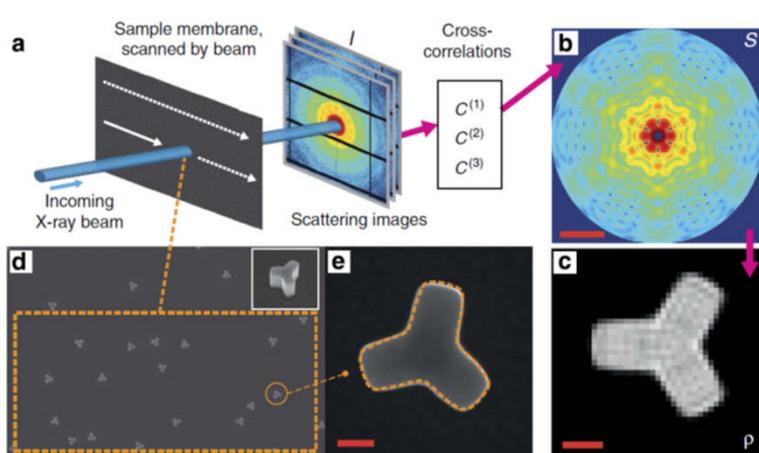
# Single-particle structure from disordered 2D ensembles



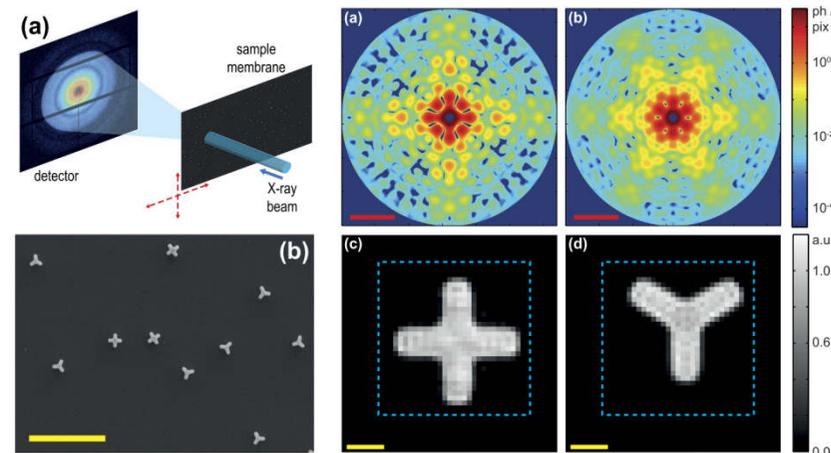
R.P. Kurta et al., *New J. Phys.* 15, 013059 (2013)  
R.P. Kurta et al., *Adv. Chem. Phys.* 161, Ch. 1 (2016)



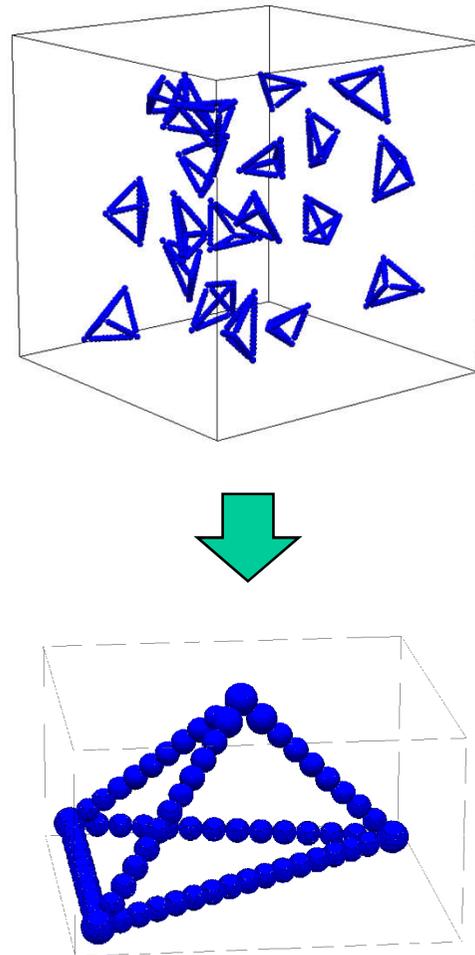
R.P. Kurta, *J. Phys. B: At. Mol. Opt. Phys.* 49, 165001 (2016)



B. Pedrini et al., *Nat. Comm.* 4, 1647 (2013)



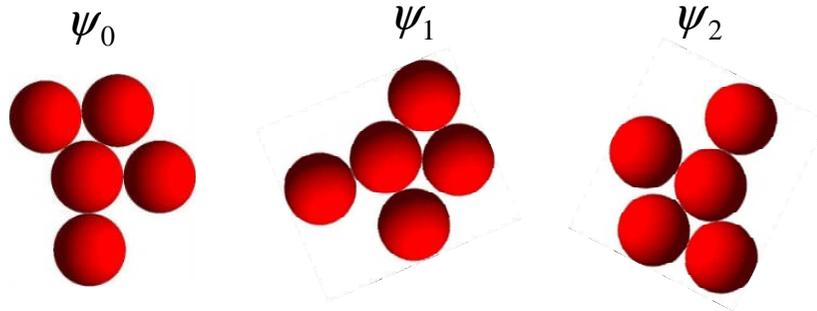
B. Pedrini et al., *Sci. Rep.* 7, 45618 (2017)



Z. Kam, *Macromolecules* 10, 927 (1977)



## 2D systems

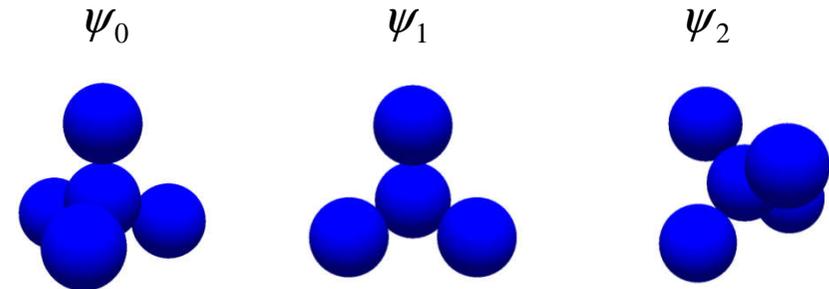


$$\langle C_{q_1, q_2}^n \rangle_M = I_{q_1, \psi_0}^{n*} \cdot I_{q_2, \psi_0}^n$$

$$\langle C_{q_1, q_2, q_3}^{n_1, n_2} \rangle_M = I_{q_1, \psi_0}^{(n_1+n_2)*} I_{q_2, \psi_0}^{n_1} I_{q_3, \psi_0}^{n_2}$$

$\langle \rangle_M$  - averaging over diffraction patterns.

## 3D systems



$$\langle C_{q_1, q_2}^n \rangle_M = \langle I_{q_1, \psi_i}^{n*} \cdot I_{q_2, \psi_i}^n \rangle_\psi$$

$$\langle C_{q_1, q_2, q_3}^{n_1, n_2} \rangle_M = \langle I_{q_1, \psi_i}^{(n_1+n_2)*} I_{q_2, \psi_i}^{n_1} I_{q_3, \psi_i}^{n_2} \rangle_\psi$$

$\langle \rangle_\psi$  - averaging over different projections.



