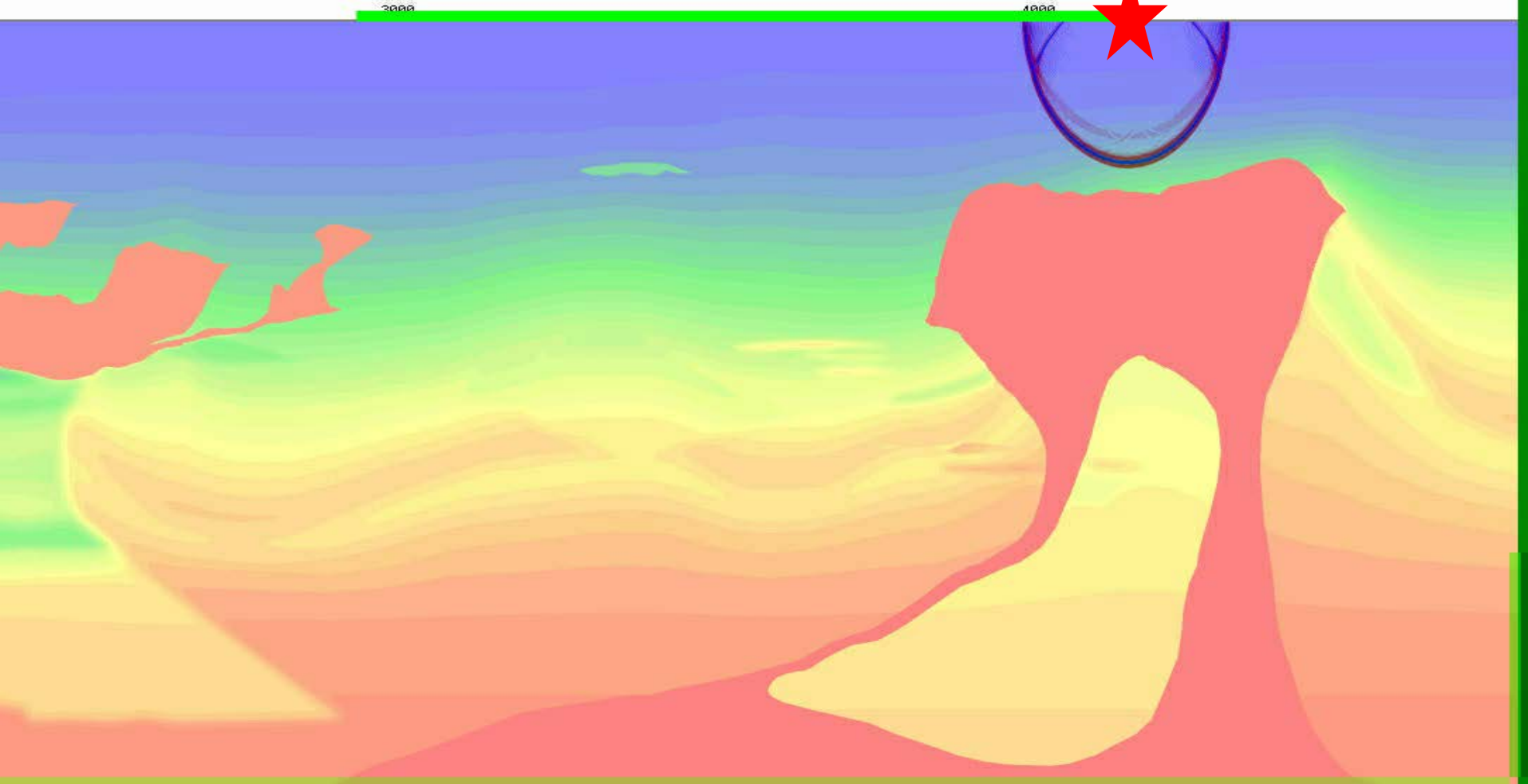


# Multicomponent elastic imaging: new insights from the old equations

Yunyue Elita Li\*, Yue Du, Jizhong Yang, and Arthur Cheng

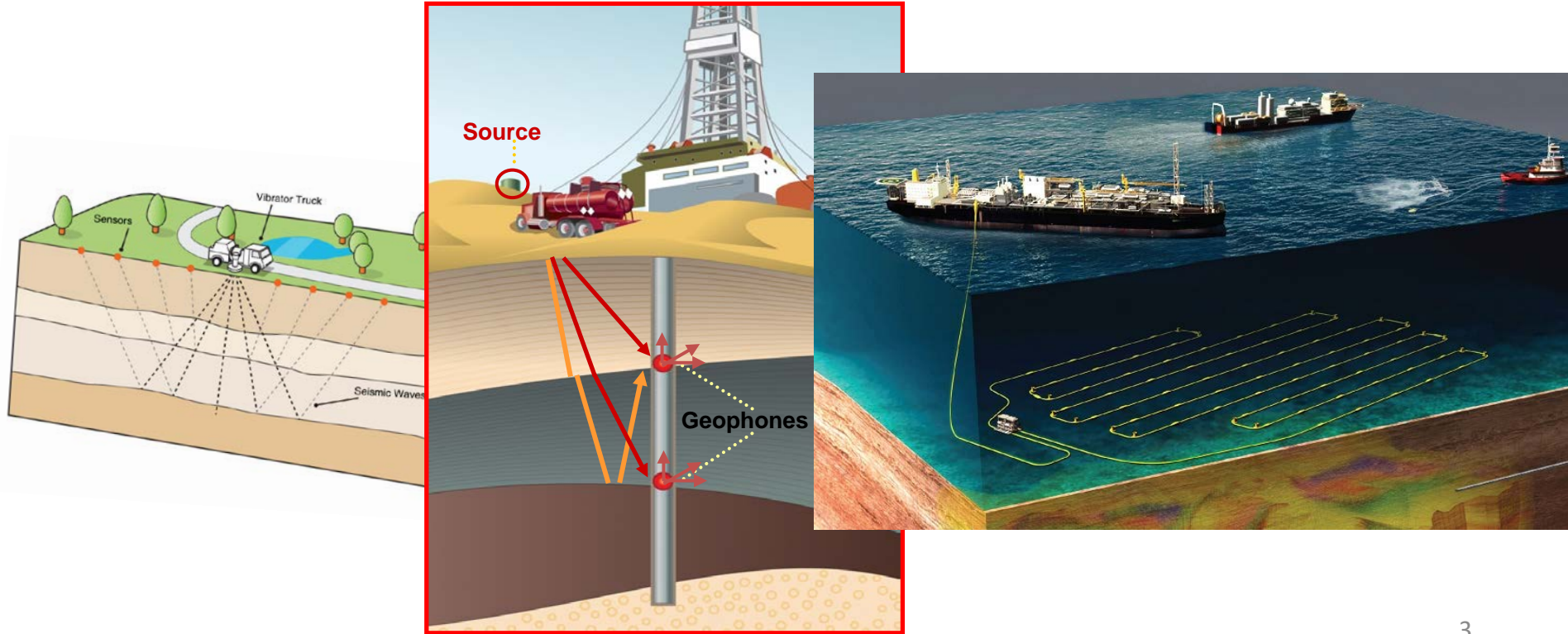
Singapore Geophysics Project

National University of Singapore

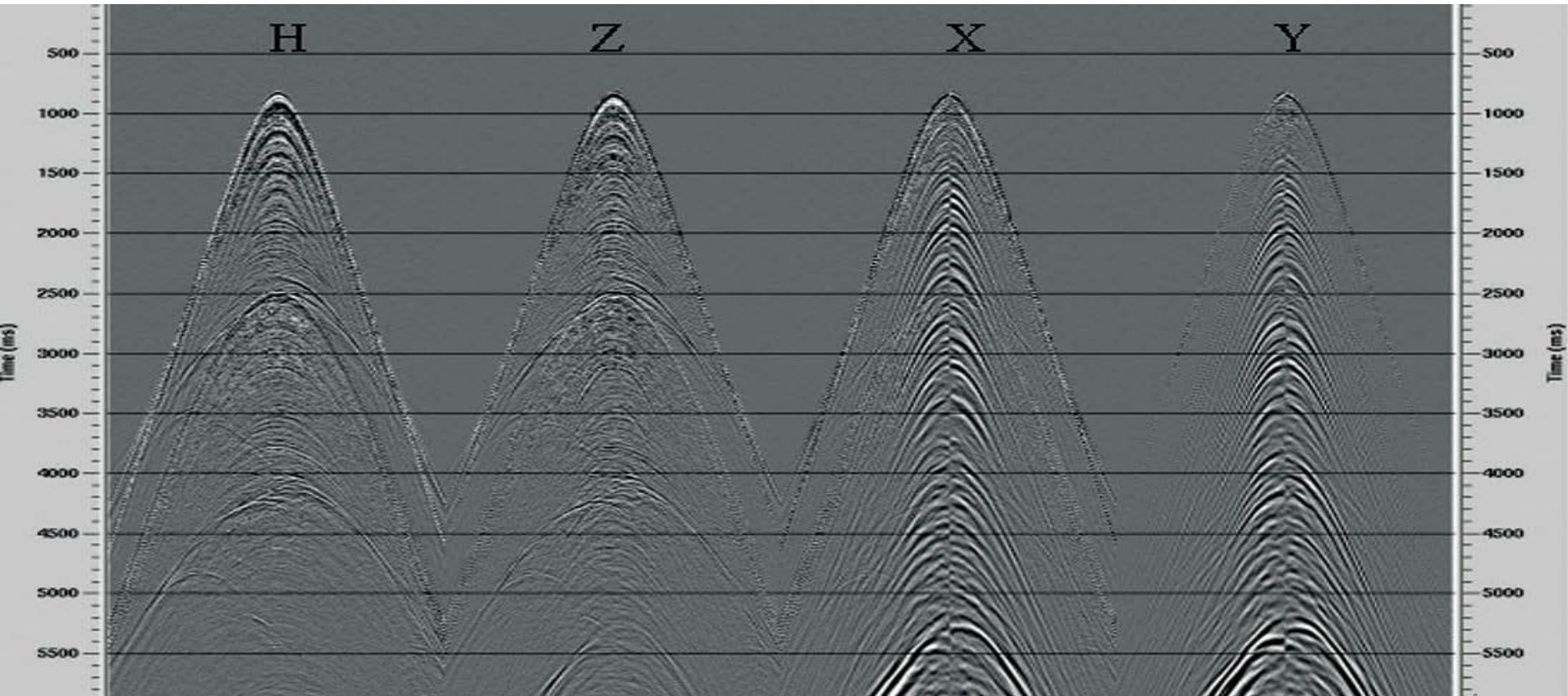


Simulation of a field scale seismic wave acquisition experiment

