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Microwave Near-field Imaging in Real Time

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APPLICATIONS OF MICROWAVE NEAR-FIELD IMAGING

- penetration into optically obscured objects (clothing, walls, luggage, living tissue...)
 - \succ the lower the frequency the better the penetration
 - \succ frequency bands from 500 MHz into the mm-wave bands (≤300 GHz)
- compact relatively cheap electronics esp. in the low-GHz range
- diverse suite of image reconstruction methods

VARIOUS APPLICATIONS



APPLICATIONS: LUGGAGE INSPECTION, NDT

[Ghasr et al., "Wideband microwave camera for real-time 3-D imaging," IEEE Trans. AP, 2017]



20 GHz to 30 GHz frequency range

Prof. Zoughi's team at Missouri University of Science & Technology



[https://youtu.be/RE-PPXmtTeA]

Fig. 15. Example of video camera utility for imaging a box cutter and a pair of scissors inside a laptop bag. (a) Picture of laptop bag in front of the camera aperture with inset showing the objects inside the bag. (b) 3-D view. (c) 2-D image slice focused on the box cutter. (d) 2-D image slice focused on the pair of scissors.

APPLICATIONS: WHOLE BODY SCANNERS



APPLICATIONS: MEDICAL IMAGING

[Song *et al.*, "Detectability of breast tumor by a hand-held impulse-radar detector: performance evaluation and pilot clinical study," *Nature Sci. Reports* 2017] (a)

Prof. Kikkawa's team at Hiroshima University, Japan



Figure 3. Dome antenna array design. (a) The top view of the antenna in *x*-*y* plane. (b) The side view of the antenna in *x*-*z* plane. (c) Top view photograph. (d) Bottom view photograph.



MICROWAVE NEAR-FIELD IMAGING: COMMERCIAL GROWTH

• mm-wave whole-body imagers for airport security inspection (> 30 GHz)

 through-wall and through-floor infrastructure inspection for contractors and home inspectors (UWB, 3 GHz to 8 GHz)

• numerous underground radar applications: detection of pipes, cables, tunnels, etc. (< 3 GHz)



https://www.youtube.com/watch?v=_YZEa1hiGO0]



OUTLINE

> specifics of near-field microwave imaging in

- data acquisition
- forward models
- inversion strategies (real-time)



DATA ACQUISITION: SCANNING

Principle: *N***-D imaging** (*N***=2,3**) **needs** *abundant* **and** *diverse N***-D data sets**

- ➤ spatial scans 1D (linear) or 2D (surface)
 - illuminate target from various angles
 - collect scattered signals at various angles/distances
 - acquisition surfaces planar, cylindrical, spherical
 - mechanical vs. electronic scanning

	mechanical scan	electronic switching
speed	low	HIGH
complexity	LOW	high
flexibility in adjusting scan parameters	GREAT	limited

- > frequency/temporal sweeps
- > polarization diversity



DATA ACQUISITION: SPATIAL SAMPLING

- ➤ spatial data diversity
 - each sample must add independent information improves uniqueness
 - linearly dependent data may lead to ill-posed inversion problems
 - over-sampling: pros and cons
 - staying below but close to the *maximum* spatial sampling step ensures diversity

$$\Delta \zeta \leq \Delta \zeta_{\max} \approx \frac{\lambda_{\min}}{4\sin\alpha}, \ \zeta \equiv x, y$$

• *effective near-field wavelengths* may be shorter than $\lambda = 2\pi / k_{b}$

$$\lambda_{\rm eff,min} = \frac{2\pi}{k_{x,y}^{\rm max}}$$

$$\tilde{S}(k_x, k_y) = \mathcal{F}_{2D}\left\{S(x, y)\right\}$$



DATA ACQUISITION: FREQUENCY SAMPLING

- Frequency data diversity in frequency-sweep measurements
 - stay below but close to the *maximum frequency sampling step*
 - ensures that back-scattered signals (after IFFT) from all targets at distances $\leq R_{\text{max}}$ do not overlap

$$\Delta f \le \Delta f_{\max} = \frac{1}{2T_{\max}} \approx \frac{v_{\rm b}}{4R_{\max}}$$



- temporal data diversity in time-domain measurements
 - > stay below but close to the *maximum time sampling step*
 - > ensures that all frequency components of the pulsed signals are fully used (Nyquist)

$$\Delta t \le \Delta t_{\max} \approx \frac{T_{\min}}{2} = \frac{1}{2f_{\max}}$$

FORWARD MODEL OF SCATTERING: S-PARAMETER DATA EQUATION

[Nikolova et al., APS-URSI 2016][Beaverstone et al., IEEE Trans. MTT, 2017]

• scattering from penetrable objects (isotropic scatterer)