## ✓ Spherical harmonics

Z. Kam, *Macromolecules* 10, 927 (1977)  
 Z. Kam, *J. Theor. Biol.* 82, 1 (1980)  
 D. K. Saldin *et al.*, *J. Phys.: Cond. Matt.* 21, 134014 (2009)  
 D. Starodub *et al.*, *Nat. Comm.* 3, 1276 (2012)

## ✓ Icosahedral harmonics

D. K. Saldin *et al.*, *Opt. Express* 19, 17318 (2011)

## ✓ Cylindrical harmonics

D. K. Saldin *et al.*, *Acta Cryst. A* 66, 32 (2010)  
 G. Chen *et al.*, *J. Synch. Rad.* 19, 695 (2012)

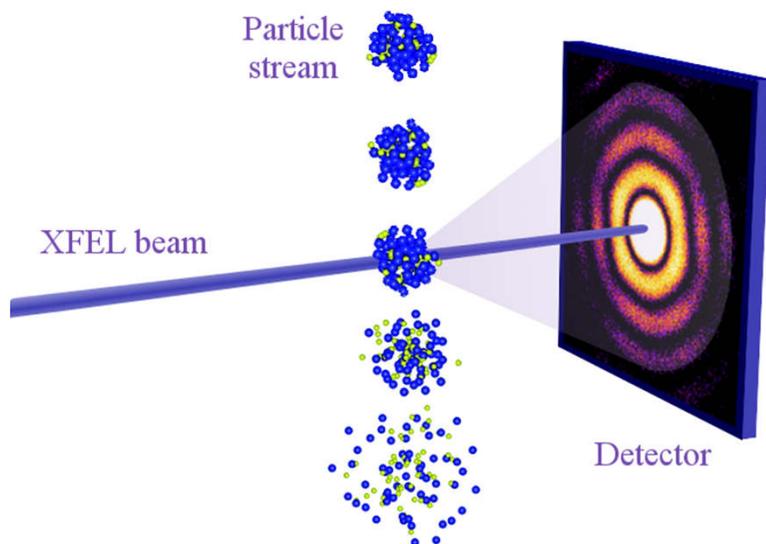
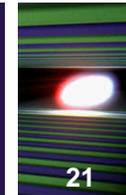
## ✓ Zernike polynomials

H. Liu *et al.*, *Acta Cryst. A* 68, 561 (2012)  
 H. Liu *et al.*, *Acta Cryst. A* 69, 365 (2013)

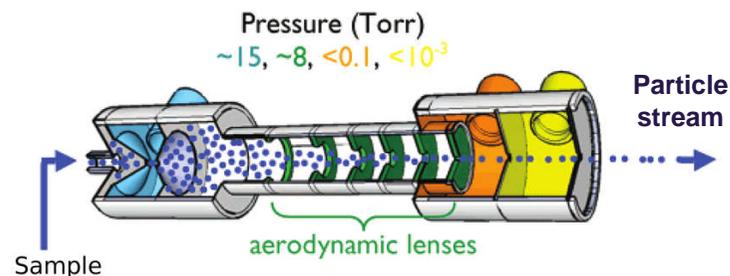
- Indirect methods: fitting or model assumptions are applied
- Better than conventional SAXS: additional information is available from the analysis



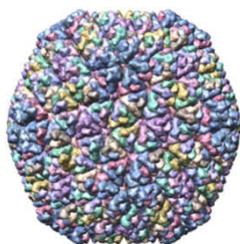
# SPI experiments on aerosolized virus particles at the AMO instrument, LCLS



Photon energy:  $E=1.6\text{keV}$   
 Sample-detector distance: 581 mm  
 Detector: pnCCD  
 Sample injection: aerodynamic lens stack system with a GDNV



## RDV



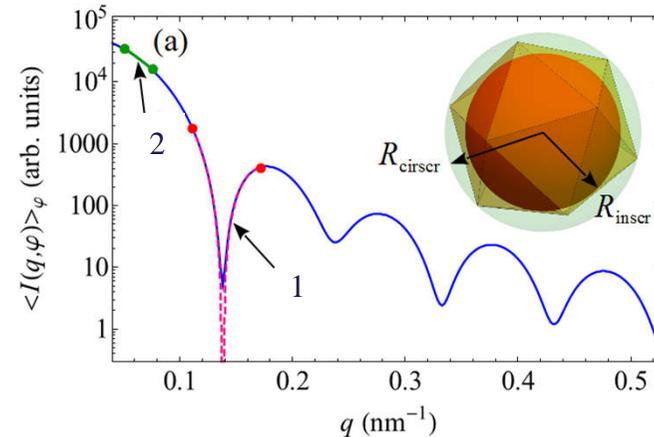
The Rice Dwarf Virus was the first studied plant pathogenic virus; an icosahedral double shelled virus, ranging from 70-75 nm in diameter.

## PR772



The Parenteral Drug Association virus filter task force has chosen PR772 as a model bacteriophage to standardize nomenclature for large-pore-size virus-retentive filters (filters designed to retain viruses larger than 50-60 nm in size).

H. K. N. Reddy *et al.*, Scientific Data 4, 170079 (2017)  
 M. Bogan *et al.*, Nano Lett. 8, 310 (2008)



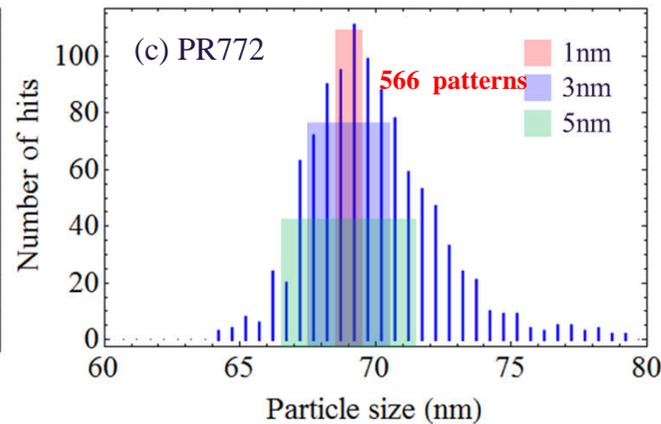
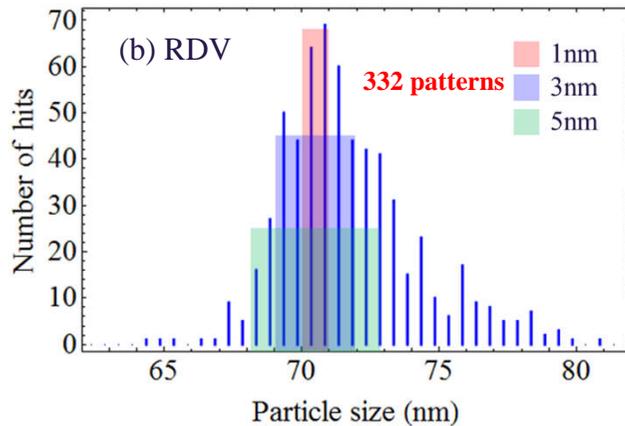
## 1) Spherical particle form factor fit

$$\langle I(q) \rangle_{\varphi} = A \left[ \frac{\sin(qR_s) - qR_s \cos(qR_s)}{q^3} \right]^2 \quad (1)$$

