Inverse Scattering Problems: To Overcome the Ill-posedness

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# > Introduction

- Inverse Scattering : Analysis and Computation
- Inverse Source Problems
- > On-going Research

# **Scattering and Inverse Scattering**



- Scattering
- Inverse scattering

# **Applications**







Stealth

#### **Geophysical inspection**

Cloaking



**Medical imaging** 



**Super-resolution** 

# **Application: Radon Transform**



**Radon transform (1917)** 



$$Rf(s,\theta) = g(s,\theta) = \int_{\langle x,\theta\rangle=s} f(x)dl = \int_{L} f(x)dt$$

**Inverse Radon transform** 

$$f(x) = \frac{1}{4\pi^2} p.v. \int_{S^1} d\theta \int \frac{\frac{d}{ds}g(s,\theta)ds}{\langle x,\theta \rangle - s}$$



Radon, J.: *Uber die Bestimmung von Funktionen durch ihre Integralwerte Langs Gewisser Mannigfaltigkeiten.* Ber. Saechsische Akad. Wiss. 29, 262 (1917)

# **Application: Radon Transform**

### **1979 Nobel Prize in Medicine Computed Tomography (CT)**





A. Cormack

G. Hounsfield



(From Wikipedia)

### **Integral geometry**

$$T = \int_{\gamma} \frac{1}{c(x)} ds =$$
Travel time



## **Application: Seismic Inversion**



$$\begin{pmatrix} \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \Delta \end{pmatrix} p(x, y, z, t) = 0 \text{Data: } p(x, y, z, t)|_{z=0} \\ \text{To determine } c(x, y, z)?$$

Boundary rigidity problem: Michel 81', Gromov, 83', ... Lens rigidity problem: Largely open!

# **Application: Near-Field Optics**



Diffraction limit:  $\theta_1 = \lambda/2$ 2014 Noble Prize in Chemistry Eric Betzig (92', 93') evanescent  $\theta_1 = \lambda/10$ , (95') PSF  $\Delta_{\min} = \frac{\Delta}{\sqrt{N}} = \frac{1}{\sqrt{N}} \frac{\lambda}{2n \sin \alpha}$ 

Stefan W. Hell (94', 95') STED  $\Delta_{\min} \approx \frac{\lambda}{2n \sin \alpha (\sqrt{1 + I_0 / I_{sat}})}$ 

IOP: Betzig 93'

William E. Moerner experiment



$$\begin{cases} \Delta u + k^2 u = 0\\ \frac{\partial u}{\partial r} - ikr = o(\frac{1}{r}) & \text{as } r \to \infty\\ u + u^i + u^r = 0 \end{cases}$$

Measurements :  $u(x_1, x_2, d)$ ,  $d - \max f < \lambda$ , to determine:  $f(x_1, x_2)$ 

# **Challenges for Inverse Scattering**

- ➢ Ill-posedness
- > Nonlinearity
- > Computation
- > Uncertainty

**To determine**  $\gamma$  **from**  $\Lambda$  (voltage to current map)

$$\nabla \cdot (\gamma(x)\nabla \mathbf{u}) = \mathbf{0}$$
$$\Lambda : \mathbf{u} \Big|_{\Gamma} \to \gamma \frac{\partial u}{\partial \nu} \Big|_{\Gamma}$$



Calderón A. P.: On an inverse boundary value problem, Seminar on Numerical Analysis and its Applications to Continuum Physics (Rio de Janeiro, 1980) Soc. Brasil Mater. 65-73 Rio de Janeiro



- Kohn & Vogelius 1984, CPAM
- Sylvester & Uhlmann 1987, Ann. Math.
- Nachman 1988, Ann. Math.
- Nachman 1995, Ann. Math.
- Pestov & Uhlmann 2005, Ann. Math.
- > Astala & Päivärinta 2006, Ann. Math.
- Kenig, Sjöstrand & Uhlmann 2007, Ann. Math.
- Imanuvilov, Uhlmann & Yamamoto 2010, JAMS
- Uhlmann & Vasy 2016, Invent. Math.

 Unstable! III-posedness...
 Alessandrini 1987, Mandache 2001
 Same ill-posedness for Helmholtz/Maxwell eqns at a fixed frequency! No lucky break...