FORWARD MODEL: BORN'S LINEARIZING APPROXIMATION



- total field $\mathbf{E}_{k}^{\text{tot}}(\mathbf{r}';\Delta\varepsilon(\mathbf{r}'))$ is generally unknown and depends on contrast: *data equation is nonlinear in the unknown contrast*
- Born's approximation of the total internal field linearizes the data equation:

$$\mathbf{\underline{E}}_{k}^{\text{tot}}(\mathbf{r}';\Delta\varepsilon(\mathbf{r}')) \approx \mathbf{\underline{E}}_{k}^{\text{inc}}(\mathbf{r}') \qquad \Rightarrow S_{ik}^{\text{sc}} \approx \iiint_{V_{s}} \Delta\varepsilon_{r}(\mathbf{r}') \underbrace{\overline{\mathbf{E}}_{i}^{\text{inc}}(\mathbf{r}') \cdot \overline{\mathbf{E}}_{k}^{\text{inc}}(\mathbf{r}')}_{\text{resolvent kernel}} d\mathbf{r}'$$

• main challenge of near-field imaging with linear inversion methods: *incident-field distributions in antennas' near zone are difficult to model*

(assumed known)

FORWARD MODEL: FAR-ZONE vs. NEAR-ZONE ANTENNA FIELDS

any one of the conditions below implies near-field (short-range) imaging

- object is in the near-field region of antennas
- distance from object to antennas ≤ object's size
- distance from object to $r \leq \lambda$ antennas \leq wavelength





Goran M Djuknic - US Patent 6657596, commons.wikimedia.org/w/index.php?curid=20417988

• implication: resolvent kernel depends on incident fields which do not conform to analytical free-space far-zone propagation models, e.g.,

$$\mathbf{E}^{\mathrm{inc}}(\mathbf{r}') \sim \hat{\mathbf{p}}G(\theta, \varphi) \frac{e^{-ik_{\mathrm{b}}r}}{r} \longleftarrow not \ valid$$

LINEARIZED FORWARD MODEL: TIME DOMAIN

• Born's linearizing approximation is applied in the same way

 $\frac{\text{frequency-domain (S-parameters)}}{S_{ik}^{\text{sc}} \approx \iiint_{V_{s}} \Delta \varepsilon_{r}(\mathbf{r}') \overline{\mathbf{E}}_{i}^{\text{inc}}(\mathbf{r}') \cdot \overline{\mathbf{E}}_{k}^{\text{inc}}(\mathbf{r}') d\mathbf{r}'}$



[Nikolova, Introduction to Microwave Imaging, 2017]

LINEARIZED FORWARD MODEL: BORN vs. RYTOV DATA APPROXIMATION

• Born scattered-field data approximation

$$S_{ik}^{\rm sc}(\mathbf{r} \in S_{\rm a}) \approx \underbrace{\left[S_{ik}^{\rm tot} - S_{ik}^{\rm inc}\right]_{(\mathbf{r} \in S_{\rm a})}}_{\text{data calibration step}} \approx \iiint_{V_{\rm s}} \Delta \varepsilon_r(\mathbf{r}') \overline{\mathbf{E}}_i^{\rm inc}(\mathbf{r}') \cdot \overline{\mathbf{E}}_k^{\rm inc}(\mathbf{r}') d\mathbf{r}'$$

• Rytov scattered-field data approximation

$$S_{ik}^{\text{sc}}(\mathbf{r} \in S_{a}) \approx S_{ik}^{\text{inc}} \ln \left(\frac{S_{ik}^{\text{tot}}}{S_{ik}^{\text{inc}}} \right)_{\mathbf{r} \in S_{a}} \approx \iiint_{V_{s}} \Delta \varepsilon_{r}(\mathbf{r}') \overline{\mathbf{E}}_{i}^{\text{inc}}(\mathbf{r}') \cdot \overline{\mathbf{E}}_{k}^{\text{inc}}(\mathbf{r}') d\mathbf{r}'$$

data calibration step

[Tajik et al., JPIER- B, 2017][Shumakov et al., Trans. MTT, 2018]

FORWARD MODEL: APPROXIMATIONS

> total internal field approximation (Born)

• limitations on both size and contrast of the scatterer

 $a^{2}\left|k_{\rm s}^{2}(\mathbf{r})-k_{\rm b}^{2}\right|\ll 1, \, \mathbf{r}\in V_{\rm s}$

[Nikolova, Introduction to Microwave Imaging, 2017]

- if OUT violates limit: image contains artifacts which reflect differences between $\mathbf{E}_{Tx}^{inc}(\mathbf{r}')$ and $\mathbf{E}_{Tx}^{tot}(\mathbf{r}')$ rather than contrast
- expect trouble in areas of strong multiple scattering & mutual coupling

➤ data approximation

- Born limitation on both size and contrast: $2a |k_s(\mathbf{r}) k_b| < \pi, \mathbf{r} \in S_a$
- Born neglects multiple scattering & mutual coupling between antennas and OUT
- Rytov limitation on contrast only:







LINEAR INVERSION METHODS: THE ENGINES OF REAL-TIME IMAGING

principle: linearize forward model & solve resulting linear system of equations





principle: solve nonlinear forward-model equations for BOTH contrast and total field using nonlinear optimization and/or iterative methods

$$S_{ik}^{\text{sc}} \approx \iiint_{V_{s}} \Delta \varepsilon_{r}(\mathbf{r}') \overline{\mathbf{E}}_{i}^{\text{inc}}(\mathbf{r}') \cdot \overline{\mathbf{E}}_{k}^{\text{tot}} (\mathbf{r}'; \Delta \varepsilon_{r}(\mathbf{r}')) d\mathbf{r}' \qquad \text{Maxwell's equations}$$

data contrast Green's function total internal field
(known) (unknown) (assumed known) (unknown)

RECONSTRUCTION WITH FREQUENCY-SWEEP DATA: HOLOGRAPHY

holography refers to reconstruction methods that use both the magnitude and phase of the scattered waves recorded at a surface to produce a 3D image in a <u>single</u> inversion step



type of response	number of values $S_{ik}(x',y')$	
co-pol X-X	$4 \mathbf{x} N_{\omega}$	
co-pol Y-Y	$4 \mathbf{x} N_{\omega}$	
cross-pol X-Y	$4 \mathbf{x} N_{\omega}$	
cross-pol Y-X	$4 \mathbf{x} N_{\omega}$	
TOTAL	$16 \mathbf{x} N_{\omega} = N_T$	
T number of responses acquired at each position		

HOLOGRAPHY: RESOLVENT KERNEL

$$S_{\xi}^{\rm sc}(x, y, \overline{z}; \omega) \approx \iiint_{V_{\rm s}} \Delta \varepsilon_r(x', y', z') \left[\overline{\mathbf{E}}_{\xi, \mathrm{Rx}}^{\rm inc} \cdot \overline{\mathbf{E}}_{\xi, \mathrm{Tx}}^{\rm inc} \right]_{(x', y', z'; x, y, \overline{z}; \omega)} dx' dy' dz', \ \xi = 1, \dots, N_T$$

approximate (Born) resolvent kernel $\mathcal{K}_{\xi}(\mathbf{r}'; \mathbf{r}; \omega)$

• assume kernel is translationally invariant in *x* and *y* (background is uniform or layered)

Let
$$\mathcal{K}_{0,\xi}(x', y'; z'; \omega)$$
 antennas at
 $\equiv \mathcal{K}_{\xi}(x', y', z'; 0, 0, \overline{z}; \omega)$
Then $\mathcal{K}_{\xi}(x', y', z'; x, y, \overline{z}; \omega) =$
 $\mathcal{K}_{0,\xi}(x' - x, y' - y; z'; \omega)$



 $\Rightarrow S_{\xi}^{\rm sc}(x, y, \overline{z}; \omega) \approx \iiint_{V_{\rm s}} \Delta \varepsilon_r(x', y', z') \mathcal{K}_{0,\xi}(x' - x, y' - y; z'; \omega) dx' dy' dz'$



• **examples:** analytical kernels used with far-zone reflection data $(\overline{\mathbf{E}}_{\xi,\mathrm{Rx}}^{\mathrm{inc}} = \overline{\mathbf{E}}_{\xi,\mathrm{Tx}}^{\mathrm{inc}})$

plane waves:
$$\mathcal{K}_{\xi}(x', y', z'; x, y, z_{\text{Rx}}; \omega) = \overline{\mathbf{E}}_{\text{Rx}}^{\text{inc}} \cdot \overline{\mathbf{E}}_{\text{Rx}}^{\text{inc}} \sim e^{-i2k_{\text{b}}\sqrt{(x-x')^{2}+(y-y')^{2}+(z-z')^{2}}}$$

spherical waves: $\mathcal{K}_{\xi}(x', y', z'; x, y, z_{\text{Rx}}; \omega) \sim \frac{e^{-i2k_{\text{b}}\sqrt{(x-x')^{2}+(y-y')^{2}+(z-z')^{2}}}{(x-x')^{2}+(y-y')^{2}+(z-z')^{2}}$
cylindrical waves: $H_{0}^{(2)}(2k\rho), \ \rho = \sqrt{(x-x')^{2}+(y-y')^{2}}, \ z = z' = const$

- far-field analytical kernels do not work well with near-field data
- kernels computed from simulated incident-field distributions suffer from modeling errors [Amineh *et al*, *Trans. AP*, 2011]
- **near-field kernels** are best determined through measuring the system PSF [Savelyev&Yarovoy, *EuRAD* 2012][Amineh *et al.*, *IEEE Trans. Instr.*&Meas., 2015]

HOLOGRAPHY: MEASURING THE POINT-SPREAD FUNCTION (PSF)