# Multicomponent data acquisition



# OBN acquisition: 4C data



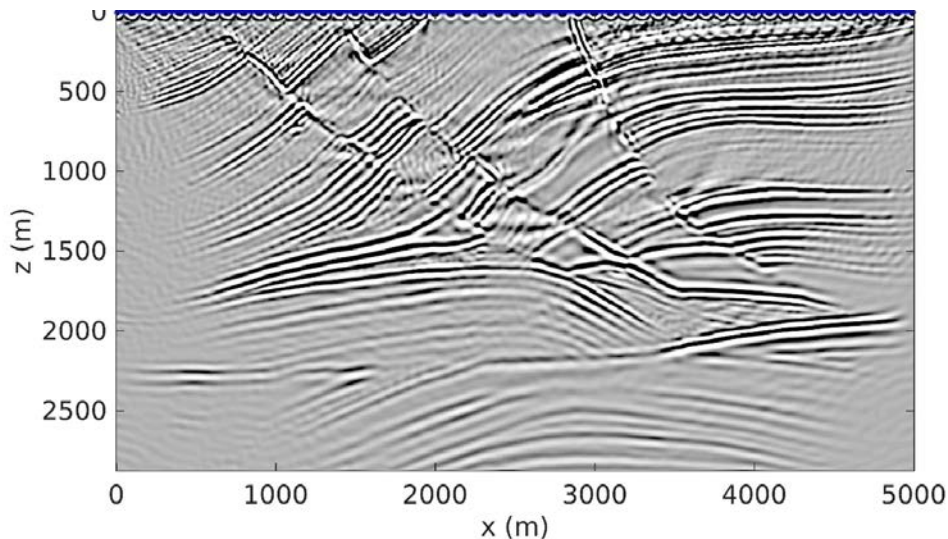
# Elastic imaging is not widely applied

- Large computational cost compared with acoustic imaging
  - 5 times in runtime and memory in 2D
  - 9 times in runtime and memory in 3D
- Deteriorated image for converted waves
  - Polarity reversal at normal incidence
  - Complicated, cumbersome, and add hock