> Remedy:

Hybrid, Multiple frequency data

### **Inverse Scattering Problems**



Models: Electromagnetic, Optics, Acoustics, Elasticity
 Measurements: boundary

Multiple Frequency

# **Stability for the IP of Wave Equation**

$$\left(\frac{1}{c^2(x)}\partial_t^2 - \Delta\right)u = 0$$
$$\Lambda: u \Big|_{\Gamma} \to \frac{\partial u}{\partial v}\Big|_{\Gamma}, \text{ to find } c(x)$$

#### **Uniqueness:**

Boundary control theory, Belishev & Kurylev 92'

Stability (partial results): Uhlmann, Lassas, Vasy et.al, 98' -- now GO, No Caustics!

For wave equation in time domain, WKB expansion for GO is:

$$u(x,t) = A(x,t)e^{i\omega\varphi(x,t)}, \ x \in \Omega, t \in [0,T],$$

where the frequency  $\omega \gg 1$ , and the phase function  $\varphi(x, t)$  is real. The solution is defined on a ray which follows the characteristic of the eikonal equation.

However, the GO solution blows up at the caustics!





Since the GO solution is local, at the caustics, the rays are intersecting and  $A(x,t) \rightarrow \infty$ 

Gaussian Beam solution is also based on WKB form. Instead of being a local solution, each GB is a global solution.

In particular, near the ray, the phase function admits the following expansion

$$\varphi(x,t) = p(t) \cdot \left(x - x(t)\right) + \frac{1}{2} \left(x - x(t)\right)^T \cdot M(t) \cdot \left(x - x(t)\right) + O(||x - x(t)||^3)$$

where the Hessian matrix M(t) has a positive definite imaginary part.

Since GB solution conserve the energy during the propagation and is linearly independent of the other GB solutions, it will not blow up at the caustics.

## **New Stability Result**

G. Bao and H. Zhang: *Sensitivity analysis of an inverse problem for the wave equations in the presence of caustics*, J. of AMS, 2014

### **New Method:**

Gaussian beam and microlocal analysis

### **Key Ingredients:**

Linearized Hamiltonian with respect to the velocity, Stability analysis of the X-ray transform, Gaussian beams/microlocal analysis First stability result for the inverse problem with caustics ; general well-posedness result

Lens rigidity, first stability result for the nonsimple metric case, ARMA, 2017 **Objective: Stable reconstruction methods:** 

- Ill-posedness
- > Nonlinarity
- Systematic initial guesses

Uncertainty principle, limitation of resolution



- Related direct imaging approaches: Linear sampling, factorization methods, transmission eigenvalues, Colton, Kirsch, Monk, Cakoni, et al
- Many issues not addressed!

### **Inverse Medium Problem**

Maxwell's Eq:  $\nabla \times (\nabla \times E^t) - k^2 (1 + q(x)) E^t = 0$  k: wavenumber q(x) > -1. supported in  $\Omega \subset \mathbb{R}^3$   $E^t = E^i + E$  $E^i$ : incidence

$$\nabla \times (\nabla \times E^i) - k^2 E^i = 0$$



*E*: scattered field

$$\nabla \times (\nabla \times E) - k^2 (1 + q(x))E = k^2 q(x)E^i$$
$$\lim_{r \to \infty} r \left[ \nabla \times E \times \frac{x}{r} - ikE \right] = 0, \quad r = |x|$$

### **Inverse Medium Problem**

### **Abstract Setting:**

$$M(q,k) = \text{Data}(k)$$

#### **Previous work:**

- Dorn, Bertete-Aguirre, Berryman, Papanicolaou (1999), Haddar and Monk (2002), Hohage (2003), Colton and Kress, W. Chew et al (2003, 2004), Sini and Thanh (2012), de Hoop, Scherzer, Qiu (2014)...
- Optimization, initial guess/stability

### **Low Frequency Approximation**

Born

$$\nabla \times (\nabla \times E) - k^2 E = k^2 q(E^i + E)$$
$$\nu \times (\nabla \times E) + ik\nu \times (\nu \times E) = 0$$