• PSF-based kernels enable **quantitative imaging in real time**



• relating PSF to kernel





Then
$$\mathcal{K}_{0,\xi}(x, y; z'; \omega) = \text{PSF}_{0z',\xi}(-x, -y, \overline{z}; \omega) / (\Delta \varepsilon_{r,\text{sp}} \Omega_{\text{sp}})$$



scattering probe

DATA EQUATION OF HOLOGRAPHY IN TERMS OF PSF

• in real space

$$S_{\xi}^{\rm sc}(x,y,\overline{z};\omega) \approx \frac{1}{\Delta \varepsilon_{r,\rm sp} \Omega_{\rm sp}} \int_{z'} \iint_{y'x'} \Delta \varepsilon_{r}(x',y',z') \cdot \mathrm{PSF}_{0z',\xi}(x-x',y-y';\omega) dx' dy' dz'$$
2D convolution

• in Fourier (or *k*) space

$$\tilde{S}_{\xi}^{\rm sc}(k_x,k_y;\overline{z};\omega) \approx \frac{\Delta x' \Delta y'}{\Delta \varepsilon_{r,\rm sp} \Omega_{\rm sp}} \int_{z'} \underbrace{\tilde{F}(k_x,k_y;z')}_{\operatorname{FT_{2D}}\{\Delta \varepsilon_r(x',y',z')\}} \cdot \widetilde{\operatorname{PSF}}_{0z',\xi}(k_x,k_y;\omega) \, dz'$$

• system of equations to solve *at each spectral position* $\kappa = (k_x, k_y)$

ADVANTAGES OF SOLVING IN FOURIER SPACE: DIVIDE AND CONQUER

$$\widetilde{S}_{\xi}^{(m)}(\mathbf{\kappa}_{ij}) \approx \sum_{n=1}^{N_z} \widetilde{f}(\mathbf{\kappa}_{ij}; z'_n) \widetilde{\mathrm{PSF}}_{0z',\xi}^{(m)}(\mathbf{\kappa}_{ij}) \qquad \begin{array}{l} m = 1, \dots, N_{\omega} \\ \xi = 1, \dots, N_T \end{array}$$

$$\Rightarrow \mathbf{A}(\mathbf{\kappa}_{ij})_{[N_T N_\omega \times N_z]} \cdot \mathbf{f}(\mathbf{\kappa}_{ij})_{[N_z \times 1]} = \mathbf{d}(\mathbf{\kappa}_{ij})_{[N_T N_\omega \times 1]}$$

$$\mathbf{\kappa}_{ij} = (i\Delta k_x, j\Delta k_y)$$

$$i = 1, \dots, N_x; \ j = 1, \dots, N_y$$

- we solve $(N_x \cdot N_y)$ such systems (on the order of 10⁴ to 10⁵)
- the size of each system is small: $N_T N_{\omega} \times N_z$ (e.g. 60×5)
- typical execution times: 2 to 3 seconds on a laptop using Matlab
- solution is orders of magnitude faster than solving in real space:

$$N_D \times N_v$$
 where $N_v = N_x N_y N_z \sim 10^6$ to 10^7
 $N_D = N_x N_y N_\omega N_T \sim 10^7$ to 10^8

FINAL STEP: BACK TO REAL SPACE

• at each plane along range (z' = const)

$$\Delta \varepsilon_r(x', y', z'_n) = \frac{\Delta \varepsilon_{r, \text{sp}} \Omega_{\text{sp}}}{\Delta x' \Delta y' \Delta z'_n} \mathcal{F}_{2D}^{-1} \left\{ \tilde{f}(\mathbf{\kappa}; z'_n) \right\}, n = 1, \dots, N_z$$

$$\varepsilon_r(x', y', z'_n) = \varepsilon_{r,b} + \Delta \varepsilon_r(x', y', z'_n)$$

EXAMPLE: METALLIC TARGETS IN AIR

[Amineh et al., Trans. Instr.&Meas., 2015]







f(GHz)	λ (mm)	$D_{\rm far}$ (mm)
3	100	12.5
8.2	37	34
20	15	83



[photo credit: Justin McCombe]

METALLIC TARGETS IN AIR – RESULTS WITH SIMULATED KERNELS

[Amineh et al., Trans. Instr.&Meas., 2015]





27

50

50

METALLIC TARGETS IN AIR – RESULTS WITH MEASURED KERNELS (PSF)

[Amineh et al., Trans. Instr.&Meas., 2015]



EXAMPLE: IMAGING TISSUE

[Tajik et al., JPIER-B 2017]



Tissue	Color Highlight	Relative Permittivity Averaged over 3 to 8 GHz
Chicken Wing		NA
Bone		21 — 10i
Skin		13 — 6i
Muscle		45 — 23i
Peanut Butter & Jam		7 — 3i
Carbon Rubber		10 — 3i
$\Delta x = \Delta y = 3 \text{ mm}$ $\Delta f = 100 \text{ MHz}$ $f \in [3, 8] \text{ GHz}$		

a layer in the OUT

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a layer in the OUT

EXAMPLE: ACQUIRING THE PSF

[Tajik et al., JPIER-B 2017]