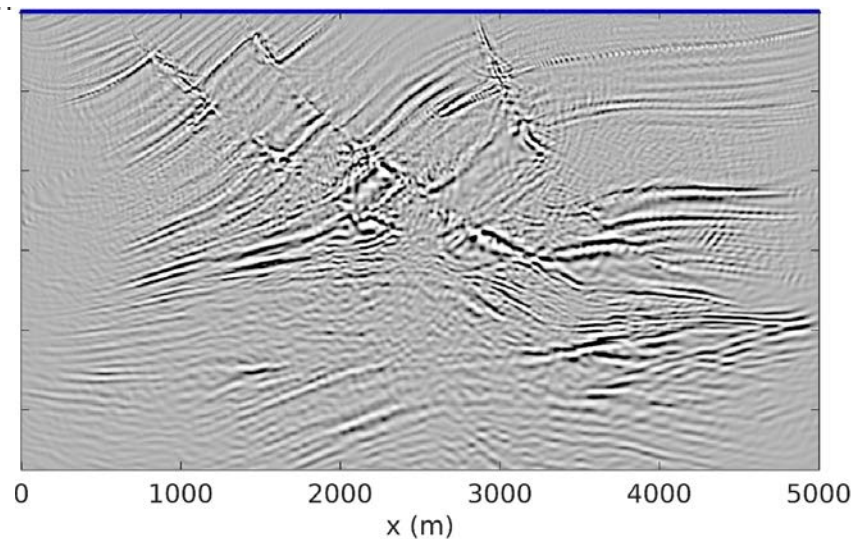


# Industry standard imaging algorithm

PP reflection image



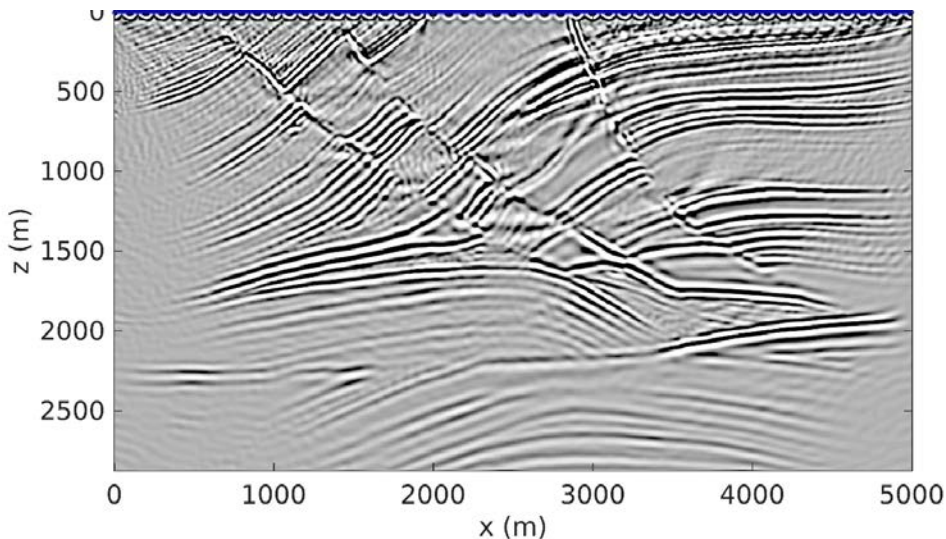
PS reflection image



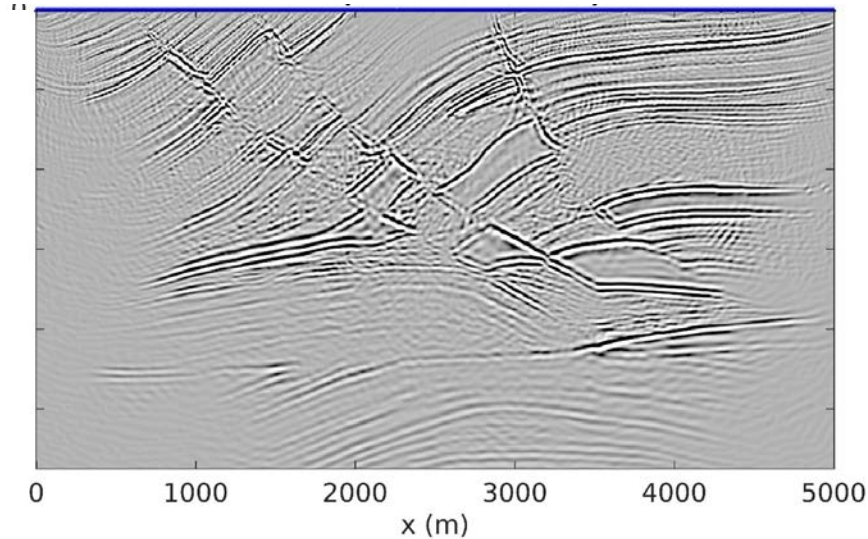
- ✧ Converted wave imaging appears noisier, less coherent, and challenging for joint interpretation
- ✧ Images are obtained with 5 times the computation and memory cost of the acoustic images

# Proposed imaging algorithm

PP reflection image



PS reflection image



- ✓ Converted wave imaging shows consistent geological features with **higher resolution**
- ✓ Imaging cost are reduced by **60%** in computation and **80%** in memory

# Outline

- Elastic wave equations
  - Revisit of the elastic wave equations
  - A new set of separated P- and S-wave equations
- The elastic imaging condition
  - PP and PS images from inverse problem formulation
- Discussions and conclusions



# Seismology 101: elastodynamic system

- Linear, isotropic, elastic medium (Aki and Richards, 1980)

Newton's Law:

$$\rho \frac{\partial^2 u_i}{\partial t^2} = \partial_j \tau_{ij} + f_i$$

$u_i$  particle displacement

$\tau_{ij}$  element of the stress tensor

$f_i$  force

Hooke's Law:

$$\tau_{ij} = \lambda \delta_{ij} \partial_k u_k + \mu (\partial_i u_j + \partial_j u_i) \quad \rho, \lambda, \mu \text{ density and Lamé constants}$$

- ✧ Need to propagate (and store) 5 fields in 2D, and 9 fields in 3D
- ✧ Cannot interpret the P- and S-wave directly from the equations

# Seismology 101: elastodynamic system

- The second-order system (Aki and Richards, 1980)

$$\rho \ddot{\mathbf{u}} = (\nabla \lambda)(\nabla \cdot \mathbf{u}) + \nabla \mu \cdot [\nabla \mathbf{u} + (\nabla \mathbf{u})^T] + (\lambda + 2\mu)\nabla \nabla \cdot \mathbf{u} - \mu \nabla \times \nabla \times \mathbf{u} + \mathbf{f}$$

$u_i$  particle displacement

$\rho, \lambda, \mu$  density and Lamé constants

- ✧ Need to propagate (and store) 3 fields in 2D, and 3 fields in 3D
- ✧ Require more strict stability condition
- ✧ Cannot interpret the P- and S-wave directly from the equations

# P- and S-wave separation in homogenous medium

- Assuming constant density and **smooth** Lamé constants

$$\ddot{\mathbf{u}} = \alpha \nabla \nabla \cdot \mathbf{u} - \beta \nabla \times \nabla \times \mathbf{u} + \mathbf{f}$$

$$\ddot{P} - \alpha \nabla^2 P = \nabla \cdot \mathbf{f}$$

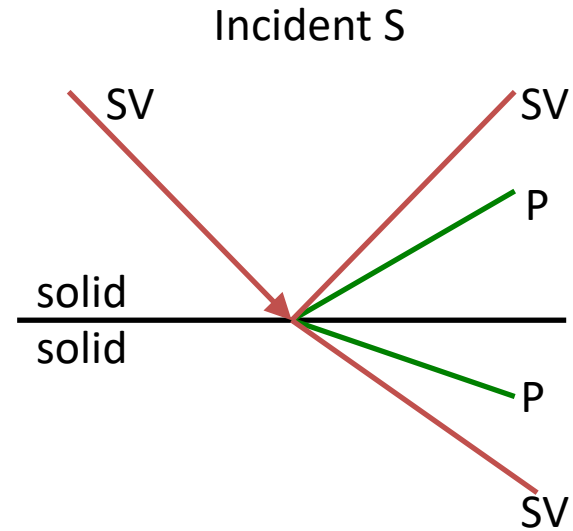
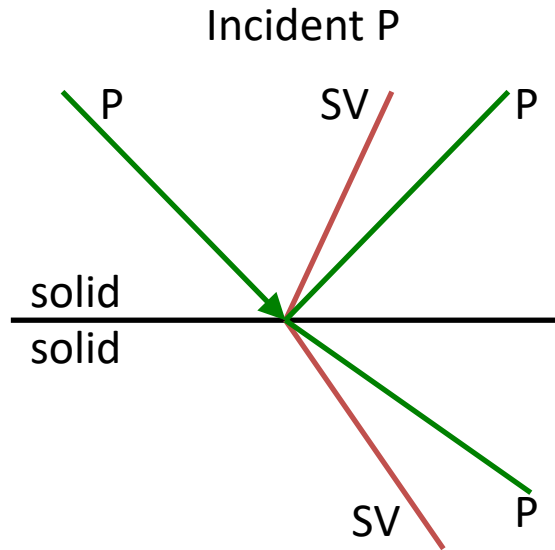
$$\ddot{\mathbf{S}} - \beta \nabla^2 \mathbf{S} = \nabla \times \mathbf{f}$$

$$\alpha = \frac{\lambda + 2\mu}{\rho} = V_p^2, \beta = \frac{\mu}{\rho} = V_s^2$$

$$P = \nabla \cdot \mathbf{u}, \quad \mathbf{S} = \nabla \times \mathbf{u}$$

- ✧ Fully decoupled P- and S-wave propagations
- ✧ Cannot interpret the mode-conversion directly from the equations

