

Efficient Methods for Homogenization of Random Heterogeneous Materials

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O Background information

O Improving convergence rate with Richardson extrapolation

O Robin boundary condition for high-contrast materials

 \boldsymbol{O} Conclusions



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Many practical problems have multiple-scale feature.

- Composite materials
- Porous media





(b) Composites at the microscale



(c) Porous media at the microscale



BVP with random rapidly oscillating coefficients

$$\begin{cases} -\nabla \cdot \left(a^{\varepsilon}(x,\omega) \nabla u^{\varepsilon}(x,\omega)\right) = f(x) & \text{in } D \subset \mathbb{R}^{d} \\ u^{\varepsilon}(x,\omega) = g(x) & \text{on } \partial D \end{cases}$$

- $\varepsilon > 0$ is a parameter describing the multiscale feature.
- ω denotes a realization of the random coefficient tensor in probability space (Ω, \mathcal{F}, P) .
- f(x) and g(x) are macroscopic functions.



• It is difficult to solve the multiscale problem with

traditional numerical methods.

(Finite element method, Boundary element method, etc.)

FEM $h \ll L_{\text{micro}} \ll L_{\text{macro}}$

• Macroscopic behavior of the heterogeneous material may be approximated by that of a fictitious homogeneous one.



Typical multiscale methods for the BVP

- Asymptotic homogenization method (*Lions* 1978)
- Variational multiscale method (*Hughes* 1995)
- Multiscale finite element method (*Hou & Efendiev* 1997)
- Heterogeneous multiscale method (*E & Engquist* 2003)



Homogenized BVP

$$\begin{cases} -\nabla \cdot \left(\overline{a} \nabla u^{0}(x)\right) = f(x) & \text{in } D \subset \mathbb{R}^{d} \\ u^{0}(x,) = g(x) & \text{on } \partial D \end{cases}$$

- \bar{a} is called effective (or homogeneous) coefficient tensor.
- \overline{a} is a constant tensor if a^{ε} is a statistically homogeneous ergodic tensor, i.e., $a^{\varepsilon}(x, \omega) = a(y, \omega), y = x/\varepsilon$.



Effective coefficients \bar{a}

- Defined in the probability space (Ω, F, P), which cannot be computed by numerical methods
- According to the ergodic theorem, \overline{a} is approximated by

$$\overline{a} \cdot e_i = \lim_{L \to \infty} \left\langle a(y, \omega) \nabla u(y, \omega) \right\rangle$$

where u is the solution of

$$-\nabla \cdot \left(a\left(y,\omega\right) \nabla u\left(y,\omega\right) \right) = 0 \quad \text{in } Y_L = \left[0,L\right]^d$$

which satisfies $\langle \nabla u \rangle = e_i$, and e_i $(i = 1, \dots, d)$ are canonical bases in \mathbb{R}^d .



Auxiliary problem in volume element

$$-\nabla \cdot \left(a\left(y,\omega\right) \nabla u\left(y,\omega\right) \right) = 0 \quad \text{in } Y_L$$

• Dirichlet boundary condition (DBC)

$$u(y,\omega) = \xi_i \cdot y \quad \text{on } \partial Y_L, i = 1, \cdots, d$$

• Neumann boundary condition (NBC)

$$a(y,\omega)\nabla u(y,\omega)\cdot n = \eta_i\cdot n \quad \text{on } \partial Y_L, i = 1,\cdots, d$$

Homogenization





Figure: The main procedure of multiscale method for random heterogeneous materials





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Some results

- With proper boundary condition, approximate effective coefficients converge to the true effective coefficients as $L \rightarrow \infty$. (*Bourgeat* 2004)
- (Numerical) convergence rates of DBC and NBC approximations are O(1/L). (*Yue* 2007)
- DBC overestimates the true effective coefficients, while NBC underestimate the true effective coefficients.
 (Dirichlet-Neumann bounds)



• Sufficiently large volume elements are always needed,

which is time-consuming and high-demanding for computer memory.



Method



Dirichlet boundary condition

$$u(y,\omega) = \xi_i \cdot y$$
 on ∂Y_L , $i = 1, \cdots, d$

It satisfies

$$\langle \nabla u(y,\omega) \rangle = \xi_i$$

Approximate effective coefficients are calculated by

$$\overline{a}_{D,L}(\omega) \cdot \xi_i = \left\langle a(y,\omega) \nabla u(y,\omega) \right\rangle$$

Mathematical expectation $\mathbb{E}\bar{a}_{D,L}$ is then computed.