[photo credit: Daniel Tajik]



calibration object with small scattering probe at center: $\varepsilon_{r,SP} \approx 18-i0$, radius 5 mm, height 10 mm

EXAMPLE: IMAGING TISSUE

[Tajik et al., JPIER-B 2017][EuCAP 2018]



QUALITATIVE IMAGING WITH SENSITIVITY MAPS

reconstruction formula

[Tu et al., Inv. Problems, 2015]

using simulated

$$D(\mathbf{r}') = \sum_{\xi=1}^{N_T} \int_{\omega} \iint_{\mathbf{r} \in S_a} \left[S_{\xi}^{\text{inc}} - \overline{S}_{\xi}^{\text{tot}} \right]_{(\mathbf{r},\omega)} \cdot \left[\frac{\partial S_{\xi}^{\text{inc}}(\mathbf{r},\omega)}{\partial \varepsilon(\mathbf{r}')} \right]^* d\mathbf{r} \, d\omega \qquad \text{adjoint sensitivity} \\ \frac{\partial \sigma (\mathbf{r}')}{\partial \varepsilon(\mathbf{r}')} = \int_{\varepsilon}^{\infty} \frac{\partial \sigma (\mathbf{r}')}{\partial \varepsilon(\mathbf{r}')} \int_{\varepsilon}^{\varepsilon} \frac{\partial \sigma$$

sensitivity map: 3D image of Fréchet derivative of ℓ_2 norm of the differences of all total and incident responses λ7

$$D(\mathbf{r}') = \mathcal{J}\left\{F[\varepsilon(\mathbf{r}')]\right\} = 0.5 \sum_{\xi=1}^{N_T} \int_{\omega} \iint_{\mathbf{r} \in S_a} \left\|S_{\xi}^{\text{tot}}(\mathbf{r}, \omega) - S_{\xi}^{\text{inc}}[\mathbf{r}, \omega; \varepsilon(\mathbf{r}')]\right\|_2^2 d\mathbf{r} d\omega$$

$$\operatorname{Re}\left\{D^{(m)}(\mathbf{r}')\right\} = \frac{dF^{(m)}}{d\varepsilon'(\mathbf{r}')} \longrightarrow \text{ indicates where contrast in } \varepsilon' \text{ exists}$$

$$\operatorname{Im}\left\{D^{(m)}(\mathbf{r}')\right\} = -\frac{dF^{(m)}}{d\varepsilon''(\mathbf{r}')} \longrightarrow \text{ indicates where contrast in } \varepsilon'' \text{ exists}$$

 \rightarrow indicates where contrast in ε' exists

$$\rightarrow$$
 indicates where contrast in ε'' exists

SENSITIVITY MAPS USING MEASURED PSFs: THE NEAR-FIELD CASE

• from simulated incident fields to measured PSFs

$$\frac{\partial S_{\xi}^{\text{inc}}(\mathbf{r},\omega)}{\partial \varepsilon(\mathbf{r}')} \approx \frac{\Delta S_{\xi}^{\text{inc}}(\mathbf{r},\omega)}{\Delta \varepsilon(\mathbf{r}')} \approx \frac{\overline{S}_{\xi,\text{sp}}^{\text{tot}}(\mathbf{r},\omega) - S_{\xi}^{\text{inc}}(\mathbf{r},\omega)}{\Delta \varepsilon_{\text{sp}}(\mathbf{r}')} = \frac{\text{PSF}_{\xi}^{\text{sc}}(\mathbf{r},\omega;\mathbf{r}')}{\Delta \varepsilon_{\text{sp}}}$$

$$\frac{\text{PSF}_{\xi}^{\text{sc}}(x, y, \overline{z}, \omega; x', y', \overline{z}') = \text{PSF}_{\xi,0}^{\text{sc}}(x - x', y - y', \overline{z}, \omega; \overline{z}')$$

$$\frac{\text{uniform background along } x \text{ and } y}{\text{top scanning plane}}$$

$$\frac{\text{antennas}}{\text{finish}(x_{Nx}, y_{Ny})}$$

SENSITIVITY MAPS USING MEASURED PSFs: THE NEAR-FIELD CASE – 2

[Nikolova, Introduction to Microwave Imaging, 2017][Shumakov et al., IEEE Trans. MTT, 2018]

• sensitivity reconstruction formula with PSFs: *scattered-power maps (SPM)*

$$-\Delta \varepsilon_{\rm sp} \cdot D(\mathbf{r}') = M(\mathbf{r}') = \sum_{\xi=1}^{N_T} \int_{\omega} \iint_{\mathbf{r} \in S_a} S_{\xi}^{\rm sc}(\mathbf{r}, \omega) \cdot \left[\text{PSF}_{\xi,0}^{\rm sc}(\mathbf{r}, \omega; \mathbf{r}') \right]^* d\mathbf{r} \, d\omega$$