## 2) Guinier-type fit: ( $q < 1.3/R_g$ )

$$\langle I(q) \rangle_{\varphi} = I(0) \exp[-q^2 R_g^2 / 3] \quad (2)$$

$$y = ax + b \quad R_g = \sqrt{-3a}$$



## Exact definition of particle size:

$$\text{Size} = 2R_{circum} \quad (3)$$

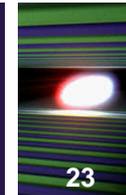
$$R_{circum} = \sqrt{15 - 6\sqrt{5}} R_{inscr} \quad (4)$$

## Empirical approximation:

$$R_{inscr} \approx (R_s + R_g) / 2 \quad (5)$$

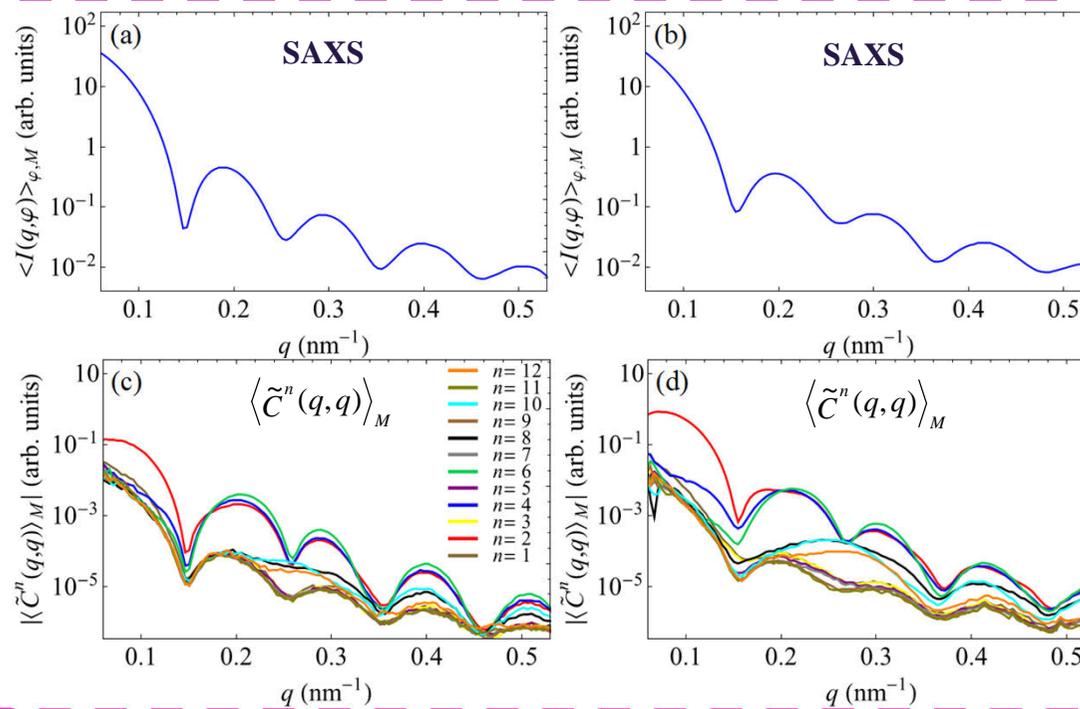
$$\text{Size} \approx 1.26(R_g + R_s) \quad (6)$$

$$(\text{or } \text{Size} \approx 2.36 R_s) \quad (7)$$

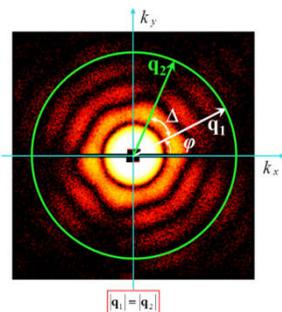


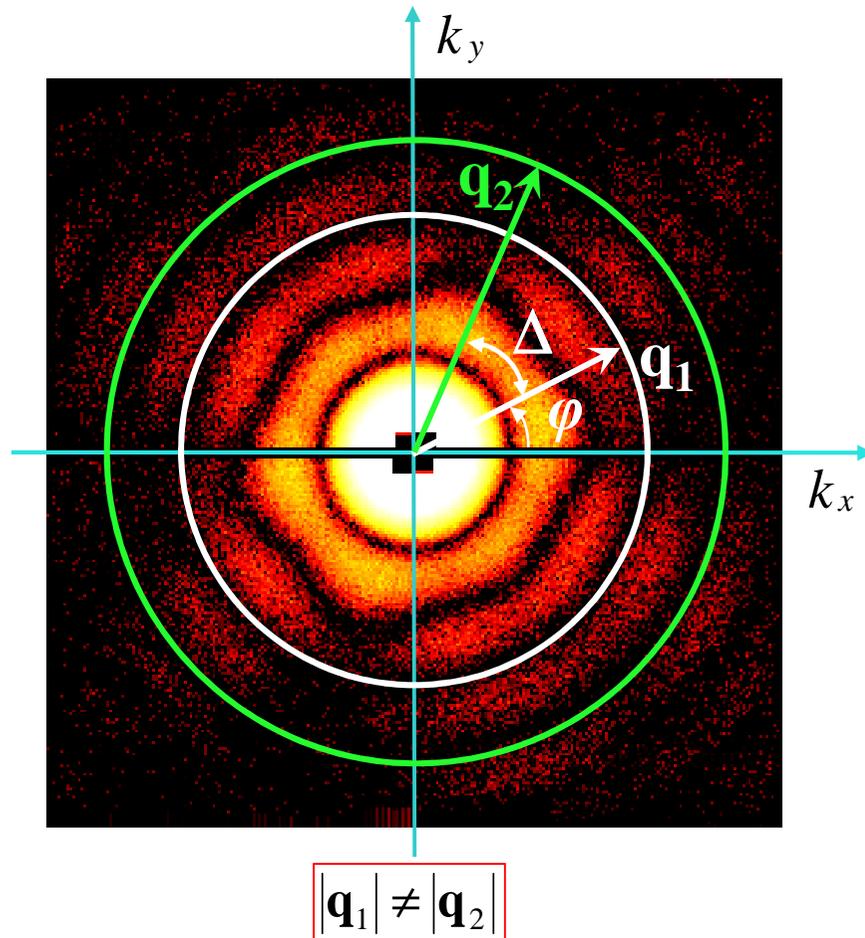
**RDV**

**PR772**



Generalization of  
SAXS:  
Higher-order SAXS





CCF defined on the same pattern  $i$ :

$$\langle C_{ii}(q_1, q_2, \Delta) \rangle_M = \left\langle \left\langle I_i(q_1, \varphi) I_i(q_2, \varphi + \Delta) \right\rangle_{\varphi} \right\rangle_M \quad (1)$$

CCF defined between different patterns  $i$  and  $j$ :

$$\langle C_{ij}(q_1, q_2, \Delta) \rangle_M = \left\langle \left\langle I_i(q_1, \varphi) I_j(q_2, \varphi + \Delta) \right\rangle_{\varphi} \right\rangle_M \quad (2)$$

Fourier components of the CCFs:

$$\langle C_{ij}^n(q_1, q_2) \rangle_M = \frac{1}{2\pi} \int_0^{2\pi} \langle C_{ij}(q_1, q_2, \Delta) \rangle_M \exp(-in\Delta) d\Delta \quad (3)$$

Difference spectrum:

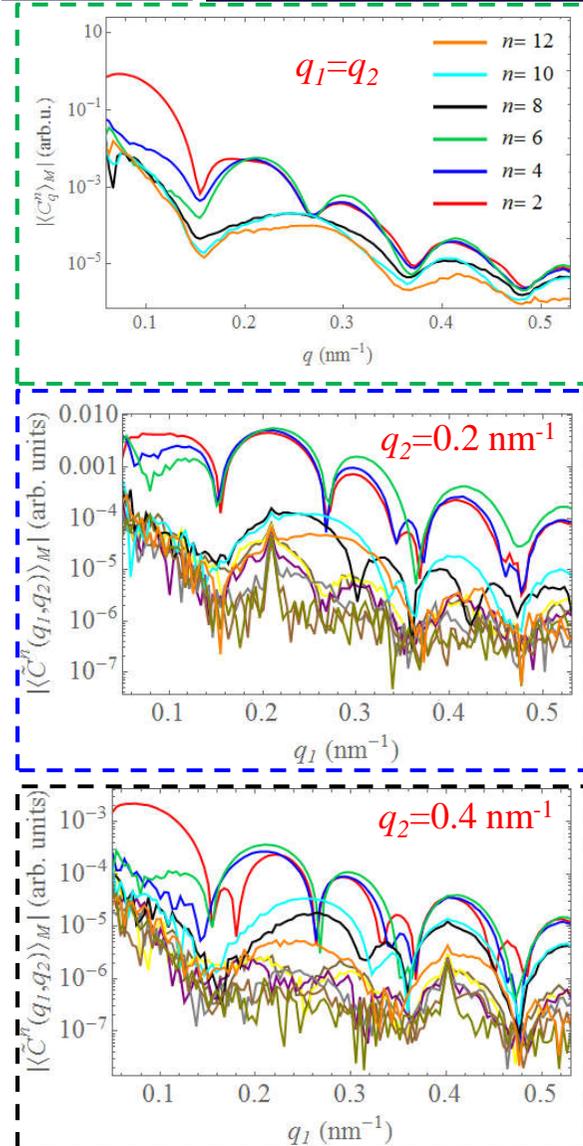
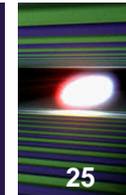
$$\langle \tilde{C}^n(q_1, q_2) \rangle_M = \langle C_{ii}^n(q_1, q_2) \rangle_M - \langle C_{ij}^n(q_1, q_2) \rangle_M \quad (4)$$

Difference spectrum is used to reduce undesirable effect of various systematic issues and improve the FXS data quality.

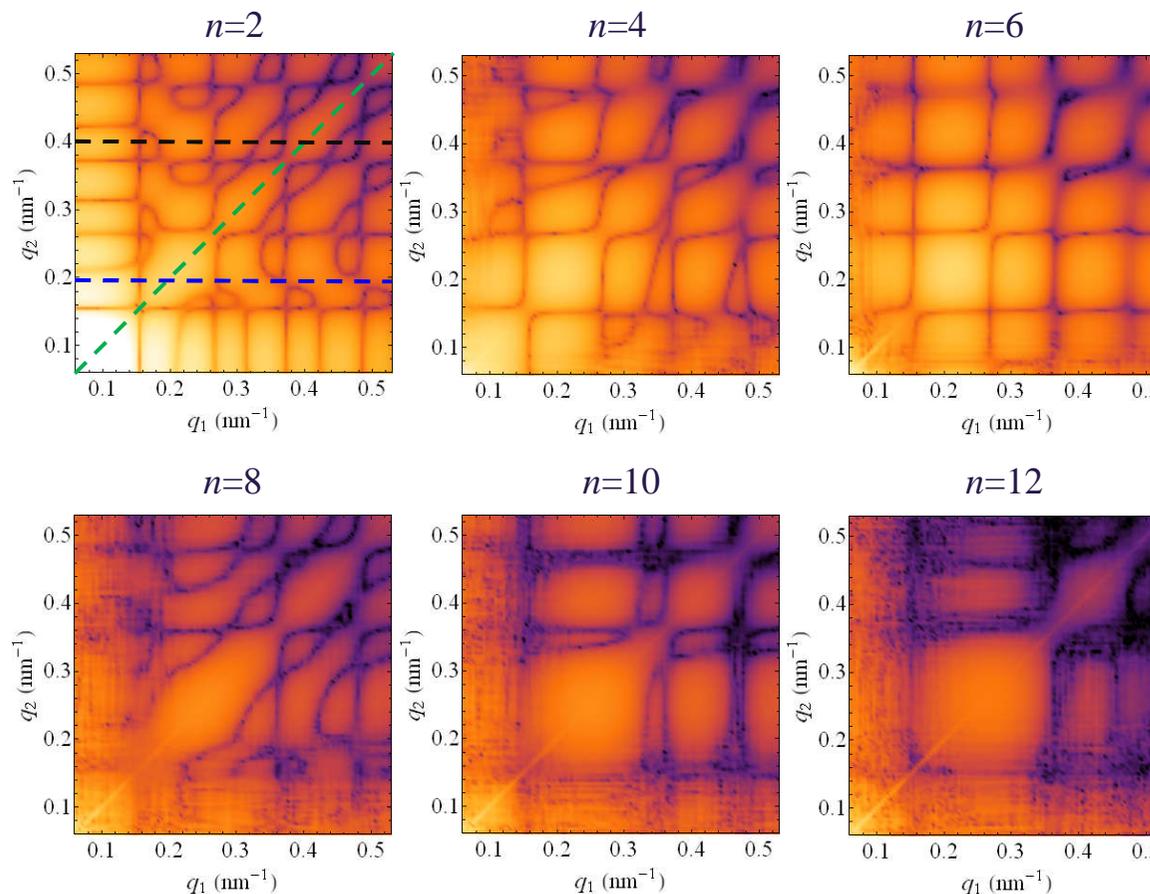
R.P. Kurta *et al.*, J. Phys.: Conf. Series 499, 012021 (2014)

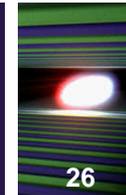
G. Chen *et al.*, J. Synch. Rad. 19, 695700 (2012)

R.P. Kurta, M. Altarelli, I.A. Vartanyants, Adv. Chem. Phys. 161, Ch.1 (2016)



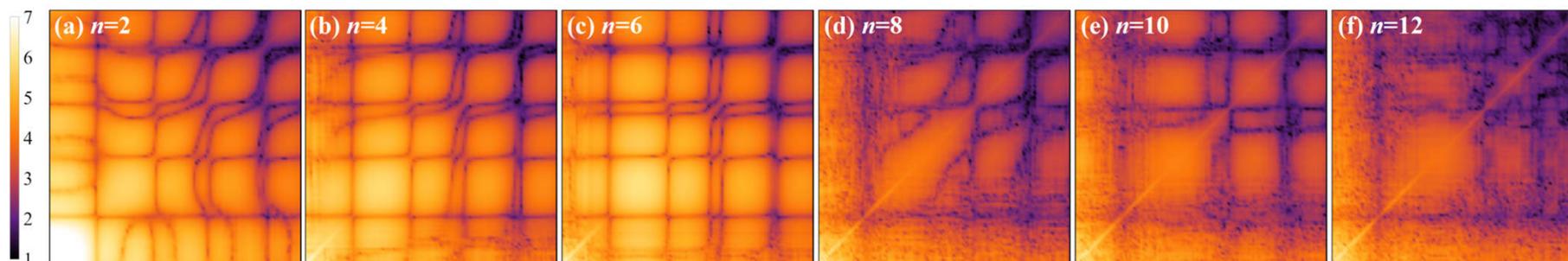
$$\left| \langle \tilde{C}^n(q_1, q_2) \rangle_M \right| = \left| \langle C_{ii}^n(q_1, q_2) \rangle_M - \langle C_{ij}^n(q_1, q_2) \rangle_M \right|$$