# Seismology 101: mode conversion



- ✧ Are these mode conversion types unconditional?
- ✓ New set of equations: clear mode conversion and its condition

# New set of separated P- and S-wave equations

$$\ddot{P} - \alpha \nabla^2 P = \nabla \cdot \mathbf{f}$$

P-wave propagation

$$+ P \nabla^2 \alpha + 2 \nabla \alpha \cdot \nabla P$$

$$- 2 P \nabla^2 \beta$$

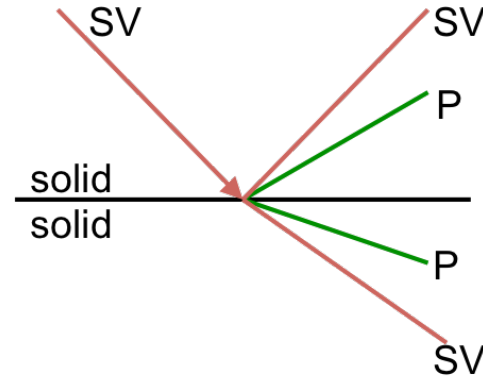
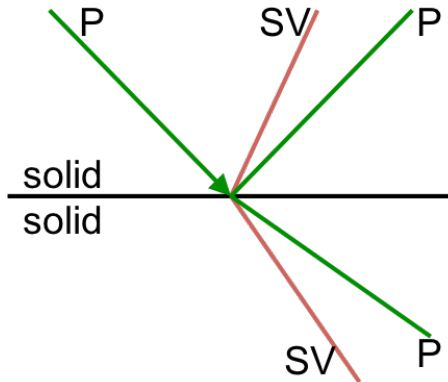
$$- 2 \nabla \beta \cdot \nabla \times \mathbf{S}$$

← Source term

← P-wave interacts with  $V_p$  boundary

← P-wave interacts with  $V_s$  boundary

← S-wave interacts with  $V_s$  boundary



# New set of separated P- and S-wave equations

$$\ddot{\mathbf{S}} - \beta \nabla^2 \mathbf{S} = \nabla \times \mathbf{f}$$

← Source term

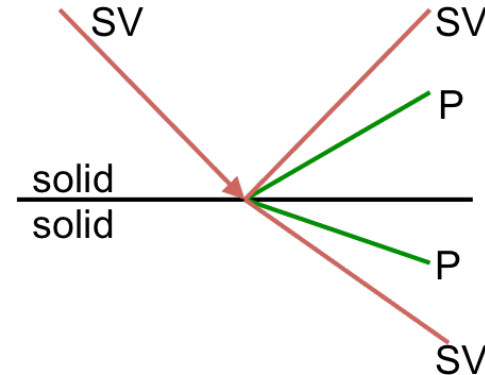
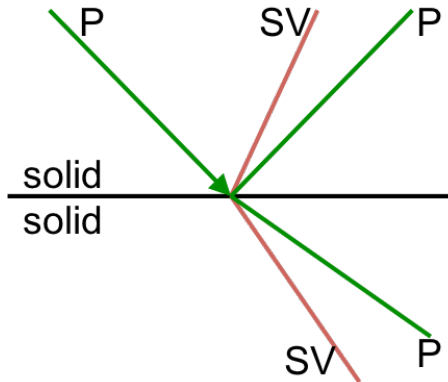
$$+ \nabla \beta \cdot \nabla \mathbf{S} - (\nabla \beta) \times (\nabla \times \mathbf{S})$$

← S-wave interacts with  $V_s$  boundary

S-wave propagation

$$+ 2(\nabla \beta) \times (\nabla P)$$

← P-wave interacts with  $V_s$  boundary



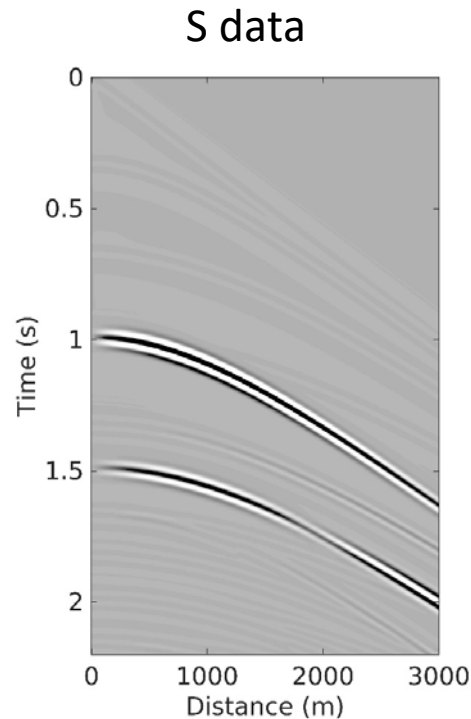
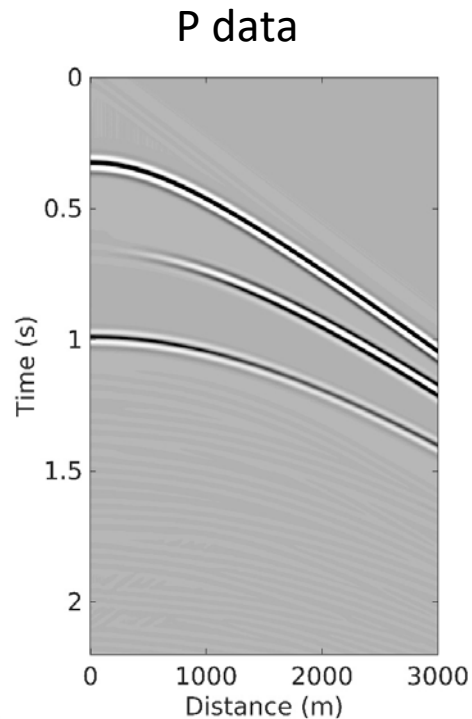
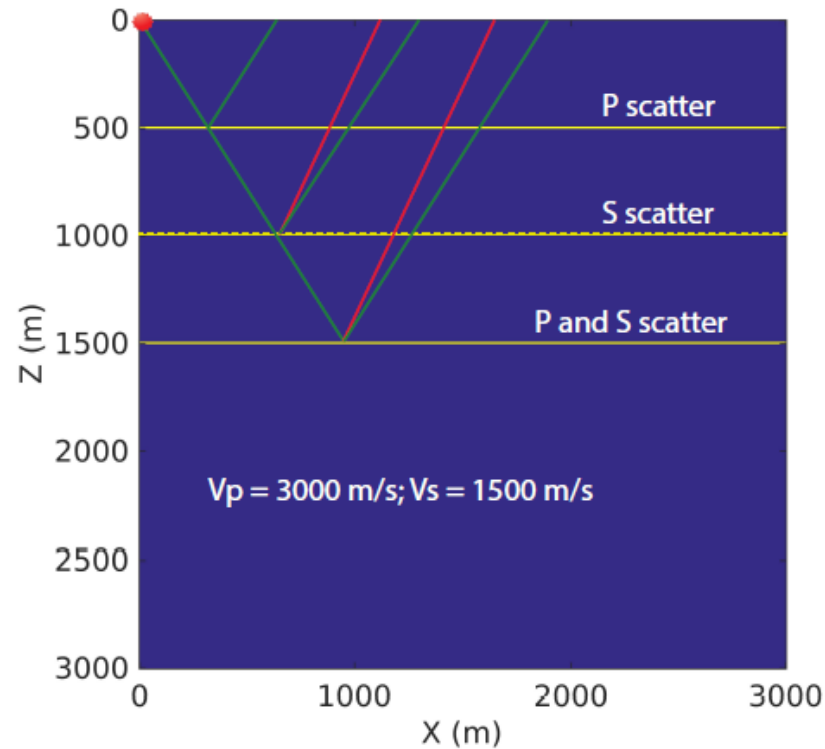


