

Growth of the number of periodic points for meromorphic maps

~~Joint work with T.-C. DINH & T.-T. TRUONG & D.-V. VU.~~

First of all, I would like to thank the organizers for the wonderful invitation. In this talk I will speak about some recent joint work with T.-C. DINH & T.-T. TRUONG & D.-V. VU. Here is the plan of my talk:

which consists of 3 parts. In the first part (I. Introduction) I will state the main questions, ~~prepare some background and~~ and state the main results. The shaded part (III. Main Tools⁺) is devoted to the two theories developed which have been developed by T.-C. Dinh & N. Sibony. ^{Sketchy proof} of intersection of positive closed currents. These theories will play a key role in talk.

In the last part (IV. Sketch of the proof) I will give a sketchy proof of the main results.

Now let me begin with the first part. Here is the settings of our problems.

- Let X be a compact Kähler manifold of dim k .
- Let $f: X \rightarrow X$ be a dominant holomorphic/meromorphic map or correspondence (Correspondence = multivalued map).
- Let $f^n := f \circ \dots \circ f$ (n times) be the iterate of order n of f .
- Let \mathbb{Q}_n be the set of isolated periodic points of period n . (i.e. isolated solutions of $f^n(x) = x$), counting with multiplicity.

We address the following problems:

Problem 1

Compute or estimate $\#\mathbb{Q}_n$

More however, we often consider the following easier variant of this problem

Problem 1'

Find a good upper bound for $\#\mathbb{Q}_n$.

We will also address ^{very} briefly the following problem although we do not have time to discuss it in length in this talk.

Problem 2:

Study the distribution of \mathbb{Q}_n as $n \rightarrow \infty$

Here is our strategy for these two problems

- ① Get a good upper bound for $\# Q_n$. (Problem 1').
- ② Construct a good family of periodic points using tools from dynamics or complex analysis.

In this way, we will have a good lower bound for $\# Q_n$ together with an equidistribution property. It is worthy noting that the problems are not completely solved even for holomorphic automorphisms in dim 2, e.g. when X is a K3 surfaces.

Now we are able to state the first main result.

Theorem A (Dinh-Sibony ('2014), Dinh-N-Truong ('2015))

Let $f: X \rightarrow X$ be a hol/mero map/correspondence on a compact Kähler manifold X . Then,

- f is an Arith-Mazur map: $\# Q_n$ grows at most exponentially fast;
- The exponential growth is bounded by the algebraic entropy ~~height~~.
- In many cases (e.g. Hénon maps), Q_n is asymptotically equidistributed with respect to a canonical invariant proba measure, as $n \rightarrow \infty$.

We should mention the result of Kaloshin and al. ('2000, '2007) for smooth non-holo maps, $\# Q_n$ can grow arbitrarily fast. The first assertion of our theorem illustrates the essential difference between holomorphic maps and non-holo ones.

Here is our second main result.

Theorem B (Dinh-N-Vu '2017).

Let $f: X \rightarrow X$ be a hol map on a compact Kähler manifold X .

Assume that the action of f^* on Hodge cohomology is simple.

Then $\# Q_n \leq d^n + o(d^n)$ as $n \rightarrow \infty$,

where d is the main dynamical degree of f .

The notions appearing in the Theorem will be explained in Part II.

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Now we come to Part II of the talk:

Part II: Background + Strategy + Difficulty

- Let $f: X \rightarrow X$ be a ~~cont~~ mere correspondence of graph $\Gamma \subset X \times X$. So Γ is an effective ~~cycle~~ $\Gamma = \sum_j \Gamma_j$, Γ_j irreducible analytic set of $\dim k$ in $X \times X$. (We only consider here finite sums and the Γ_j 's are not necessarily distinct).

and $\pi_1^*: \text{the restriction of } \pi_1 \text{ and } \pi_2 \text{ to each } \Gamma_j \text{ is surjective.}$
~~This can be seen as a multivalued meromorphic map~~
For $x \in X$, $f(x) = \pi_2(\pi_1^{-1}(x) \cap \Gamma)$ and $f^{-1}(x) = \pi_1(\pi_2^{-1}(x) \cap \Gamma)$.