• SPM reconstruction formula with planar scanning

$$M^{(m)}(x', y', \overline{z}') = \sum_{\xi=1}^{N_T} \int_{\omega} \underbrace{\iint_{S_a} \left[S_{\xi}^{\text{sc}}(x, y, \overline{z}, \omega) \right] \cdot \left[\text{PSF}_{\xi}^{\text{sc}}\left(x - x', y - y', \overline{z}, \omega; \overline{z}'\right) \right]^* dxdy \, d\omega}_{\text{cross-correlation of response } S_{\xi}^{\text{sc}} \text{ and } \text{PSF}_{\xi}^{\text{sc}} \text{ in } (x, y)}$$

- reconstruction is practically instantaneous no systems of equations are solved
- effort shifted to near-field system calibration measuring PSFs or simulating incident fields to compute response sensitivity $\partial S_{\xi}^{inc}(\mathbf{r},\omega) / \partial \varepsilon(\mathbf{r}')$

SCATTERED-POWER MAPS: SIMULATION EXAMPLE (Altair FEKO)





sample PSF: S_{11} at 4 GHz MAG/PHASE



 $f_{min} = 3 \text{ GHz}$ $f_{max} = 16 \text{ GHz}$ $\Delta f = 1 \text{ GHz}$

SCATTERED-POWER MAPS: SIMULATION EXAMPLE (F SHAPE)



• blurring typical for cross-correlation methods – diffraction limit, limited number of responses



QUANTITATIVE REAL-TIME IMAGING WITH SPM

[Tu et al., Inv. Problems, 2015][Shumakov et al., IEEE Trans. MTT, 2018]

- requires <u>measured PSFs</u> they scale accurately with the scattering-probe contrast
- reconstruction solves the linear problem



• real-time solution via inversion in Fourier space (similar to holography)

$$M(x', y', z') = \frac{1}{\Delta \varepsilon_{r,sp} \Omega_{sp}} \int_{z''} \int_{y'' x''} \Delta \varepsilon_{r}(x'', y'', z'') \cdot M_{sp@(0,0,z'')}(x' - x'', y' - y'', z') dx'' dy'' dz''$$

convolution in (x, y)

QUANTITATIVE REAL-TIME IMAGING WITH SPM: FOURIER SPACE SOLUTION

$$\tilde{M}(k_x,k_y,z_p) = \frac{\Omega_v}{\Delta \varepsilon_{\mathrm{r,sp}} \Omega_{\mathrm{sp}}} \sum_{q=1}^{N_z} \tilde{f}(k_x,k_y,z_q) \cdot \tilde{M}_{\mathrm{sp}@(0,0,z_q)}(k_x,k_y,z_p), \ p = 1,\dots,N_z$$

• small square system of equations to solve *at each spectral position* $\kappa = (k_x, k_y)$

$$\begin{split} \mathbf{M}_{(\mathbf{\kappa})} \mathbf{x}_{(\mathbf{\kappa})} &= \mathbf{m}_{(\mathbf{\kappa})} \\ \mathbf{x}_{(\mathbf{\kappa})} &= \begin{bmatrix} \tilde{f}(\mathbf{\kappa}, z_1) \cdots \tilde{f}(\mathbf{\kappa}, z_{N_z}) \end{bmatrix}^T \\ \mathbf{m}_{(\mathbf{\kappa})} &= \begin{bmatrix} \tilde{M}(\mathbf{\kappa}, z_1) \cdots \tilde{M}(\mathbf{\kappa}, z_{N_z}) \end{bmatrix}^T \mathbf{M}_{(\mathbf{\kappa})} = \begin{bmatrix} \tilde{M}_{\mathrm{sp}@(0,0,z_1)}(\mathbf{\kappa}, z_1) & \cdots & \tilde{M}_{\mathrm{sp}@(0,0,z_{N_z})}(\mathbf{\kappa}, z_1) \\ \vdots & \ddots & \vdots \\ \tilde{M}_{\mathrm{sp}@(0,0,z_1)}(\mathbf{\kappa}, z_{N_z}) & \cdots & \tilde{M}_{\mathrm{sp}@(0,0,z_{N_z})}(\mathbf{\kappa}, z_{N_z}) \end{bmatrix} \end{split}$$

• final step: back to (*x*,*y*) space

$$\Delta \varepsilon_r(x', y', z'_n) = \frac{\Delta \varepsilon_{r, \text{sp}} \Omega_{\text{sp}}}{\Omega_v} \mathcal{F}_{2D}^{-1} \left\{ \tilde{f}(\mathbf{\kappa}; z'_n) \right\}, n = 1, \dots, N_z$$