# Insights from the equations

$$\begin{aligned}\ddot{P} - \alpha \nabla^2 P &= P \nabla^2 \alpha + 2 \nabla \alpha \cdot \nabla P - 2 P \nabla^2 \beta - 2 \nabla \beta \cdot \nabla \times \mathbf{S} + \nabla \cdot \mathbf{f} \\ \ddot{\mathbf{S}} - \beta \nabla^2 \mathbf{S} &= \nabla \beta \cdot \nabla \mathbf{S} - (\nabla \beta) \times (\nabla \times \mathbf{S}) + 2 (\nabla \beta) \times (\nabla P) + \nabla \times \mathbf{f}\end{aligned}$$

- ✓ New set of equations: coupled but separated for P- and S-propagations in heterogeneous (Lamé) media (constant density)
- ✓ Wave-medium interactions can be directly interpreted
- ✓ Mode-conversion only happens at S-wave discontinuities!
- ✓ Discontinuities only in  $V_p$  are transparent to S-wave

# Elastic simulations in heterogeneous media



(Removed direct arrival)

# Outline

- Elastic wave equations
  - Revisit of the elastic wave equations
  - A new set of separated P- and S-wave equations
- The elastic imaging condition
  - PP and PS images from inverse problem formulation
- Discussions and conclusions

# Imaging condition

image = source wavefield **meets** scattered wavefield

- ✧ Wavefields only recorded on the boundary
  - ✧ Source: source signature
  - ✧ Scattered: receiver recordings
- ✧ How does the wavefields meet?
  - ✧ P-wave: scalar
  - ✧ S-wave: vector
- ✓ Approximate wavefields by solving wave equations
  - ✓ Source: forward propagation
  - ✓ Scattered: backward propagation
- ✓ Formulate imaging problem as an inverse problem
  - ✓ P-wave: take a gradient
  - ✓ S-wave: take a curl

# Imaging as an inverse problem

- Match the modeled P-wave data with the recorded P-wave data

$$J_p(\alpha, \beta) = \frac{1}{2} \|d_p - d_{p_0}\|_2^2$$

- Conventional PP-image

$$\begin{aligned} \nabla_{\alpha} J_p &= \left( \frac{\partial P}{\partial \alpha} \right)^* \bigg|_{\alpha=\alpha_0, \beta=\beta_0} (d_p - d_{p_0}) \\ &= 4 \boxed{(\nabla^2 P_0)^*} \boxed{(\Pi_p)^{-*} \delta d_p} \end{aligned}$$

$$I_{pp} = \int_t dt \quad \boxed{\text{Forward propagated source P-wavefield}} \otimes \boxed{\text{Backward propagated "scattered" P-wavefield}}$$

Li et. al., Geophysics, 2018

# Imaging as an inverse problem

- Match the modeled S-wave data with the recorded S-wave data

$$J_s(\alpha, \beta) = \frac{1}{2} \|\mathbf{d}_s - \mathbf{d}_{s_0}\|_2^2$$

- Converted PS-image

$$\begin{aligned} \nabla_{\beta} J_s &= \left( \frac{\partial \mathbf{S}}{\partial \beta} \right)^* \bigg|_{\alpha=\alpha_0, \beta=\beta_0} (\mathbf{d}_s - \mathbf{d}_{s_0}) \\ &= -2(\nabla P_0)^* \cdot (\nabla \times \Pi_s^{-*} \delta \mathbf{d}_s) \end{aligned}$$

$$I_{ps} = \int_t dt \operatorname{grad} \left[ \text{Forward propagated source P-wavefield} \right] \otimes \operatorname{curl} \left[ \text{Backward propagated "scattered" S-wavefield} \right]$$

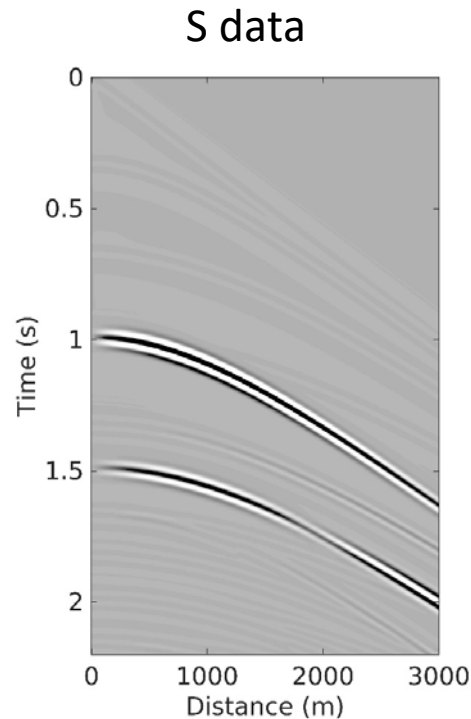
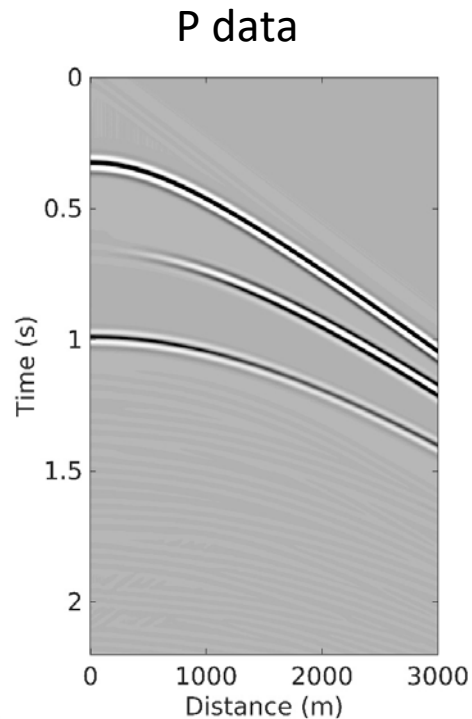
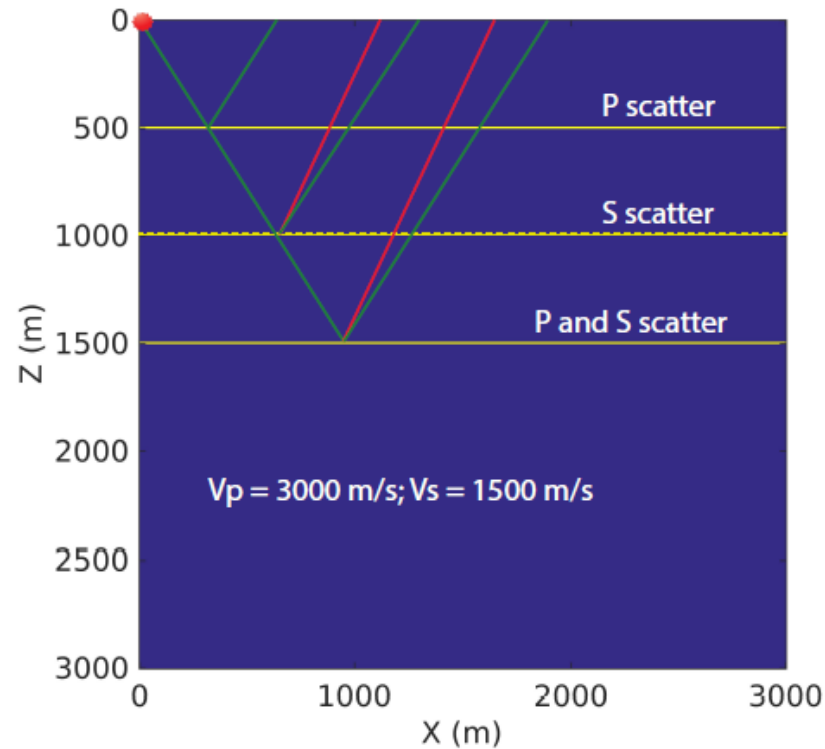