- If ϕ is a smooth (p, q) -form on X , we can define the pull-back $f^*(\phi) := (\pi_2)_* (\pi_1^*(\phi) \wedge [\Gamma])$. The push-forward of ϕ by f by f^* .
- They are L^1 (p, q) -form on X , not continuous in general.
- We cannot iterate this operation. However, f^*, f_* commute with the operator $\partial, \bar{\partial}$.
- Consequently, f^* induces a linear map on Hodge cohomology $f^*: H^{p, q}(X) \rightarrow H^{p, q}(X)$, we can iterate this.

Def: For $0 \leq p \leq k$, the limit

$$d_p := \lim_{n \rightarrow \infty} \| (f^n)^* : H^{p, q}(X) \rightarrow H^{p, q}(X) \|^{\frac{1}{n}}.$$

- is called "the dynamical degree" of order p of f .
- $h_a := \max \log d_p$ is the algebraic entropy of f .

It is relevant to recall here a classical result due to Dinh & Sibony ('2005).

Dinh-Sibony ('2005) ~~for $0 \leq p \leq k$~~

- For $0 \leq p \leq k$, d_p exists, is finite, is bi-meromorphic invariant.

- The topological entropy of f satisfies $h_f(f) \leq h_a(f)$. (Gromov '1977)

Def: The exponential growth of f ~~is~~ ^{the number of isolated periodic points of f is}
$$\lim_{n \rightarrow \infty} \frac{1}{n} \log \# Q_n$$

By Thm A,

$$\lim_{n \rightarrow \infty} \frac{1}{n} \log \# Q_n \leq d_a.$$

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No ~~is~~ is m

Khovanskii-Teissier-Gromov: ($s \mapsto \log d_s$ is concave on $0 < s < k$)
 So there are $0 < p \leq p' \leq k$ such that
 $d_0 < \dots < d_p = \dots = d_{p'} > \dots > d_k$.
 $d_{p'}^2 \geq d_{s_1} d_{s_2}$)

Def:

- * f is with dominant topological degree if $p=p'=k$, i.e., $d_k > d_s$ $\forall 0 < s \leq k-1$.
- * The action of f^* on Hodge cohomology is simple if this action admits a unique eigenvalue of modulus d_p which is moreover a simple eigenvalue. In this case $d := d_p$ is the main dynamical degree.

Ex: Hol maps on \mathbb{P}^k .

Difficulty + Strategy

- * Periodic points of period n correspond to the intersection of the graph Γ_n of f^n with the diagonal $\Delta = \{(x, x) : x \in X\}$. It may have positive dimension.
- * Recall that Q_n is the set of isolated periodic points of period n , counting with multiplicity.
- * By Lefschetz fixed points theorem: If $\dim(\Gamma_n \cap \Delta) = 0$ then $\# Q_n = |\Gamma_n| \cup |\Delta|$.

It is then easy to get an estimate on $\# Q_n$ using the dynamical degrees. As you could see, the

Main difficulty for Prob 1 is the contribution of positive dimension components of the set of periodic points. In other words, we need to control some intersection with dimension excess.

Strategy of the proof of Theorem B

- Consider the graph Γ_n of f^n as a positive closed current in $X \times X$.

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- Step 1 Prove that the limit $\Gamma_\infty = \lim_{n \rightarrow \infty} \overline{[\Gamma_n]}$ exists in the sense of current. More concretely, we prove by the super-potential theory that the two main dynamical currents $T_+ = \lim_{n \rightarrow \infty} \frac{(f^n)(w)}{d^n}$ and $T_- = \lim_{\text{boudary}} \frac{(f^n)(w^{k-p})}{d^n}$ exist. Hölder continuous al by the fact that the action of f^* on Hodge cohomology.
- We show that $\Gamma_\infty \cap T_-$ is well-defined "simply" in the sense of super-potential theory.
 - We show that $T_+ \cap T_- \rightarrow$ horizontal dimension $\dim_K T_+ \otimes T_-$ of tangent currents.

Step 2: We deal with the dimension excess in our strategy as follows.