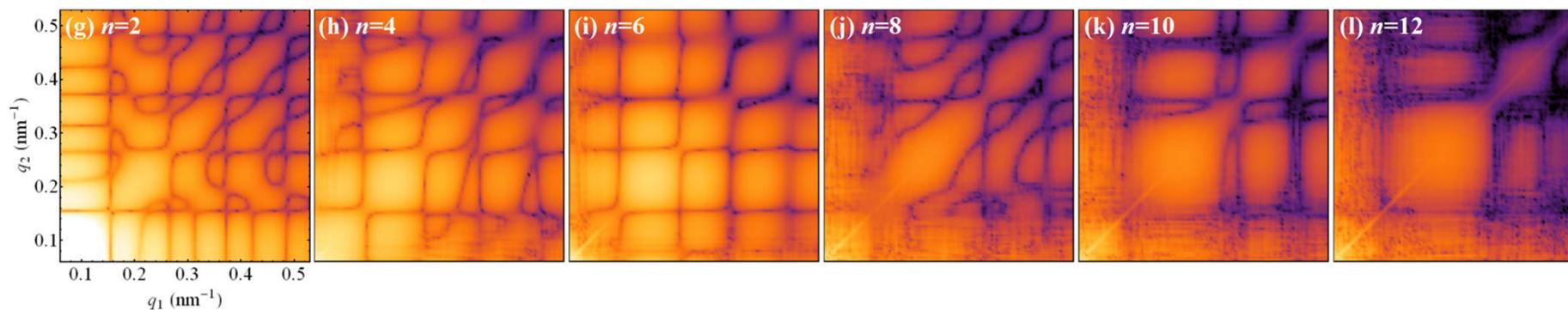


$$\text{Amplitudes: } \left| \left\langle \tilde{C}^n(q_1, q_2) \right\rangle_M \right|$$

## RDV

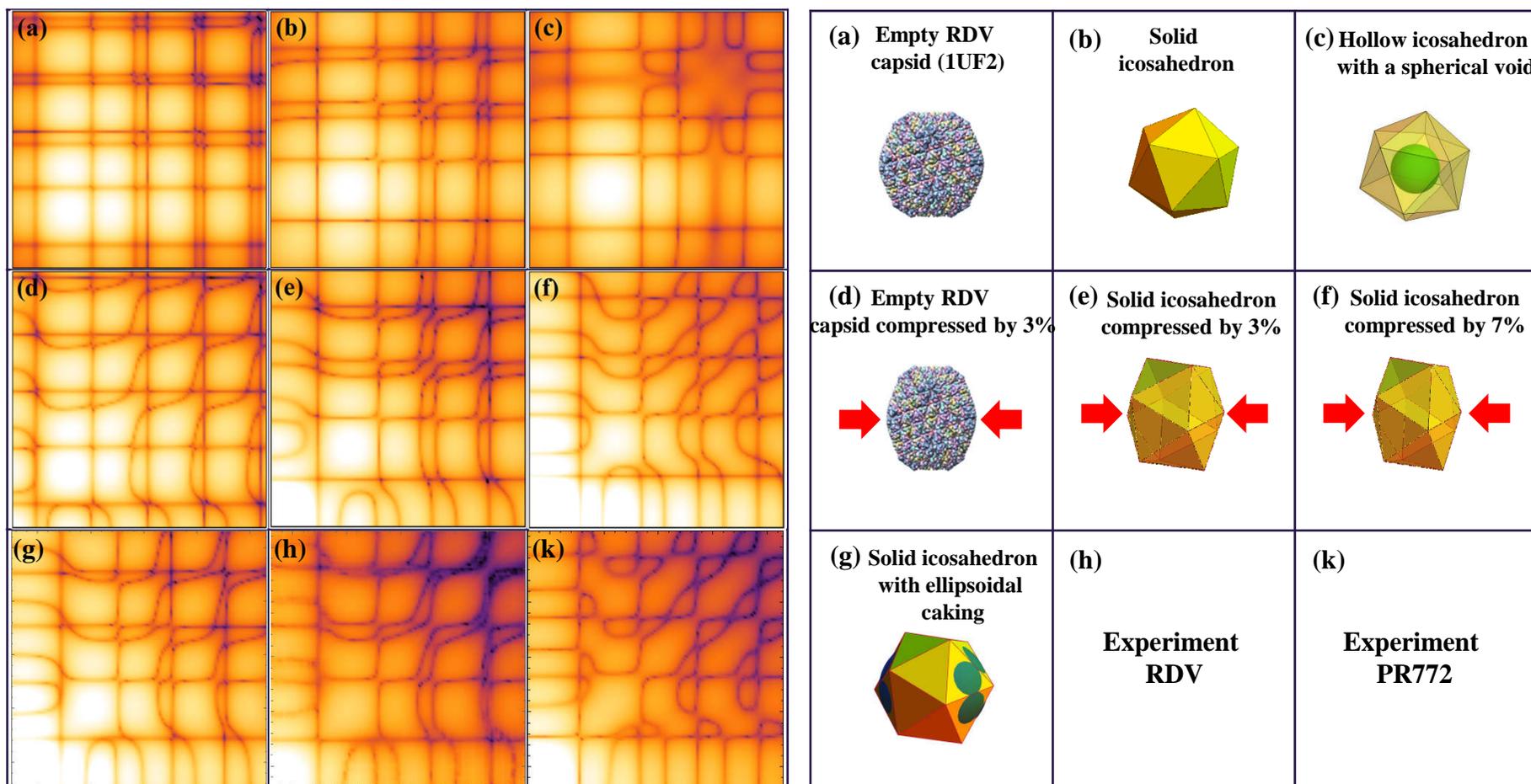


## PR772



$$\left| \left\langle \tilde{C}^{n=2}(q_1, q_2) \right\rangle_M \right|$$

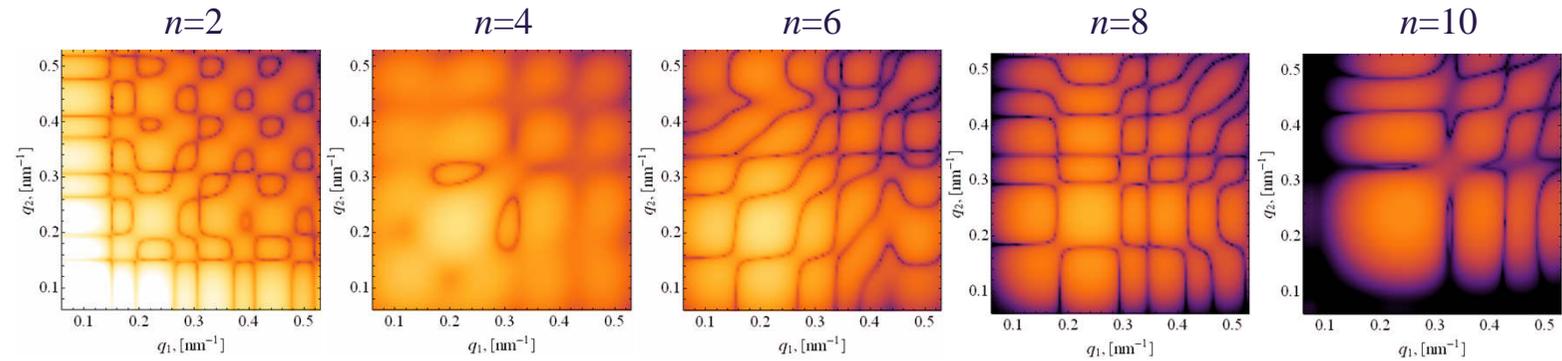
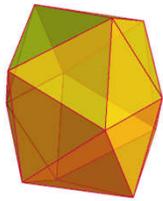
Model structures