Neumann boundary condition

$$a(y,\omega)\nabla u(y,\omega)\cdot n = \eta_i\cdot n \quad \text{ on } \partial Y_L, i = 1,\cdots, d$$

It satisfies

$$\langle a(y,\omega)\nabla u(y,\omega)\rangle = \eta_i$$

Approximate effective coefficients are calculated by

$$\overline{a}_{N,L}(\omega) \cdot \langle \nabla u(y,\omega) \rangle = \eta_i$$

Mathematical expectation $\mathbb{E}\bar{a}_{N,L}$ is then computed.

Method



Theorem (Wu, Nie, Yang 2014)

Let $\mathbb{E}\bar{a}_{D,L}$ and $\mathbb{E}\bar{a}_{N,L}$ be the approximate effective coefficients under DBC and NBC respectively, and let \bar{a}_{∞} be the true effective coefficients, then we have

$$\mathbb{E}\bar{a}_{D,1} \geq \cdots \geq \mathbb{E}\bar{a}_{D,L} \geq \mathbb{E}\bar{a}_{D,2L} \geq \cdots \geq \bar{a}_{\infty}$$

and

$$\mathbb{E}\bar{a}_{N,1} \leq \cdots \leq \mathbb{E}\bar{a}_{N,L} \leq \mathbb{E}\bar{a}_{N,2L} \leq \cdots \leq \bar{a}_{\infty}$$



Assume the first-order convergence rates of DBC and NBC approximations are satisfied, then we have

$$\mathbb{E}\,\overline{a}_{_{BC,L}} - \overline{a}_{_{\infty}} = C\,\frac{1}{L} + O\left(\frac{1}{L^{\beta}}\right), \quad \text{ where } \beta > 1$$

Since

$$\mathbb{E}\,\overline{a}_{BC,2L} - \overline{a}_{\infty} = C\frac{1}{2L} + O\left(\frac{1}{2^{\beta}L^{\beta}}\right)$$

We get Richardson extrapolation sequence

$$R_{BC,L} - \overline{a}_{\infty} = \left[2\mathbb{E}\,\overline{a}_{BC,2L} - \mathbb{E}\,\overline{a}_{BC,L}\right] - \overline{a}_{\infty} = O\left(\frac{1}{L^{\beta}}\right)$$

Method





Figure: The main procedure of multiscale method combined with Richardson extrapolation technique



Microstructure

- (RMDF) random morphology description function, proposed by Vel et al. in 2010.
- Closely resemble actual micrographs manufactured by techniques including plasma spraying and powder processing.



(a) Microstructures generated by computer with different volume fractions of Al



(b) Actual Al/Al₂O₃ micrographs



Example

We predict effective thermal conductivity of actual AI/AI_2O_3 with different volume fractions of AI.



Figure: Microstructures with different sizes



Table: Properties of constituent materials

Material	AI	Al ₂ O ₃
<i>k</i> (W/mK)	233.0	30.0

Table: Convergence rates of approximate effective coefficients

V _{AI} (%)	10	30	50	70	90
DBC	1.08	0.93	0.95	0.99	1.04
NBC	1.02	0.99	0.82	0.95	0.98





Figure: Comparison of approximate effective coefficients by homogenization and extrapolation with increasing volume element size























Example

We predict effective mechanical properties of actual AI/AI_2O_3 with different volume fractions of Al.

The contrast ratio between the elastic modulus of Al_2O_3 and Al is 5.6.

Material	Al	Al ₂ O ₃
<i>E</i> (GPa)	70	393
ν	0.30	0.22

Table: Properties of constituent materials











Example

We predict effective mechanical properties of actual stainless steel/epoxy with different volume fractions of stainless steel. The contrast ratio between the elastic modulus of stainless steel and epoxy is high (more than 100).

Material	Stainless Steel (SS)	Resin
<i>E</i> (GPa)	193.8	1.31
ν	0.29	0.40

 Table:
 Properties of constituent materials





Figure: Comparison of Homogenization, Extrapolation and Variational bounds (Voigt-Reuss bounds and Hashin-Shtrikman bounds)





Figure: Comparison of Homogenization, Extrapolation and Variational bounds (Voigt-Reuss bonds and Hashin-Shtrikman bounds)





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- Many random heterogeneous materials have a high contrast of constituent properties.
 - a. Composite materials (e.g., CNT-reinforced polymers)
 - b. Porous media
- The high contrast leads to very broad Dirichlet-Neumann upper and lower bounds.