- We prove that the dimension excess of $\Gamma_\infty \cap \Delta$ is as expected 0.
- We need an upper-semicontinuity for the intersection. This, combined with the previous assertion and the limit $\overline{[\Gamma_n]} \rightarrow \Gamma_\infty$, implies that the dimension excess for $\overline{[\Gamma_n]} \cap \Delta$ can be asymptotically 0 as $n \rightarrow \infty$. Consequently, the density of dimension 0 of $\overline{[\Gamma_n]} \cap \Delta$ is asymptotically smaller than the density of dimension 0 of $[\Gamma_\infty] \cap \Delta$. In other words,

$$\lim_{n \rightarrow \infty} \frac{\#\mathcal{Q}_n}{d^n} \leq [\Delta]^0$$

This implies $\#\mathcal{Q}_n = d^n + o(d^n)$. \square .

Now we speak of the two main tools: super-potentials and density of positive closed currents.

III Main Tools:

III.1 Super-potentials: (Dinh-Sibony 2009-2010)

The starting point is that the pluripotential theory is well-developed for positive closed currents of bidegree $(1,1)$, thanks to the notion of p.s.h. functions. More precisely, if T is a positive closed $(1,1)$ current, then we can work locally $T = \bar{d}d^c u$ in the sense of currents, where u is p.s.h. function and $d^c = \frac{i}{2\pi}(\partial - \bar{\partial})$. Working with a pointwise defined function is more comfortable than working directly with a current. For example, it allows more operations like multiplying a positive current S by u . One just has to assume that u is integrable with respect to the trace measure of S . But when T is of higher bidegree, the potential is just an L^1 -form, one cannot consider their wedge-product with S if S is singular.

Super-potentials are functions which play the role of quasi-potentials for positive closed current of arbitrary bidegree. For $0 \leq p \leq k$, let

\mathcal{C}_p be the convex cone of positive closed currents on X .

\mathcal{D}_p be the real vector space generated by \mathcal{C}_p .

\mathcal{D}_p° = subspace of \mathcal{D}_p of currents belonging to the class

$O_{\text{in}} \text{H}^{1,0}(X, \mathbb{R})$. \mathcal{D}_p° = smooth forms of \mathcal{D}_p° .

Consider the \star -norm on \mathcal{D}_p :

- for $S \in \mathcal{C}_p$: $\|S\|_\star := |\langle S, \omega^{k-p} \rangle|$, which is the mass of S .

- for $S \in \mathcal{C}_p$: $\|S\|_\star = \inf_{S^\pm \in \mathcal{C}_p} (\|S^+\| + \|S^-\|)$.

$$S = S^+ - S^-$$

$$S^\pm \in \mathcal{C}_p$$

Consider now a current $R \in \mathcal{D}_p^\circ$ and a $(k-p, k-p)$ -current

U_R , which is a potential of R i.e. $\boxed{\bar{d}d^c U_R = R}$.

Let $\alpha = (\alpha_1, \dots, \alpha_h)$ with $h = \dim \text{H}^{1,0}(X, \mathbb{R})$ be a fixed family of real smooth closed $(1,0)$ -forms on X such that the family of classes

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$\{x_1, \dots, x_h\}$ is a basis of $H^k(X, \mathbb{R})$. By adding to U_p a suitable closed smooth form, we can assume that $\langle U_p, x_i \rangle = 0$ for $i = 1, \dots, h$. In this case we say that U_p is α -normalized.

Def: Let $T \in \mathcal{D}_p$. The α -normalized super-potential U_T of T is the function $U_T : \mathcal{D}_{k-p+1}^0 \rightarrow \mathbb{R}$ given by $U_T(R) := \langle T, U_R \rangle$, where U_R is an α -normalized smooth path of R . We say that T has bounded/continuous/Hölder superpotential if U_T can be extended to a function on \mathcal{D}_{k-p+1} which is bounded/continuous/Hölder with respect to the \star -topology. $U_T(R)$ does not depend on the choice of an α -normalized U_R .

Assume that T has a continuous super-potential. Take any $S \in \mathcal{D}_q$. Define $\{T\} = \sum_{i=1}^h a_i \{x_i\}$.