R.P. Kurta *et al.*, PRL 119, 158102 (2017)

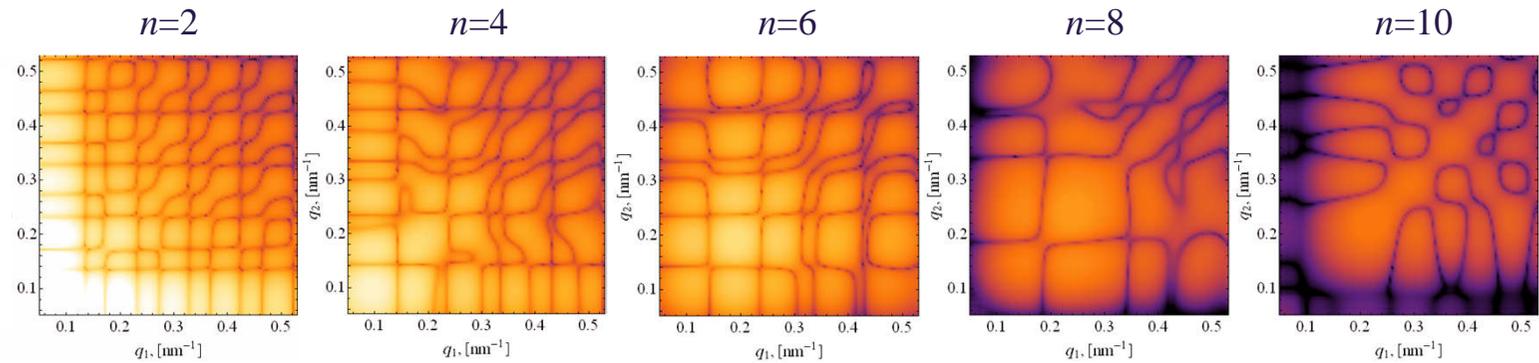
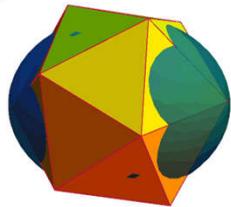
## Particle distortion