# Elastic imaging using acoustic propagators

$$\begin{aligned}\nabla_{\alpha} J_p &= \left( \frac{\partial P}{\partial \alpha} \right)^* \bigg|_{\alpha=\alpha_0, \beta=\beta_0} (d_p - d_{p_0}) \\ &= 4 (\nabla^2 P_0)^* (\Pi_p)^{-*} \delta d_p\end{aligned}\qquad \begin{aligned}\nabla_{\beta} J_s &= \left( \frac{\partial \mathbf{S}}{\partial \beta} \right)^* \bigg|_{\alpha=\alpha_0, \beta=\beta_0} (\mathbf{d}_s - \mathbf{d}_{s_0}) \\ &= -2 (\nabla P_0)^* \cdot (\nabla \times \Pi_s^{-*} \delta \mathbf{d}_s)\end{aligned}$$

- Migration velocity models are often smooth
- Wave-equations reduce to fully decoupled P- and S-wave equations for their potential fields
- They can be efficiently solved using acoustic propagators

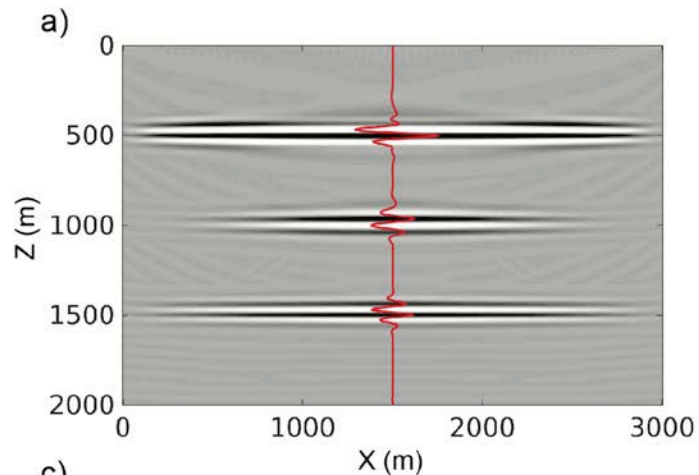
# Elastic simulations in heterogeneous media



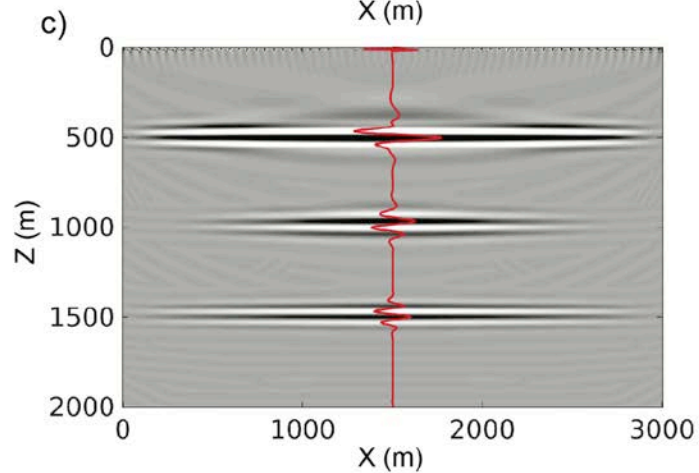
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Using acoustic  
propagators

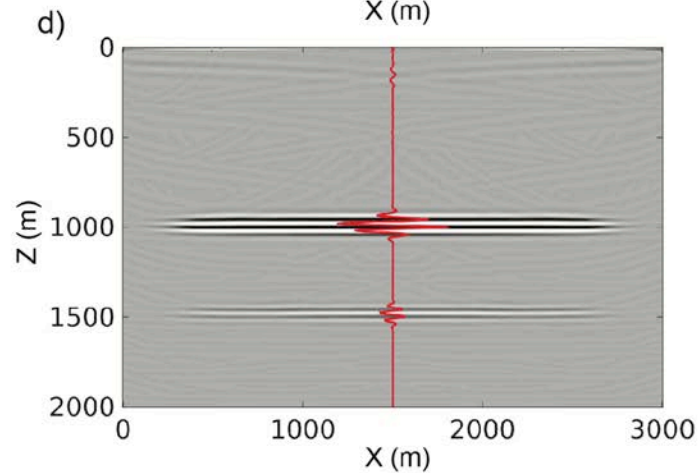
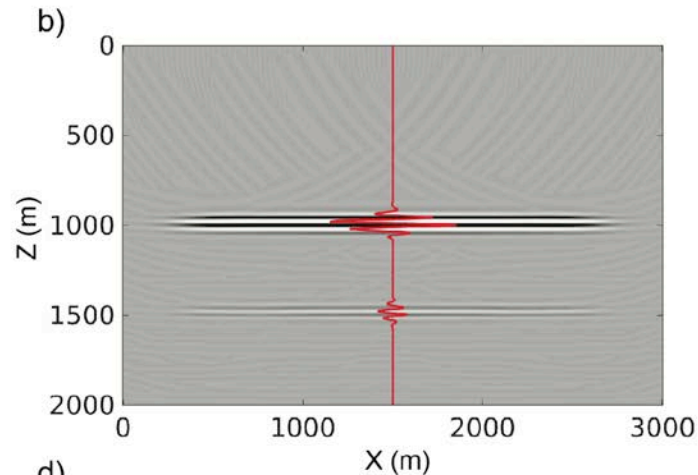
PP Image