QUANTITATIVE SPM: SIMULATION EXAMPLE (F SHAPE)

[Nikolova, Introduction to Microwave Imaging, 2017]



QUANTITATIVE SPM: EXPERIMENTAL EXAMPLE



[Nikolova, Introduction to Microwave Imaging, 2017]

• linearized time-domain resolvent kernel

• resolvent kernel and PSF all over again: measure response with each antenna pair and position of a point scatterer at $\mathbf{r}' = 0$

TIME DOMAIN FORWARD MODEL WITH PSF – 2

• assuming uniform background: $\mathcal{K}_{\mathbf{r}'}(\mathbf{r}_{\mathrm{Rx}},\mathbf{r}_{\mathrm{Tx}};t) \approx \mathcal{K}_0(\mathbf{r}_{\mathrm{Rx}},\mathbf{r}_{\mathrm{Tx}};t - \Delta t(\mathbf{r}'))$



• example: analytical far-zone kernel some reference time $\mathcal{K}_{\mathbf{r}'}(\mathbf{r}_{\mathrm{Rx}}, \mathbf{r}_{\mathrm{Tx}}; t) \sim \mathrm{PSF}_{0}(\mathbf{r}_{\mathrm{Rx}}, \mathbf{r}_{\mathrm{Tx}}, t) \sim \delta\left(t - \frac{|\mathbf{r}' - \mathbf{r}_{\mathrm{Tx}}|}{v_{\mathrm{b}}} - \frac{|\mathbf{r}' - \mathbf{r}_{\mathrm{Rx}}|}{v_{\mathrm{b}}} - \frac{1}{v_{\mathrm{b}}}\right) / r^{2}$

DELAY AND SUM (DAS): THE CROSS-CORRELATION EXPLANATION

• cross-correlation – a measure of signal similarity

steering filter for antenna pair at **r** toward voxel **r**'

• with large number of responses, $X(\mathbf{r}',t) \sim \kappa(\mathbf{r}')$ as autocorrelation term dominates \mathbf{r}'' integral

autocorrelation term:
$$\kappa(\mathbf{r}') \cdot \sum_{\xi=1}^{N_T} \iint_{\mathbf{r} \in S_a} \mathrm{PSF}_{\xi,0}(t,\mathbf{r},\mathbf{r}') \otimes \mathrm{PSF}_{\xi,0}(t,\mathbf{r},\mathbf{r}') d\mathbf{r}$$

cross-correlation terms: $\kappa(\mathbf{r}'') \cdot \sum_{\xi=1}^{N_T} \iint_{\mathbf{r} \in S_a} \mathrm{PSF}_{\xi,0}(t,\mathbf{r},\mathbf{r}') \otimes \mathrm{PSF}_{\xi,0}(t,\mathbf{r},\mathbf{r}'') d\mathbf{r}$

DAS: SIGNAL-FLOW SCHEMATIC



DAS: CONCEPTUAL EXAMPLE



DAS: CONCEPTUAL EXAMPLE -2



CONCLUDING REMARKS

- we have just grazed the surface of an extensive subject
- real-time microwave imaging is rapidly growing and developing
 - hardware antennas & RF/radar electronics
 - calibration methods
 - inversion methods









IMAGE SPATIAL RESOLUTION – WHAT TO EXPECT

[Nikolova, Introduction to Microwave Imaging, 2017]

• lateral (or cross-range) resolution

$$\delta_{x,y} \ge \frac{\lambda_{\text{eff,min}}}{4} = \frac{\pi}{2k_{x,y}^{\text{max}}}$$

• depth (or range) resolution

$$\delta_{z} \geq \frac{\lambda_{\text{eff,min}}}{2} = \frac{\pi}{k_{z}^{\text{max}}} \approx \frac{v_{\text{b}}}{2B}$$

• wide viewing angles are critically important: wide-beam antennas, large scanned apertures

APPLICATIONS: NONDESTRUCTIVE TESTING

[Sheen et al., "Near-field three-dimensional radar imaging techniques and applications," Applied Optics 2010]

Pacific Northwest National Laboratory, Washington, USA



X-BAND (8 TO 12 GHz) SCANNER WITH 16×16 ELECTRONICALLY SWITCHED ARRAY: ABSORBER INSPECTION

APPLICATIONS: THROUGH-WALL IMAGING

[Depatla et al., "Robotic through-wall imaging," IEEE A&P Mag. 2017]

Prof. Mostofi's team at the University of California Santa Barbara



Area of interest – top view

3D binary ground-truth image of the unknown area to be imaged (2.96 m x 2.96 m x 0.4 m)

Our 3D image of the area, based on 3.84 % measurements





1.50 m

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LIMITATIONS OF BORN'S APPROXIMATION: TOTAL INTERNAL FIELD

• limitations on both the size and the contrast of the scatterer

$$a^2 \left| k_{\rm s}^2(\mathbf{r}) - k_{\rm b}^2 \right| \ll 1, \, \mathbf{r} \in V_{\rm s}$$



 if OUT violates the limits: images contain artifacts which reflect differences between E^{tot}_{Tx}(r') and E^{inc}_{Tx}(r') rather than contrast