#### **Important identity:**

$$\int_{\Omega} q(x) p_1 \cdot p_2 e^{ikx \cdot (n_1 + n_2)} dx = \frac{i}{k} \int_{\Gamma} (v \times E) \cdot ((n_2 + v) \times p_2) e^{ikx \cdot n_2} ds - \int_{\Omega} q(x) p_2 \cdot E e^{ikx \cdot n_2} dx.$$

#### Linearization:

$$\int_{\Omega} q(x)e^{ikx\cdot(n_1+n_2)}dx = \frac{i}{(p_1\cdot p_2)k}\int_{\Gamma} (\nu \times E) \cdot ((n_2+\nu) \times p_2)e^{ikx\cdot n_2}ds$$

$$\hat{q}(\xi) = \int_{\Omega} q(x)e^{ikx \cdot (n_1 + n_2)} dx \qquad \xi = k(n_1 + n_2) \qquad |\xi| \le 2k$$

### **Born + Recursive linearization**

Born  $q_{k_0}$ Do  $i = 1, 2, \dots$  (wave number)  $q_{k}^0 = q_{k_{i-1}}$ Do j = 1, ..., m (incidence)  $\delta q_{j} = \frac{1}{\beta_{k}} DM_{j}^{*}(q_{k_{i}}^{j-1})R_{j}(q_{k_{i}}^{j-1})$  $q_{k}^{j} = q_{k}^{j-1} + \delta q_{j}$ 

End

End

### **Numerical Results**

#### 2D smooth medium





 $\boldsymbol{q}_{k_0}$ 





### **Numerical Results**



 $\eta = 6.6$ 



 $\eta = 4.8$ 





 $\eta = 3.0$ 

 $\eta = 0$ 

# Features on the Algorithm I

- Multiple frequency data is crucial, spectral information
- > Systematic selection of initial guesses
- Convergence, analysis is related to the stability analysis
- > Other types of inverse problems in wave propagation
- > Uncertainty based continuation method

# Features on the Algorithm II

- > Not distorted Born approximation (W. Chew et al)
- Not direct frequency hopping,...
- It is a continuation method!

At each frequency, optimization, regularization are involved, the problem needs not be solved precisely.

# **Uncertainty Based Continuation Method**

Inverse medium scattering:

2-D Full aperture, Chen & Rokhlin 97', Chen 97', B., Chen, & Ma 00'; Limited aperture, B. & Liu 03', B. & Li 07'; Fixed frequency, B. & Li 05', 06', 07'; Convergence, B. & Triki 08'; 3-D, B. & Li 04', 05', 07', 08', 09'

- Inverse source problems: B., Lin, &Triki 10', 11', B., Lu, Rundell, Xu, 15'
- Inverse obstacle scattering:

Coifman, Goldberg, Hrycak, & Rokhlin 99', B., Hou, & Li 07', B., Lin 10', 11', B. & Li 12', B., Li, & Lv 13'

Inverse scattering problems with multi-frequency data B., Li, Lin, Triki, Topical Review, Inverse Problems, 15'

### Model problem: 2D Helmholtz equation

$$\begin{cases} \Delta u + k^2 u = S & \text{in } R^2 \\ \frac{\partial u}{\partial r} - iku = o(r^{-1/2}) & \text{as } r \to +\infty \end{cases}$$

- (1) S compact supported
- (2) Boundary measurements

$$\left\{u_k, \frac{\partial u_k}{\partial n}\right\} \text{ for } k \in [k_{\min}, k_{\max}]$$

**ISP:**  
**From** 
$$\left\{ u_k, \frac{\partial u_k}{\partial n} \right\}$$
 for  $k \in [k_{\min}, k_{\max}]$ 

### to determine S

Bleistein & Cohen 77' He & Romanov 98', Ammari, Gang Bao, Fleming 02' Albanese & Monk 06' Davaney et al 04', 07', ...

Difficulties: Nonuniqueness/ill-posedness at fixed frequency

### Uniqueness: B., Lin, Triki, JDE, 10'

Let  $\{k_j\}_{j=1}^{\infty}$  be a bounded and strictly monotone sequence, then  $\left\{u_k, \frac{\partial u_k}{\partial n}\right\}_{k=1}^{\infty}$  uniquely determine S.

### Stability:

There is a critical number, when the highest frequency exceeds the number. The stability is Lipschitz; otherwise it is logrithmic.