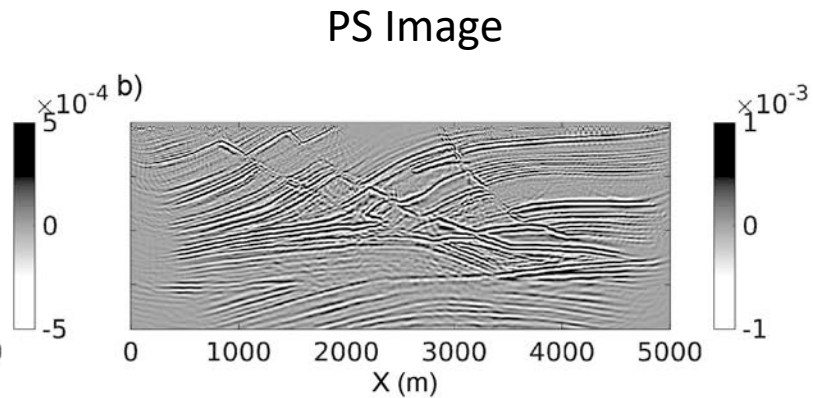
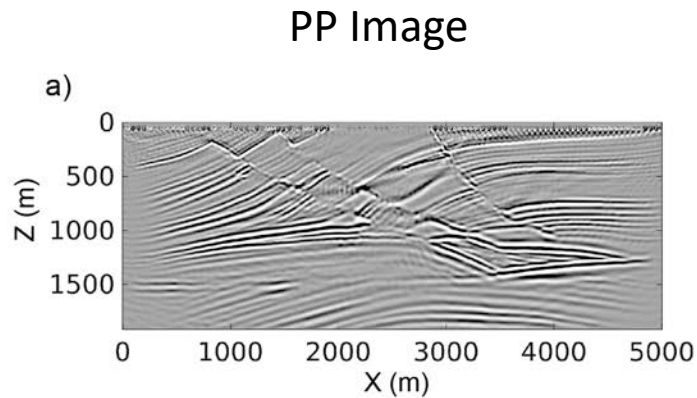
Using elastic  
propagators



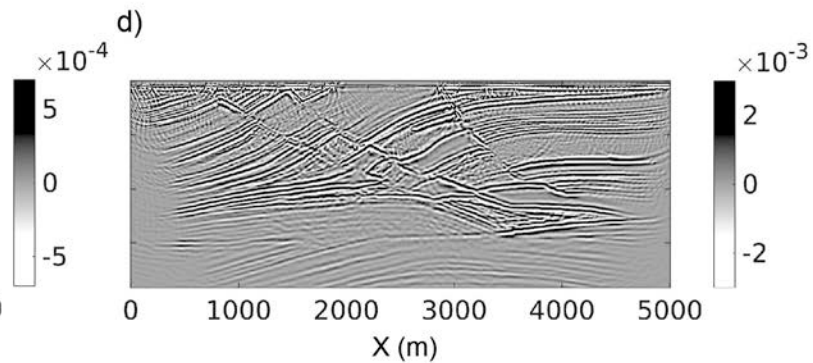
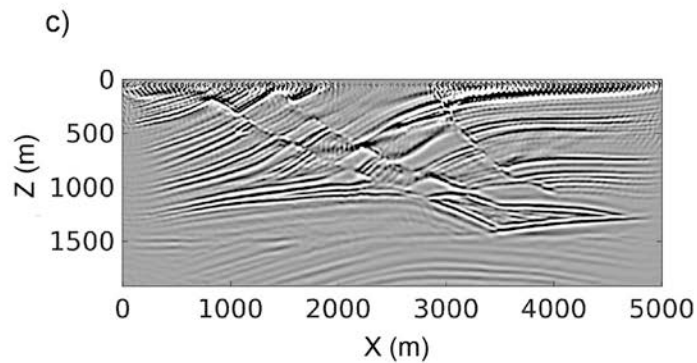
PS Image



Using acoustic  
propagators



Using elastic  
propagators



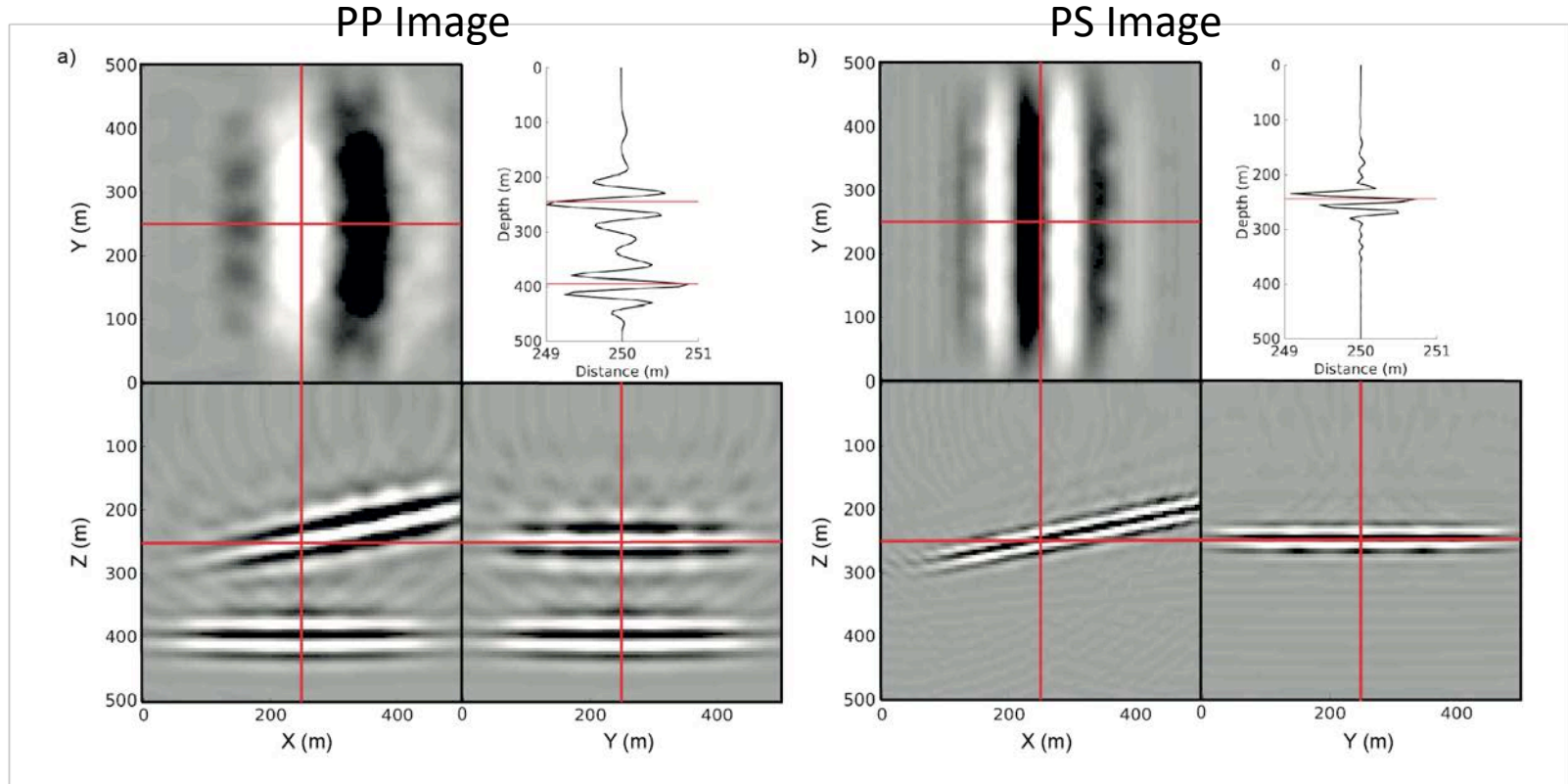
# Comparison of the computational costs

Cost \ Using	Acoustic propagator	Elastic propagators
Memory	$nx*nz*3$	$nx*nz*3*5$
Floating-point operations	$O(nx*nz)$	$O(nx*nz*5)$
# of simulations	2	1