Mixed Dirichlet-Neumann boundary condition (DNBC)

approximations lie within the Dirichlet-Neumann bounds. (*Yue* 2007)

$$\begin{split} u(y,\omega) &= \xi_i \cdot y \quad \text{on } \partial_1 Y_L \\ a(y,\omega) \nabla u(y,\omega) \cdot n &= 0 \quad \text{on } \partial_2 Y_L \\ \end{split}$$

$$\begin{split} & \Gamma_4 \\ & \Gamma_4 \\ & \Gamma_1 \\ & \Gamma_1 \\ & \Gamma_3 \\ \end{split}$$

$$\end{split}$$

$$\begin{split} & \text{Where } \partial Y_L &= \overline{\partial_1 Y_L} \cup \overline{\partial_2 Y_L}, \\ & \text{And } i = 1, \cdots, d. \\ \end{split}$$



- Boundary conditions destroy the ergodicity of solutions to the auxiliary problem.
- When the ergodicity can be satisfied, we have $\mathbb{E}\bar{a}_L = \bar{a}_\infty$ for any *L*.



Figure: Volume elements with different sizes

Continuous properties

$$u_{\omega_1} = u_{\omega_2}$$
$$a\nabla u_{\omega_1} \cdot n_{\omega_1} = -a\nabla u_{\omega_2} \cdot n_{\omega_2}$$



Robin boundary condition (RBC)

$$a(y,\omega)\nabla u(y,\omega)\cdot n + \lambda u(y,\omega) = \eta_i \cdot n + \lambda \xi_i \cdot y \quad \text{on } \partial Y_i$$

where $\lambda \in (0, \infty)$ is the adjusting factor, and $i = 1, \dots, d$.

Theorem (Convergence)

Let $\mathbb{E}\bar{a}_{R,L}$ be the approximate effective coefficient tensor computed by the auxiliary problem with Robin boundary condition. Then

 $\lim_{L\to\infty}\bar{a}_{R,L}(\omega)=\bar{a}_{\infty} \quad \text{a.s.}$

Method





Figure: The main procedure of multiscale method combined with Robin boundary condition



Example

We consider effective coefficients of random checker-board microstructures. Each unit cell is occupied by matrix material or reinforcement material with probability p and 1-p, (0<p<1).



Figure: Random checker-board microstructures with different sizes

Methods for Homogenization of RHMs



For two-phase random heterogeneous materials,

$$a(y,\omega) = \begin{cases} a_1 \cdot I, & y \in Y_L^m \\ a_2 \cdot I, & y \in Y_L^r \end{cases}$$

has a high contrast, i.e.,

$$r = \frac{a_2}{a_1} \gg 1$$

Here, *I* is the identity matrix.





Figure: Approximate effective coefficients with different boundary conditions and different contrast ratios ($\lambda = 40$ in RBC)





Figure: Effect of adjusting factor in RBC on the accuracy of approximate effective coefficients





Figure: Approximate effective coefficients with different adjusting factor in RBC



- High contrast leads to large condition number of stiffness matrix, which may reduce the numerical accuracy of approximate effective coefficients.
- We discuss the effect of condition number on the accuracy of approximate effective coefficients.

Auxiliary problem with DBC

$$\begin{cases} -\nabla \cdot \left(a\left(y,\omega\right) \nabla u\left(y,\omega\right) \right) = 0 & \text{in } Y_L \\ u\left(y,\omega\right) = \xi_i \cdot y & \text{on } \partial Y_L, i = 1, \cdots, d \end{cases}$$

More Discussion





Figure: Effect of contrast ratio and mesh size on the condition number of stiffness matrix (constituent with high property in the center)

More Discussion





Figure: Effect of contrast ratio and mesh size on the condition number of stiffness matrix (constituent with low property in the center)

More Discussion



$\operatorname{cond}(A) = Crh^{-2}$							
		h = 1/ccond	/36 <i>N</i> ite	h=1/cond	/ 42 N _{ite}	h=1cond	/48 <i>N</i> ite
<i>r</i> = 100	CGM ICCG M	24472 1168	293 25	33306 1590	366 29	43106 2076	440 33
<i>r</i> = 1000	CGM ICCG M	218510 10325	474 29	289075 13517	612 33	363644 16902	768 37

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- Largest condition number is no more that 10⁶, while double-precision floating-point system guarantees more than 15 significant digits of freedom.
- Condition number has little effect on the accuracy of approximate effective coefficients.





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Richardson extrapolation is an effective technique for predicting effective coefficients of random heterogeneous materials.

• Since smaller volume elements are used, a lot of

computation time and computer memory could be saved.



Robin boundary condition is proposed for predicting effective coefficients of random heterogeneous materials.

- It provides much better approximate effective coefficients than other boundary conditions.
- It is more flexible than other boundary conditions because of the adjusting factor.
- It is more suitable for random heterogeneous materials with high contrast.

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Thank you for your attention !