### **Numerical Results for ISP**



### (a) The source function S(x)



(b) Reconstruction of S(x)



### **Inverse Obstacle Problems**





# Rough surface B., P. Li, et al , 13', 14'...



- > Inverse medium problems: Stability
- Inverse problems with uncertainty: Inverse random source problems, rough surface scattering
- Inverse scattering for elastic waves
- Inverse problems via AI

# Stability for Multi-frequency IMP 1-D

1-D Case:

$$\varphi''(x,k) + k^2(1+q(x))\varphi(x,k) = 0, x \in (0,1)$$

IP: Given the reflection coefficients  $k \in [0, k_0]$  to determine q(x)

### Theorem(B., Triki , 17')

Let 
$$q, \tilde{q} \in C_0^m(0,1), ||q||, ||\tilde{q}|| \le M, q, \tilde{q} > q_0 > -1.$$
  
Then for any  $k_0 \ge k_{M,q_0},$ 
$$||q - \tilde{q}||_{L^{\infty}} \le C_{M,q_0} \left( ||d(k) - \tilde{d}(k)||_{L^1(-k_0,k_0)} + \frac{1}{k_0^{m-1}} \right)$$

### Key: Trace formula

# **Inverse Scattering for Elastic Waves**

### Model

 $abla \cdot \sigma(u) + \rho \omega^2 u = 0$   $\sigma(u) = C : \varepsilon(u)$ Kupadze radiation condition Lame parameters:  $\lambda, \mu$ 

### **Shear + pressure**

Forward problem

Solution 
$$u^s = u - u^i$$

Inverse Problem

Reconstruct medium/obstacle/source from boundary measurements



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# **Research for Scattering in Elasticity**

### > Scattering:

Finite element method, transparent BC
B., G. Hu, J. Sun and T. Yin, JMPA, 18'.
Boundary integral equation method
B., L. Xu and T. Yin, JCP, 17'
Time-domain
B., Y. Gao, P. Li, ARMA, 18'

### > Inverse scattering:

Uniqueness/stability: Beretta, Yamamoto, Uhlmann, De Hoop... Object: Recursive linearization algorithm for inverse elastic scattering problems with multi-frequency data. Inverse source problems: stability analysis, B., P. Li et al.

#### Model problem: elasticity equation

$$\begin{cases} \Delta^* u + \omega^2 u = f(x) & \text{ in } B_R, \\ Tu = \mathcal{T} u & \text{ on } \Gamma_R. \end{cases}$$

 $\Delta^* = \mu \Delta + (\lambda + \mu)$  grad div is the Lamé operator

**ISP**. Let f be a complex function with a compact support  $\Omega \subset B_R$ . The ISP is to determine f from the data  $u(x, \omega), x \in \Gamma_R, \omega \in (0, K)$ , where K > 1 is a constant.

# **Stability for ISP**

$$\|u(\cdot,\omega)\|_{\Gamma_R}^2 = \int_{\Gamma_R} \left(|\mathcal{T}u(x,\omega)|^2 + \omega^2 |u(x,\omega)|^2\right) ds_x$$
$$\mathbb{F}_M = \{f \in H^m(\Omega)^3 : \|f\|_{H^m(\Omega)^3} \le M, \text{ supp } f = \Omega \subset B_R\}$$

#### Theorem (B.-Li-Zhao)

Let  $f \in \mathbb{F}_M$  and u be the solution of the scattering problem corresponding to f. Then

$$\|f\|_{L^{2}(B_{R})^{3}}^{2} \lesssim \epsilon^{2} + \frac{M^{2}}{\left(\frac{K^{\frac{2}{3}}|\ln\epsilon|^{\frac{1}{4}}}{(R+1)(6m-15)^{3}}\right)^{2m-5}},$$

where

$$\epsilon = \left(\int_0^K \omega^2 \|u(\cdot,\omega)\|_{\Gamma_R}^2 \mathrm{d}\omega\right)^{\frac{1}{2}}.$$

- Inverse scattering problems are an exciting and fast growing area of mathematics driven by interdisciplinary applications.
- Many basic math questions to be addressed; effective computational methods are in demand.
- Interactions of crossed-disciplinary efforts, Prof. Bolomey...
- New problems are emerging; AI, deep learning, big data,...