Memory saving up to 80%, run time saving 60%

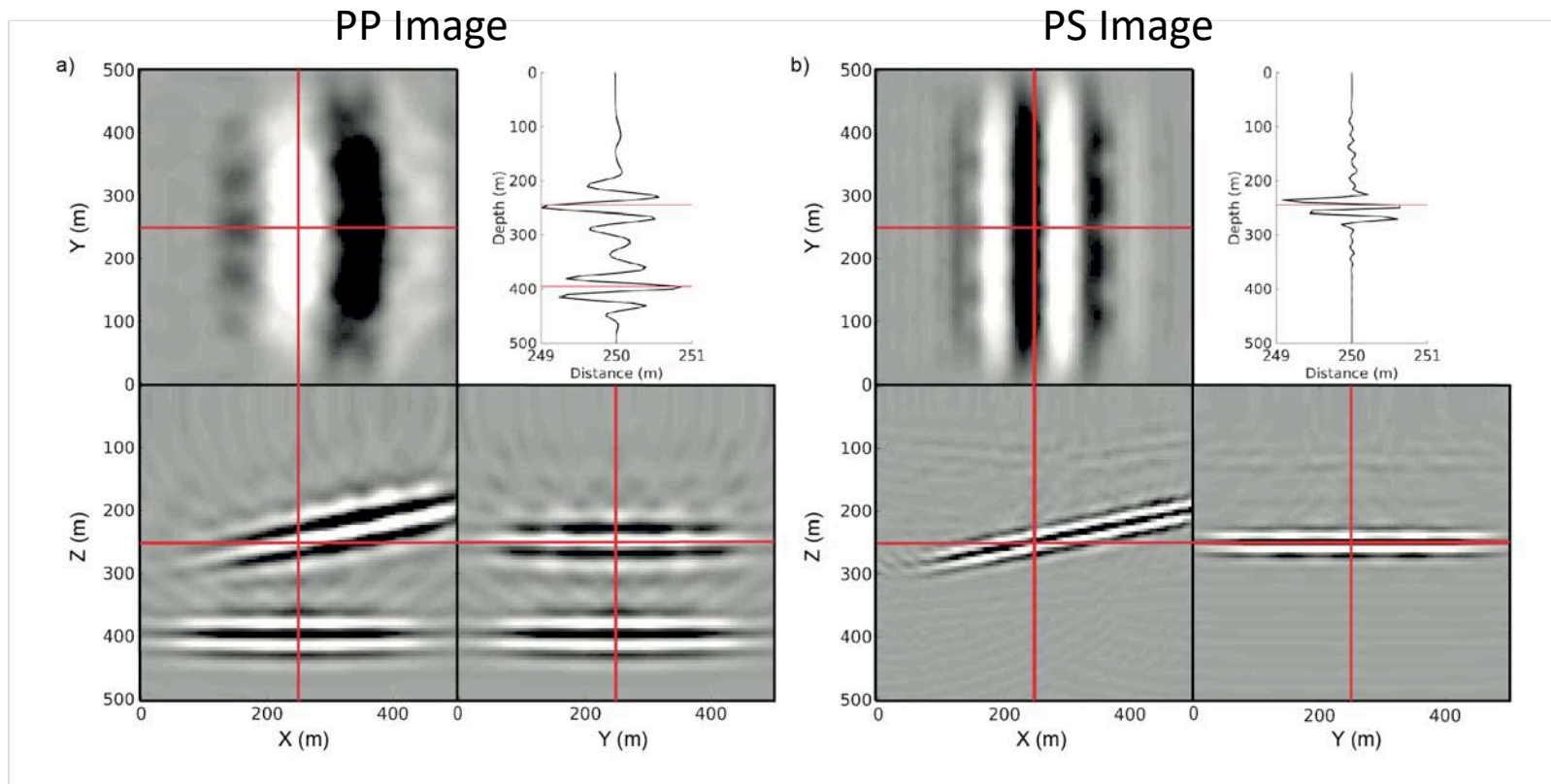
Run time saving up to 80%, memory saving 60%

# Elastic imaging in 3D using acoustic prop.





# Elastic imaging in 3D using elastic prop.



# Comparison of the computational costs

Cost \ Using	Acoustic propagator	Elastic propagators
Memory	$nx*ny*nz*3$	$nx*ny*nz*3*9$
Floating-point operations	$O(nx*ny*nz)$	$O(nx*ny*nz*9)$
# of simulations	4	1

Memory saving up to 88.9%, run time saving 55.6%

Run time saving up to 88.9%, memory saving 55.6%

# Discussions and conclusions

- We derive a new set of coupled, but separated wave equations for P- and S-wave propagation
- This work provides a rigorous theoretical basis for the vector image conditions
- Better interpretation of the PP and PS images based on fundamental wave physics

# Discussions and conclusions

- Advantages of using acoustic propagators for elastic imaging
  - Lower memory and computational cost
  - Free of the artifacts caused by the unphysical wave mode conversion:
    1. Artifacts near the receiver locations
    2. Imprints of S-wave velocity model – “in-situ” mode conversions

# Limitations

- Constant density assumption
  - P- and S-waves are fully coupled at all density discontinuities
  - Images are contaminated with density contrasts
- P- and S-data separation in the recorded data
  - Potential data are needed for this formulation
  - Inverse problem to solve for the separated fields

# Complete set of equations for constant density media

$$\underbrace{\ddot{P} - \alpha \nabla^2 P}_{\text{P propagation}} = \underbrace{P \nabla^2 \alpha + 2 \nabla \alpha \cdot \nabla P}_{\text{P scatter at } V_p \text{ contrast}} + \underbrace{2(\nabla \nabla \beta \cdot \nabla \nabla (\nabla^{-2} P) - P \nabla^2 \beta)}_{\text{P scatter at } V_s \text{ contrast}} - \underbrace{2[\nabla \beta \cdot \nabla \times \mathbf{S} + \nabla \nabla \beta \cdot \nabla \nabla \times (\nabla^{-2} \mathbf{S})]}_{\text{SP mode conversion at } V_s \text{ contrast}} + \underbrace{\nabla \cdot \mathbf{f}}_{\text{source}}, \quad (26)$$

$$\underbrace{\ddot{\mathbf{S}} - \beta \nabla^2 \mathbf{S}}_{\text{S propagation}} = \underbrace{2 \nabla \beta \times \nabla P + 2 \nabla \nabla \beta \star \nabla \nabla (\nabla^{-2} P)}_{\text{PS mode conversion at } V_s \text{ contrast}} + \underbrace{\nabla \beta \cdot \nabla \mathbf{S} - (\nabla \beta) \times (\nabla \times \mathbf{S})}_{\text{S scatter at } V_s \text{ contrast}} - \underbrace{\nabla \nabla \beta \star \{\nabla \nabla \times (\nabla^{-2} \mathbf{S}) + [\nabla \nabla \times (\nabla^{-2} \mathbf{S})]^T\}}_{\text{S scatter at } V_s \text{ contrast}} + \underbrace{\nabla \times \mathbf{f}}_{\text{source}}. \quad (27)$$