[Nikolova, Introduction to Microwave Imaging, 2017]

NUMERICAL ASPECTS OF THE SOLUTION IN FOURIER SPACE

$$\tilde{S}_{\xi}^{(m)}(\kappa_{ij}) \approx \sum_{n=1}^{N_z} \tilde{f}(\kappa_{ij}; z'_n) \widetilde{\text{PSF}}_{0z',\xi}^{(m)}(\kappa_{ij}) \qquad \begin{array}{l} m = 1, \dots, N_{\omega} \\ \xi = 1, \dots, N_T \end{array}$$

$$\mathbf{K}_{ij} = (i\Delta k_x, j\Delta k_y)$$
$$i = 1, \dots, N_x; j = 1, \dots, N_y$$

solve $(N_x \cdot N_y)$ such systems of size: $N_T N_\omega \times N_z$

the data vector:

$$\mathbf{d}(\kappa_{ij}) = \left[\mathbf{d}_{1}^{T}(\kappa_{ij}) \cdots \mathbf{d}_{N_{T}}^{T}(\kappa_{ij})\right]_{N_{T}N_{\omega}\times 1}^{T}, \mathbf{d}_{\xi}^{T}(\kappa_{ij}) = \left[\tilde{S}_{\xi}^{(1)}(\kappa_{ij}) \cdots \tilde{S}_{\xi}^{(N_{\omega})}(\kappa_{ij})\right]_{N_{\omega}\times 1}^{T}, \xi = 1, \dots, N_{T}$$

the contrast vector:

$$\mathbf{f}(\kappa_{ij}) = [\tilde{f}(\kappa_{ij}; z'_1) \cdots \tilde{f}(\kappa_{ij}; z'_{N_z})]_{N_z \times \mathbb{Z}}^T$$

the system (PSF) matrix:

$$\mathbf{A}(\kappa_{ij}) = \begin{bmatrix} \mathbf{A}_{1}(\kappa_{ij}) \\ \vdots \\ \mathbf{A}_{N_{T}}(\kappa_{ij}) \end{bmatrix} \text{ where } \mathbf{A}_{\xi}(\kappa_{ij}) = \begin{bmatrix} \widetilde{\mathsf{PSF}}_{0z'_{1},\xi}^{(1)}(\kappa_{ij}) \cdots \widetilde{\mathsf{PSF}}_{0z'_{N_{z}},\xi}^{(1)}(\kappa_{ij}) \\ \vdots & \ddots & \vdots \\ \widetilde{\mathsf{PSF}}_{0z'_{1},\xi}^{(N_{\omega})}(\kappa_{ij}) \cdots \widetilde{\mathsf{PSF}}_{0z'_{N_{z}},\xi}^{(N_{\omega})}(\kappa_{ij}) \end{bmatrix}$$

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SYNTHETIC FOCUSING – 2



THE FORWARD MODEL OF HOLOGRAPHY

• the *S*-parameter forward model

$$\underbrace{S_{\xi}^{\mathrm{sc}}(x, y, \overline{z}; \omega)}_{\text{data}} \approx \iiint_{V_{\mathrm{s}}} \Delta \varepsilon_{r}(x', y', z') \Big[\overline{\mathbf{E}}_{\xi, \mathrm{Rx}}^{\mathrm{inc}} \cdot \overline{\mathbf{E}}_{\xi, \mathrm{Tx}}^{\mathrm{inc}} \Big]_{(x', y', z'; x, y, \overline{z}; \omega)} dx' dy' dz', \ \xi = 1, \dots, N_{T}$$

• ξ (response type) replaces (*i*,*k*)

TYPICAL RESPONSE TYPES

ξ	(<i>i</i> , <i>k</i>)	response
1	(1,1)	S_{11}
2	(2,2)	<i>S</i> ₂₂
3	(1,2) or (2,1)	$S_{12} = S_{21}$
$N_T = 3$		

THE FORWARD MODEL OF HOLOGRAPHY: PLANAR SCAN

-position of Tx/Rx antenna pair

• the *S*-parameter forward model



- $S_{\xi}^{\rm sc}(\underline{x, y, \overline{z}}; \omega) \approx \iiint_{V_{\varepsilon}} \Delta \varepsilon_r(\underline{x', y', z'}) \left[\overline{\mathbf{E}}_{\xi, \mathrm{Rx}}^{\rm inc} \cdot \overline{\mathbf{E}}_{\xi, \mathrm{Tx}}^{\rm inc} \right]_{(\underline{x', y', z'}; \underline{x}, y, \overline{z}; \omega)} dx' dy' dz', \ \xi = 1, \dots, N_T$
 - data