0.750

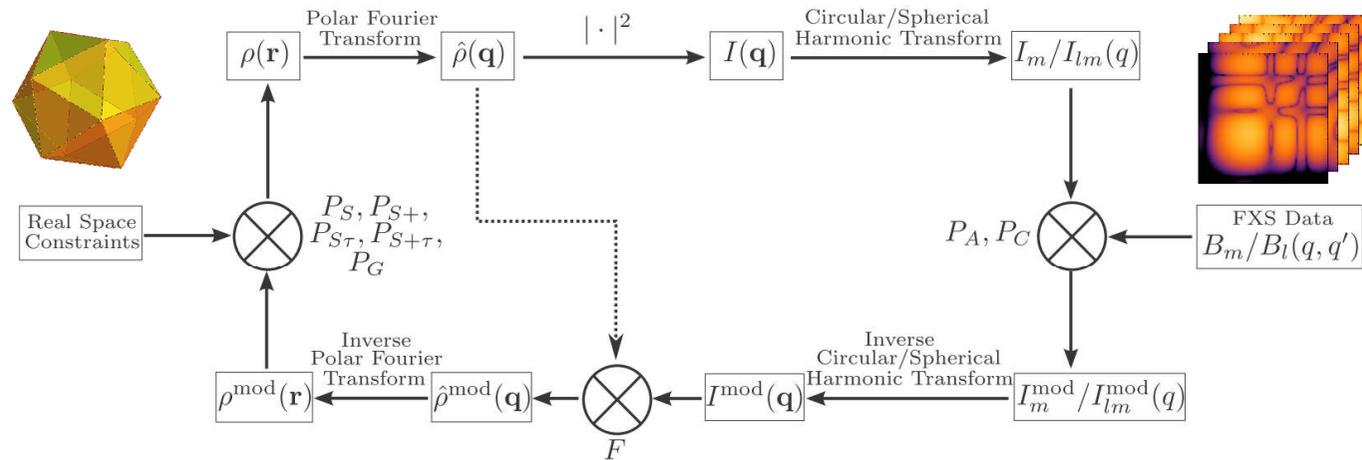


## Particle caking

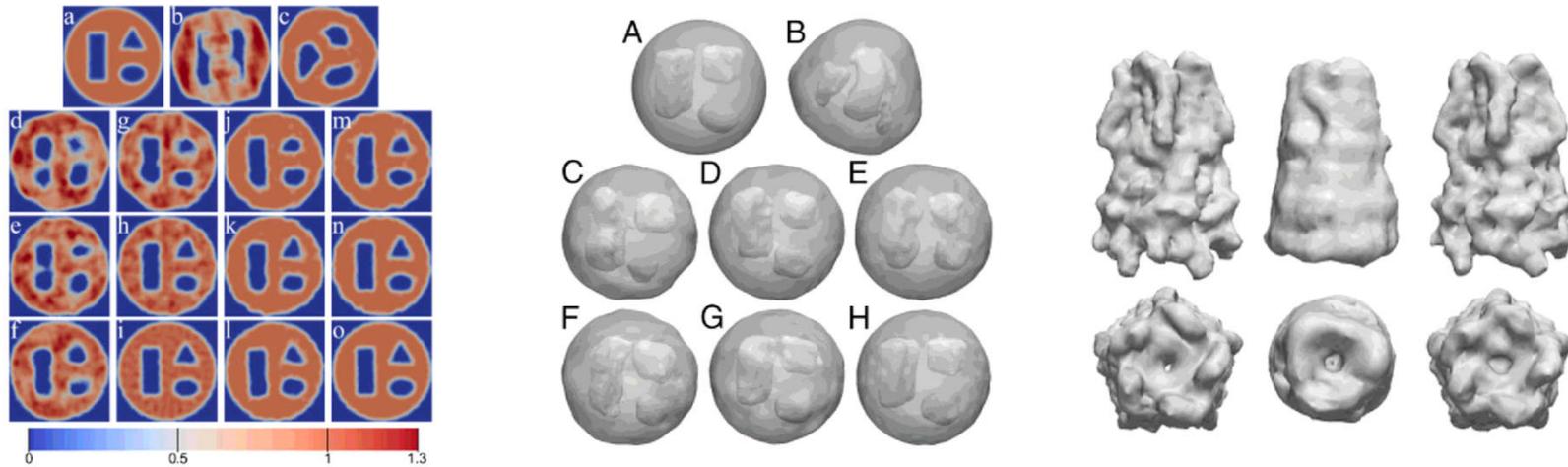
375



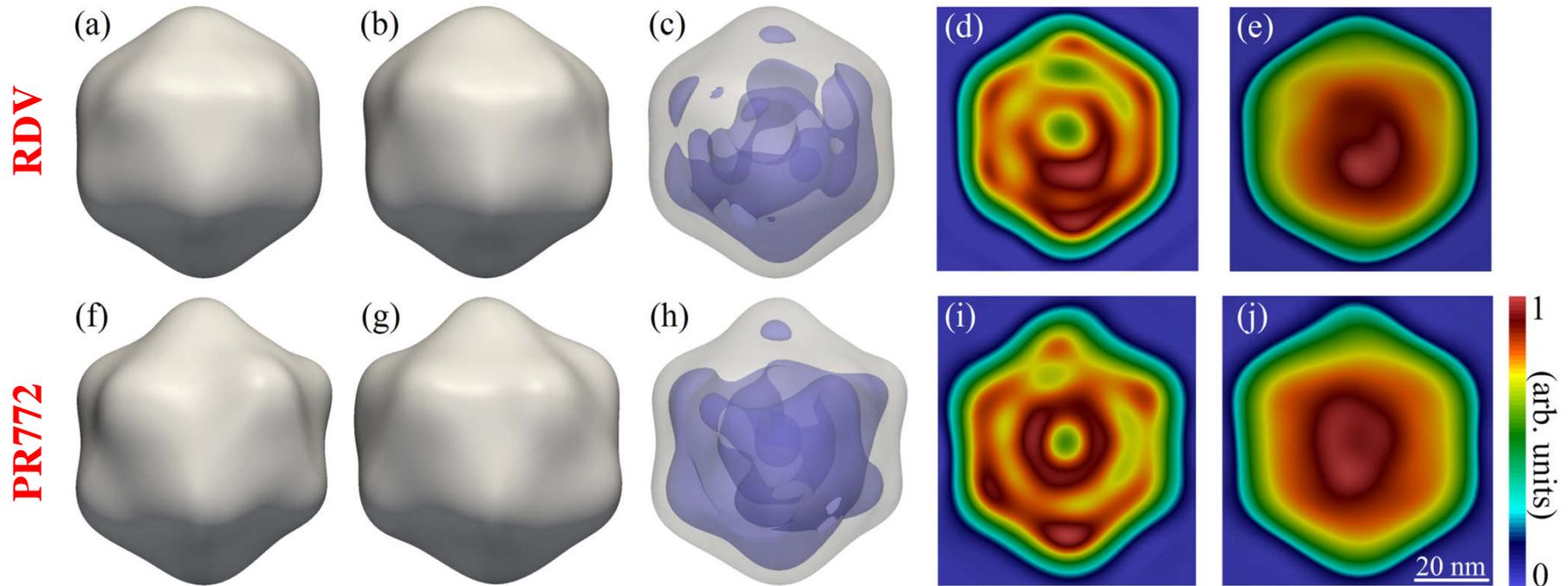
# Multitiered iterative phasing (MTIP) for structure recovery



**Fig. S7.** Flowchart of the multitiered iterative phasing procedure. The mod superscript denotes values that have been modified to agree with the data. One uses circular harmonic transforms,  $I_m$ , and  $B_m$  for the 2D case and spherical harmonic transforms,  $I_{lm}$ , and  $B_l$  for the 3D case.



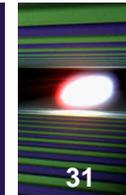
J. J. Donatelli, P. H. Zwart, J. A. Sethian, PNAS 112 (33), 10286 (2015)



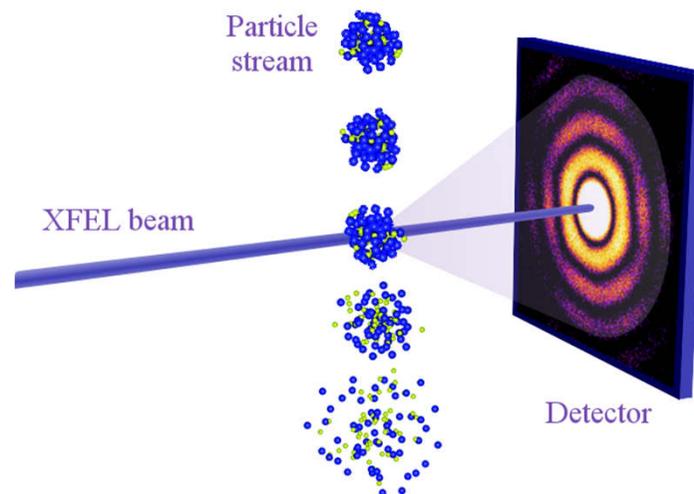
### Reconstructed images of RDV and PR772.

Two different views (corresponding to a 72 degree rotation about the top axis) of the reconstructed RDV (a,b) and PR772 (f,g) particles, as well as density plots showing nonuniformities in the internal distribution of material inside RDV (c) and PR772 (h), 2D slices through the center of the reconstructed densities for RDV (d) and PR772 (i), and 2D projections of the reconstructed densities for RDV (e) and PR772 (j).

# SPI experiments on PR772 virus particles at the AMO instrument, LCLS



31



### AMO\_86615 (August 2015):

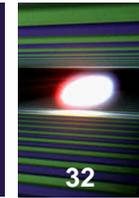
Photon energy:  $E=1.6$  keV  
 Sample-detector distance: 581 mm  
 Detector: pnCCD  
 Sample: PR772, RDV  
 Sample delivery: GDVN

### AMO\_06516 (April 2016):

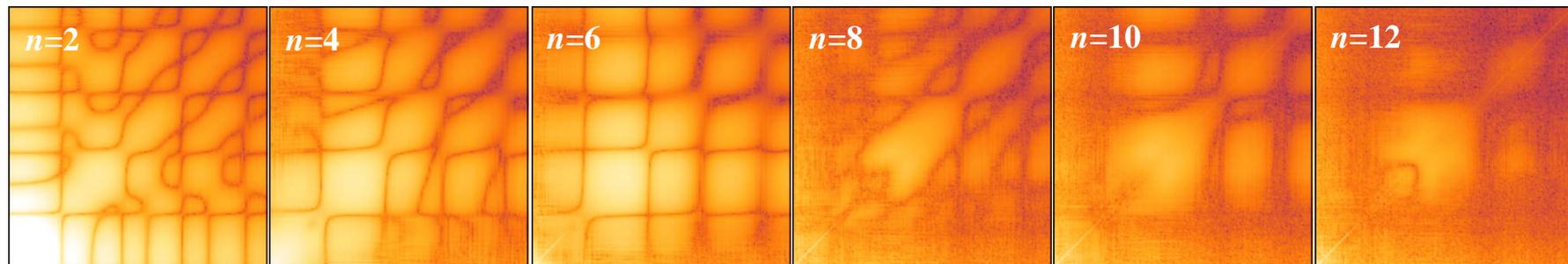
Photon energy:  $E=1.7$  keV  
 Sample-detector distance: 257 mm  
 Detector: pnCCD  
 Sample: PR772  
 Sample delivery: GDVN