Define $T \wedge S$ to be the real $(p+q, p+q)$ current satisfying

$$\langle T \wedge S, \phi \rangle := U_T \left(\frac{\partial}{\partial \phi} \phi \wedge S \right) + \sum_{j=1}^h a_j \langle x_j, \phi \wedge S \rangle$$

$$\begin{aligned} \langle T \wedge S, \phi \rangle & \text{ is a real smooth } (k-p-q, k-p-q) \text{-form.} \\ &= \underbrace{\langle T - \sum_{j=1}^h a_j x_j, \phi \wedge S \rangle}_{\in \mathcal{D}_p^0} = \langle dd^c U_T, \phi \wedge S \rangle = \langle U_T, dd^c(\phi \wedge S) \rangle \\ &= \langle U_T, dd^c \phi \wedge S \rangle = \langle U_T(dd^c \phi), S \rangle, \end{aligned}$$

Now we come to

III.2 Density of positive closed currents along a manifold

Consider a manifold $V \subset X$

Consider $n: E \rightarrow V$ the normal vector bundle of V in X .

Denote by \bar{E} the natural compactification of E , i.e.

$\bar{E} = \mathbb{P}(E \oplus \mathbb{C})$, where \mathbb{C} is the trivial line bundle over X .

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For $\lambda \in \mathbb{C}^*$ denote by A_λ the multiplication by λ on the fibers of E .
 $A_\lambda(u) = \lambda u$ for $u \in E_a$ and $a \in V$.

Let $T \in \mathcal{E}_p$. Then the limit $(A_\lambda)_x T$ does not exist in general. However,

Theorem (Bis + Sibony '2012)

All limit values of $(A_\lambda)_x T$ are in the same color class in $H^{*,*}(E, \mathbb{R})$.

We call this color class the density of T along V .

Remark When V is a single point a , $\bar{E} = \mathbb{P}^k$ at.

$\{\text{the density of } T \text{ along } V\} \equiv \mathcal{D}\{w_{FS}\}$, $\lambda \in \mathbb{R}^+$

χ is the Lelong number of T at a .

We have an analytic counterpart of the Fulton intersection theory.

We also the density of positive closed currants enjoys a upper-semi continuous property. as in the case of Viehweg's

Finally we explain some points in the proof of Theorem B.

There are two new results of independent interest.

Prop 1: Let $T \in \mathcal{E}_p$, $S \in \mathcal{E}_q$ and T has a continuous superpotential. Then $T \wedge S = T \wedge S$.

This proposition allows us

Prop 2: Domination principle for positive closed currents.

Let $T \in \mathcal{E}_p$ s.c. The superpotential of T is bounded/continuous/Hölder continuous. Let $S \in \mathcal{E}_p$ s.c. $S \leq T$.

Then the superpotential of S is also bounded/continuous/Hölder continuous. This prop is needed in order to show that:

$$d^{-n} [\Gamma_n] \xrightarrow{n \rightarrow \infty} T^+ \otimes T^-$$

$\dim V = l$. $\dim X = k$.

Subject :

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R target curve
 $\{R\} = K^V(T)$ target div of T .

$H^*(\bar{E}, \mathbb{C})$ free module $H^*(V, \mathbb{C})$ -module
generated by $1, h, \dots, h^{k-l}$.
- h fan logical $(1, 1)$ class a \bar{E} .

$$\underline{\text{class}}(R) = \sum_{K_j^V(T)} n^*(K_j^V(T)) \otimes h^{j-l+g}$$
$$K_j^V(T) \in H_{j,j}(V, \mathbb{C})$$

Send curve T/h max j . $-K_j^V(T) \neq 0 \rightarrow$ horizontal divisor

$$\boxed{T_n \rightarrow T \xrightarrow{\quad} \text{horizontal dim } (T) = \lim_{n \rightarrow \infty} K_s^V(T_n) \leq K_s^V(T)}$$

$$T \wedge S = A(T \otimes S) \xrightarrow{\quad} R \text{ unique.}$$

$$(T \wedge S) \text{ unique on } \Delta \underset{\parallel}{\Delta} \underset{\times}{\Delta} : \boxed{\pi^*(T \wedge S) = R}$$