### AMO\_11416 (August 2016):

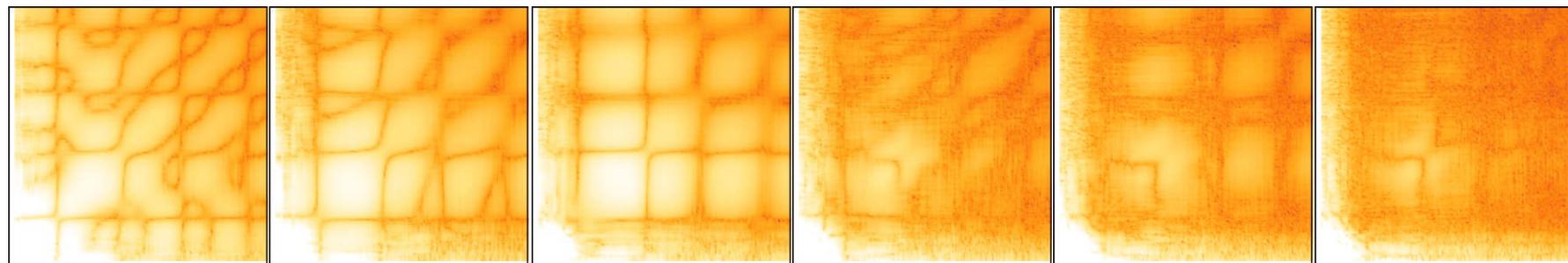
Photon energy:  $E=1.7$  keV  
 Sample-detector distance: 205 mm  
 Detector: pnCCD  
 Sample: PR772  
 Sample delivery: GDVN



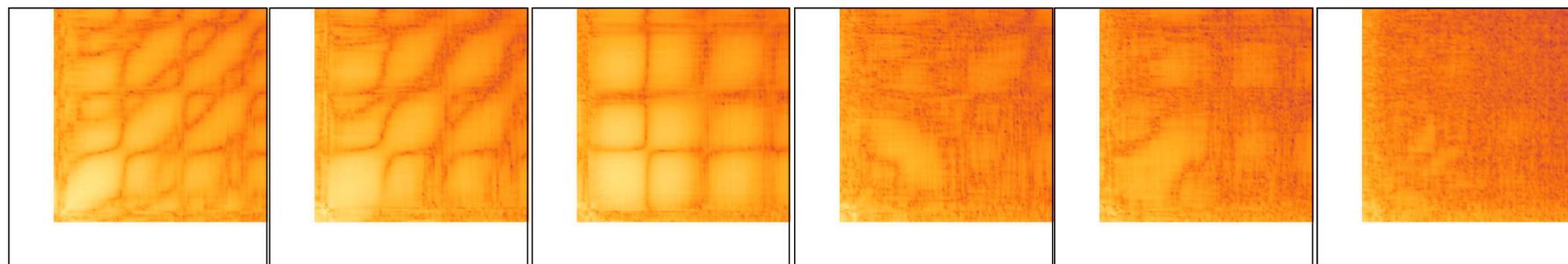
## AMO\_86615 (PD=3nm, M=2705 patterns)



## AMO\_06516 (PD=3nm, M=621 patterns)

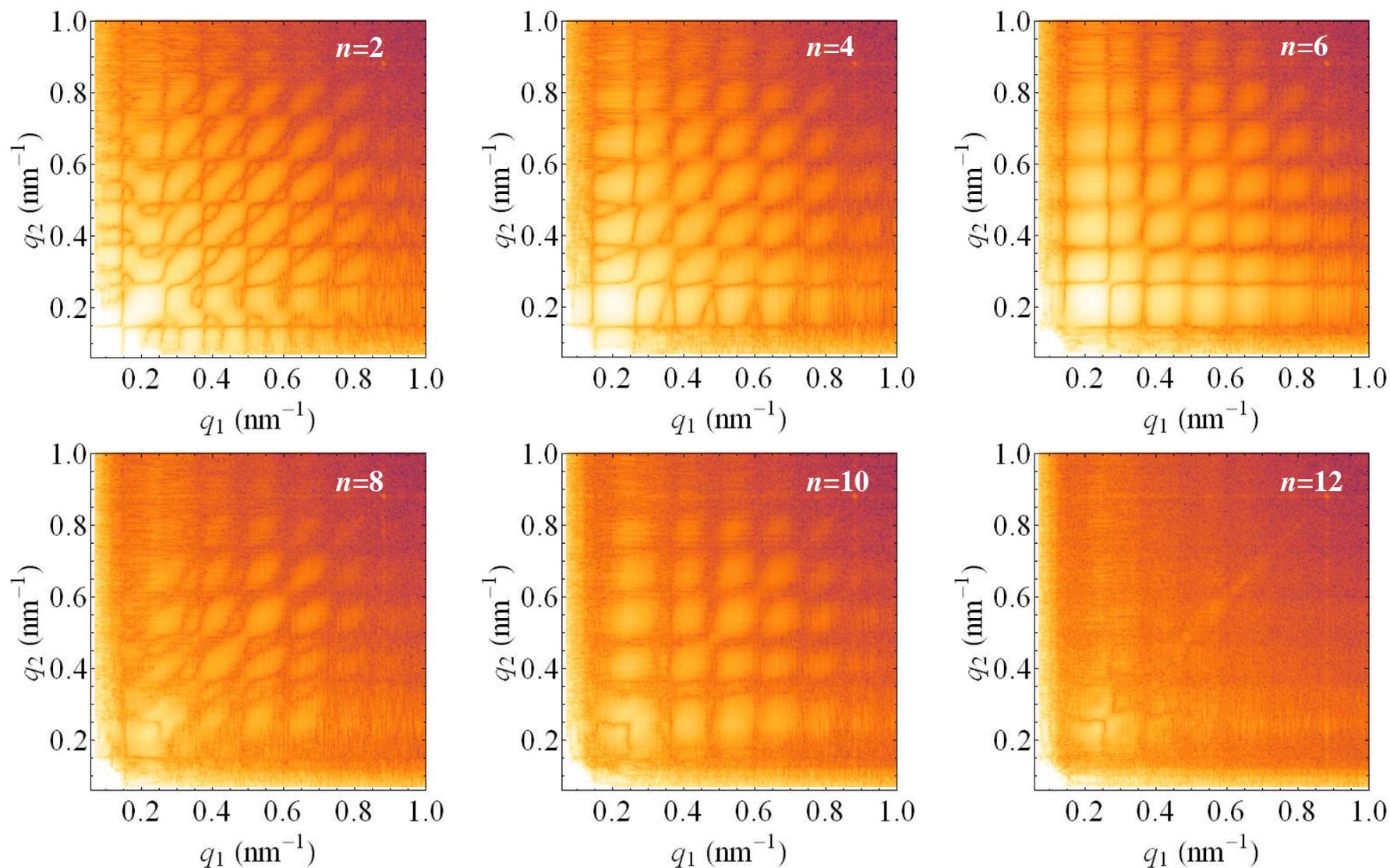
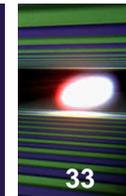


## AMO\_11416 (PD=3nm, M=115 patterns)



# Experimental 2D correlation maps:

AMO\_06516 data, full range





- FXS approach offers an alternative way for structural analysis of biological particles with an XFEL;
- FXS is, in the simplest case, a generalization of SAXS; In more general case it may give several orders of magnitude increase of information content as compared to SAXS;
- Cross-correlation functions are valuable statistical means, which can be conveniently used for model-based and *ab-initio* structure recovery;
- First application of FXS to biological particles at an XFEL demonstrates substantial potential of the technique for the future studies of structure and dynamics of biomaterials.



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E. Weckert

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more from the  
SPI  
initiative

## Thank you for your attention